# Robin double-phase problems with singular and superlinear terms 

Nikolaos S. Papageorgiou ${ }^{\text {a }}$, Vicenţiu D. Rădulescu ${ }^{\text {b,c,* }}$, Dušan D. Repovš ${ }^{\text {d,e }}$<br>${ }^{\text {a }}$ National Technical University, Department of Mathematics, Zografou Campus, Athens 15780, Greece<br>${ }^{\mathrm{b}}$ Faculty of Applied Mathematics, AGH University of Science and Technology, 30-059 Kraków, Poland<br>${ }^{\text {c }}$ Department of Mathematics, University of Craiova, 200585 Craiova, Romania<br>${ }^{\text {d }}$ Faculty of Education and Faculty of Mathematics and Physics, University of Ljubljana, 1000 Ljubljana, Slovenia<br>${ }^{\text {e }}$ Institute of Mathematics, Physics and Mechanics, 1000 Ljubljana, Slovenia

## A R T I C L E I N F O

## Article history:

Received 14 July 2020
Accepted 5 September 2020
Available online xxxx

## Keywords:

Nonhomogeneous differential operator
Nonlinear regularity theory
Truncation
Strong comparison principle
Positive solutions


#### Abstract

We consider a nonlinear Robin problem driven by the sum of $p$-Laplacian and $q$-Laplacian (i.e. the $(p, q)$-equation). In the reaction there are competing effects of a singular term and a parametric perturbation $\lambda f(z, x)$, which is Carathéodory and $(p-1)$-superlinear at $x \in \mathbb{R}$, without satisfying the Ambrosetti-Rabinowitz condition. Using variational tools, together with truncation and comparison techniques, we prove a bifurcation-type result describing the changes in the set of positive solutions as the parameter $\lambda>0$ varies.


(C) 2020 Elsevier Ltd. All rights reserved.

## 1. Introduction

Let $\Omega \subseteq \mathbb{R}^{N}$ be a bounded domain with a $C^{2}$-boundary $\partial \Omega$. In this paper, we study the following nonlinear Robin problem

$$
\left\{\begin{array}{c}
-\Delta_{p} u(z)-\Delta_{q} u(z)+\xi(z) u(z)^{p-1}=u(z)^{-\gamma}+\lambda f(z, u(z)) \text { in } \Omega, \\
\frac{\partial u}{\partial n_{p q}}+\beta(z) u^{p-1}=0 \text { on } \partial \Omega, u>0, \lambda>0,0<\gamma<1,1<q<p
\end{array}\right\}
$$

For every $r \in(1, \infty)$, we denote by $\Delta_{r}$ the $r$-Laplace differential operator defined by

$$
\Delta_{r} u=\operatorname{div}\left(|D u|^{r-2} D u\right) \text { for all } u \in W^{1, r}(\Omega) .
$$

The differential operator of $\left(P_{\lambda}\right)$ is the sum of $p$-Laplacian and $q$-Laplacian. Such an operator is not homogeneous and it appears in the mathematical models of various physical processes. We mention the

[^0]works of Cherfils \& Ilyasov [1] (reaction-diffusion systems) and Zhikov [2] (elasticity theory). The potential function $\xi \in L^{\infty}(\Omega)$ satisfies $\xi(z) \geqslant 0$ for almost all $z \in \Omega$. In the reaction (the right-hand side of $\left(P_{\lambda}\right)$ ), we have the combined effects of two nonlinearities of different nature. One nonlinearity is the singular term $u^{-\gamma}$ and the other nonlinearity is the parametric term $\lambda f(z, x)$, where $f(z, x)$ is a Carathéodory function (that is, for all $x \in \mathbb{R}$, the mapping $z \mapsto f(z, x)$ is measurable and for almost all $z \in \Omega$, the mapping $x \mapsto f(z, x)$ is continuous), which exhibits ( $p-1$ )-superlinear growth near $+\infty$ but without satisfying the usual in such cases Ambrosetti-Rabinowitz condition (the AR-condition for short). In the boundary condition, $\frac{\partial u}{\partial n_{p q}}$ denotes the conormal derivative corresponding to the ( $p, q$ )-Laplace differential operator. Then according to the nonlinear Green's identity (see Gasinski \& Papageorgiou [3, p. 210]), we have
$$
\frac{\partial u}{\partial n_{p q}}=\left(|D u|^{p-2} D u+|D u|^{q-2} D u, n\right) \text { for all } u \in C^{1}(\bar{\Omega})
$$
with $n(\cdot)$ being the outward unit normal on $\partial \Omega$. The boundary coefficient $\beta \in C^{0, \alpha}(\partial \Omega)$ (with $0<\alpha<1$ ) satisfies $\beta(z) \geqslant 0$ for all $z \in \partial \Omega$.

In the past, nonlinear singular problems were studied only in the context of Dirichlet equations driven by the $p$-Laplacian (a homogeneous differential operator). We mention the works of Giacomoni, Schindler \& Takač [4], Papageorgiou, Rădulescu \& Repovš [5,6], Papageorgiou \& Smyrlis [7], Papageorgiou \& Winkert [8], and Perera \& Zhang [9]. Nonlinear elliptic problems with unbalanced growth have been studied recently by Papageorgiou, Rădulescu and Repovš [10-12]. Double-phase transonic flow problems with variable growth have been considered by Bahrouni, Rădulescu and Repovš [13]. A comprehensive study of semilinear singular problems can be found in the book of Ghergu \& Rădulescu [14].

Using variational methods based on the critical point theory together with suitable truncation and comparison techniques, we prove a bifurcation type result, describing in a precise way the dependence of the set of positive solutions of $\left(P_{\lambda}\right)$ on the parameter. So, we produce a critical parameter value $\lambda^{*}>0$ such that for all $\lambda \in\left(0, \lambda^{*}\right)$, problem $\left(P_{\lambda}\right)$ has at least two positive solutions, for $\lambda=\lambda^{*}$ problem $\left(P_{\lambda}\right)$ has at least one positive solution and for $\lambda>\lambda^{*}$ there are no positive solutions for problem $\left(P_{\lambda}\right)$.

## 2. Mathematical background and hypotheses

Let $X$ be a Banach space. By $X^{*}$ we denote the topological dual of $X$. Given $\varphi \in C^{1}(X, \mathbb{R})$, we say that $\varphi(\cdot)$ satisfies the "C-condition", if the following property holds

> "Every sequence $\left\{u_{n}\right\}_{n \geqslant 1} \subseteq X$ such that
> $\left\{\varphi\left(u_{n}\right)\right\}_{n \geqslant 1} \subseteq \mathbb{R}$ is bounded and $\left(1+\left\|u_{n}\right\|\right) \varphi^{\prime}\left(u_{n}\right) \rightarrow 0$ in $X^{*}$ as $n \rightarrow \infty$, admits a strongly convergent subsequence."

This is a compactness type condition on the functional $\varphi$, which leads to the minimax theory of the critical values of $\varphi(\cdot)$.

The two main spaces in the analysis of problem $\left(P_{\lambda}\right)$ are the Sobolev space $W^{1, p}(\Omega)$ and the Banach space $C^{1}(\bar{\Omega})$. By $\|\cdot\|$ we denote the norm of the Sobolev space $W^{1, p}(\Omega)$. We have

$$
\|u\|=\left[\|u\|_{p}^{p}+\|D u\|_{p}^{p}\right]^{\frac{1}{p}} \text { for all } u \in W^{1, p}(\Omega)
$$

The Banach space $C^{1}(\bar{\Omega})$ is ordered with positive (order) cone given by

$$
C_{+}=\left\{u \in C^{1}(\bar{\Omega}): u(z) \geqslant 0 \text { for all } z \in \bar{\Omega}\right\} .
$$

This cone has a nonempty interior

$$
D_{+}=\left\{u \in C_{+}: u(z)>0 \text { for all } z \in \bar{\Omega}\right\} .
$$

We will also consider another order cone (closed convex cone) in $C^{1}(\bar{\Omega})$, namely the cone

$$
\hat{C}_{+}=\left\{u \in C^{1}(\bar{\Omega}): u(z) \geqslant 0 \text { for all } z \in \bar{\Omega},\left.\frac{\partial u}{\partial n}\right|_{\partial \Omega \cap u^{-1}(0)} \leqslant 0\right\}
$$

This cone has a nonempty interior

$$
\operatorname{int} \hat{C}_{+}=\left\{u \in C^{1}(\bar{\Omega}): u(z)>0 \text { for all } z \in \Omega,\left.\frac{\partial u}{\partial n}\right|_{\partial \Omega \cap u^{-1}(0)}<0\right\}
$$

To take care of the Robin boundary condition, we will also use the "boundary" Lebesgue spaces $L^{q}(\partial \Omega)(1 \leqslant q \leqslant \infty)$. More precisely, on $\partial \Omega$ we consider the $(N-1)$-dimensional Hausdorff (surface) measure $\sigma(\cdot)$. Using this measure on $\partial \Omega$ we can define in the usual way the Lebesgue spaces $L^{q}(\partial \Omega)(1 \leqslant q \leqslant \infty)$. We know that there exists a continuous, linear map $\gamma_{0}: W^{1, p}(\Omega) \rightarrow L^{p}(\partial \Omega)$, known as the "trace map" such that

$$
\gamma_{0}(u)=\left.u\right|_{\partial \Omega} \text { for all } u \in W^{1, p}(\Omega) \cap C(\bar{\Omega})
$$

So, the trace map extends the notion of boundary values to all Sobolev functions. We have

$$
\operatorname{im} \gamma_{0}=W^{\frac{1}{p^{\prime}}, p}(\partial \Omega)\left(\frac{1}{p}+\frac{1}{p^{\prime}}=1\right) \text { and } \operatorname{ker} \gamma_{0}=W_{0}^{1, p}(\Omega)
$$

The trace map $\gamma_{0}$ is compact into $L^{q}(\partial \Omega)$ for all $q \in\left[1, \frac{(N-1) p}{N-p}\right)$ if $N>p$ and into $L^{q}(\partial \Omega)$ for all $q \geqslant 1$ if $p \geqslant N$. In the sequel, for the sake of notational simplicity, we drop the use of the trace map $\gamma_{0}(\cdot)$. All restrictions of Sobolev functions on $\partial \Omega$ are understood in the sense of traces.

For every $r \in(1,+\infty)$, let $A_{r}: W^{1, r}(\Omega) \rightarrow W^{1, r}(\Omega)^{*}$ be defined by

$$
\left\langle A_{r}(u), h\right\rangle=\int_{\Omega}|D u|^{r-2}(D u, D h)_{\mathbb{R}^{N}} d z \text { for all } u, h \in W^{1, r}(\Omega)
$$

The following proposition summarizes the main properties of this map (see Gasinski \& Papageorgiou [3]).

Proposition 1. The map $A_{r}(\cdot)$ is bounded (that is, it maps bounded sets to bounded sets) continuous, monotone (hence maximal monotone, too) and of type $(S)_{+}$, that is, if $u_{n} \xrightarrow{w} u$ in $W^{1, r}(\Omega)$ and $\limsup \operatorname{sum}_{n \rightarrow \infty}\left\langle A_{r}\left(u_{n}\right), u_{n}-u\right\rangle$, then $u_{n} \rightarrow u$ in $W^{1, r}(\Omega)$.

Evidently, the $(S)_{+}$-property is useful in verifying the C-condition.
Now we introduce the conditions on the potential function $\xi(\cdot)$ and on the boundary coefficient $\beta(\cdot)$.
$H(\xi): \xi \in L^{\infty}(\Omega)$ and $\xi(z) \geqslant 0$ for almost all $z \in \Omega$.
$H(\beta): \beta \in C^{0, \alpha}(\partial \Omega)$ with $0<\alpha<1$ and $\beta(z) \geqslant 0$ for all $z \in \partial \Omega$.
$H_{0}: \xi \not \equiv 0$ or $\beta \not \equiv 0$.

Remark 1. When $\beta \equiv 0$ we have the usual Neumann problem.

The next two propositions can be found in Papageorgiou \& Rădulescu [15].

Proposition 2. If $\xi \in L^{\infty}(\Omega), \xi(z) \geqslant 0$ for almost all $z \in \Omega$ and $\xi \not \equiv 0$, then $c_{0}\|u\|^{p} \leqslant\|D u\|_{p}^{p}+$ $\int_{\Omega} \xi(z)|u|^{p} d z$ for some $c_{0}>0$ and all $u \in W^{1, p}(\Omega)$.

Proposition 3. If $\beta \in L^{\infty}(\partial \Omega), \beta(z) \geqslant 0$ for $\sigma$-almost all $z \in \partial \Omega$ and $\beta \not \equiv 0$, then $c_{1}\|u\|^{p} \leqslant$ $\|D u\|_{p}^{p}+\int_{\partial \Omega} \beta(z)|u|^{p} d \sigma$ for some $c_{1}>0$ and all $u \in W^{1, p}(\Omega)$.

In what follows, let $\gamma_{p}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ be defined by

$$
\gamma_{p}(u)=\|D u\|_{p}^{p}+\int_{\Omega} \xi(z)|u|^{p} d z+\int_{\partial \Omega} \beta(z)|u|^{p} d \sigma \text { for all } u \in W^{1, p}(\Omega) .
$$

If hypotheses $H(\xi), H(\beta), H_{0}$ hold, then from Propositions 2 and 3 we can infer that

$$
\begin{equation*}
c_{2}\|u\|^{p} \leqslant \gamma_{p}(u) \text { for some } c_{2}>0 \text { and all } u \in W^{1, p}(\Omega) . \tag{1}
\end{equation*}
$$

As we have already mentioned in the Introduction, our approach also involves truncation and comparison techniques. So, the next strong comparison principle, a slight variation of Proposition 4 of Papageorgiou \& Smyrlis [7], will be useful.

Proposition 4. If $\hat{\xi} \in L^{\infty}(\Omega)$ with $\hat{\xi}(z) \geqslant 0$ for almost all $z \in \Omega, h_{1}, h_{2} \in L^{\infty}(\Omega)$,

$$
0<c_{3} \leqslant h_{2}(z)-h_{1}(z) \text { for almost all } z \in \Omega,
$$

and the functions $u_{1}, u_{2} \in C^{1}(\bar{\Omega}) \backslash\{0\}, u_{1} \leqslant u_{2}, u_{1}^{-\gamma}, u_{2}^{-\gamma} \in L^{\infty}(\Omega)$ satisfy

$$
\begin{aligned}
& -\Delta_{p} u_{1}-\Delta_{q} u_{1}+\hat{\xi}(z) u_{1}^{p-1}-u_{1}^{-\gamma}=h_{1} \text { for almost all } z \in \Omega, \\
& -\Delta_{p} u_{2}-\Delta_{q} u_{2}+\hat{\xi}(z) u_{2}^{p-1}-u_{2}^{-\gamma}=h_{2} \text { for almost all } z \in \Omega,
\end{aligned}
$$

then $u_{2}-u_{1} \in \operatorname{int} \hat{C}_{+}$.
Consider a Carathéodory function $f_{0}: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ satisfying

$$
\left|f_{0}(z, x)\right| \leqslant a_{0}(z)\left[1+|x|^{r-1}\right] \text { for almost all } z \in \Omega \text { and all } x \in \mathbb{R},
$$

with $a_{0} \in L^{\infty}(\Omega)$ and $1<r \leqslant p^{*}=\left\{\begin{array}{ll}\frac{N p}{N-p} & \text { if } p<N \\ +\infty & \text { if } N \leqslant p\end{array}\right.$ (the critical Sobolev exponent corresponding to p).

We set $F_{0}(z, x)=\int_{0}^{x} f_{0}(z, s) d s$ and consider the $C^{1}$-functional $\varphi_{0}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\left.\varphi_{0}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} F_{0}(z, u) d z \text { for all } u \in W^{1, p}(\Omega) \text { (recall that } q<p\right) .
$$

The next proposition can be found in Papageorgiou \& Rădulescu [16] and essentially is an outgrowth of the nonlinear regularity theory of Lieberman [17].

Proposition 5. If $u_{0} \in W^{1, p}(\Omega)$ is a local $C^{1}(\bar{\Omega})$-minimizer of $\varphi_{0}$, that is, there exists $\rho_{0}>0$ such that

$$
\varphi_{0}\left(u_{0}\right) \leqslant \varphi_{0}\left(u_{0}+h\right) \text { for all }\|h\|_{C^{1}(\bar{\Omega})} \leqslant \rho_{0},
$$

then $u_{0} \in C^{1, \alpha}(\bar{\Omega})$ for some $\alpha \in(0,1)$ and $u_{0}$ is also a local $W^{1, p}(\Omega)$-minimizer of $\varphi_{0}$, that is, there exists $\rho_{1}>0$ such that

$$
\varphi_{0}\left(u_{0}\right) \leqslant \varphi_{0}(u+h) \text { for all }\|h\| \leqslant \rho_{1} .
$$

The next fact about ordered Banach spaces is useful in producing upper bounds for functions and can be found in Gasinski \& Papageorgiou [18, p. 680] (Problem 4.180).

Proposition 6. If $X$ is an ordered Banach space with positive (order) cone $K$,

$$
\operatorname{int} K \neq \emptyset \text { and } e \in \operatorname{int} K
$$

then for every $u \in X$ we can find $\lambda_{u}>0$ such that $\lambda_{u} e-u \in K$.

Under hypotheses $H(\xi), H(\beta), H_{0}$, the differential operator $u \mapsto-\Delta_{p} u+\xi(z)|u|^{p-2} u$ with the Robin boundary condition, has a principal eigenvalue $\hat{\lambda}_{1}(p)>0$ which is isolated, simple and admits the following variational characterization:

$$
\begin{equation*}
\hat{\lambda}_{1}(p)=\inf \left\{\frac{\gamma_{p}(u)}{\|u\|_{p}^{p}}: u \in W^{1, p}(\Omega), u \neq 0\right\} . \tag{2}
\end{equation*}
$$

The infimum is realized on the corresponding one-dimensional eigenspace, the elements of which have fixed sign. By $\hat{u}_{1}(p)$ we denote the positive, $L^{p}$-normalized (that is, $\left\|\hat{u}_{1}(p)\right\|_{p}=1$ ) eigenfunction corresponding to $\hat{\lambda}_{1}(p)>0$. The nonlinear Hopf theorem (see, for example, Gasinski \& Papageorgiou [3, p. 738]) implies that $\hat{u}_{1}(p) \in D_{+}$.

Let us fix some basic notation which we will use throughout this work. So, if $x \in \mathbb{R}$, we set $x^{ \pm}=$ $\max \{ \pm x, 0\}$ and the for $u \in W^{1, p}(\Omega)$ we define $u^{ \pm}(z)=u(z)^{ \pm}$for all $z \in \Omega$. We know that

$$
u^{ \pm} \in W^{1, p}(\Omega), u=u^{+}-u^{-},|u|=u^{+}+u^{-} .
$$

If $\varphi \in C^{1}\left(W^{1, p}(\Omega), \mathbb{R}\right)$, then by $K_{\varphi}$ we denote the critical set of $\varphi$, that is,

$$
K_{\varphi}=\left\{u \in W^{1, p}(\Omega): \varphi^{\prime}(u)=0\right\} .
$$

Also, if $u, y \in W^{1, p}(\Omega)$, with $u \leqslant y$, then we define

$$
\begin{aligned}
& {[u, y]=\left\{h \in W^{1, p}(\Omega): u(z) \leqslant h(z) \leqslant y(z) \text { for almost all } z \in \Omega\right\},} \\
& {[u)=\left\{h \in W^{1, p}(\Omega): u(z) \leqslant h(z) \text { for almost all } z \in \Omega\right\},} \\
& \operatorname{int}_{C^{1}(\bar{\Omega})}[u, y]=\text { the interior in the } C^{1}(\bar{\Omega}) \text {-norm of }[u, y] \cap C^{1}(\bar{\Omega}) .
\end{aligned}
$$

Now we introduce our hypotheses on the perturbation $f(z, x)$.
$H(f): f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function such that $f(z, 0)=0$ for almost all $z \in \Omega$ and
(i) $f(z, x) \leqslant a(z)\left(1+x^{r-1}\right)$ for almost all $z \in \Omega$ and all $x \geqslant 0$ with $a \in L^{\infty}(\Omega), p<r<p^{*}$;
(ii) if $F(z, x)=\int_{0}^{x} f(z, s) d s$, then $\lim _{x \rightarrow+\infty} \frac{F(z, x)}{x^{p}}=+\infty$ uniformly for almost all $z \in \Omega$;
(iii) there exists $\tau \in\left((r-p) \max \left\{\frac{N}{p}, 1\right\}, p^{*}\right)$ such that

$$
0<\hat{\beta}_{0} \leqslant \liminf _{x \rightarrow+\infty} \frac{f(z, x) x-p F(z, x)}{x^{\tau}} \text { uniformly for almost all } z \in \Omega ;
$$

(iv) for every $\vartheta>0$, there exists $m_{\vartheta}>0$ such that

$$
m_{\vartheta} \leqslant f(z, x) \text { for almost all } z \in \Omega \text { and all } x \geqslant \vartheta ;
$$

(v) for every $\rho>0$ and $\lambda>0$, there exists $\hat{\xi}_{\rho}^{\lambda}>0$ such that for almost all $z \in \Omega$, the function $x \mapsto f(z, x)+\hat{\xi}_{\rho}^{\lambda} x^{p-1}$ is nondecreasing on $[0, \rho]$.

Remark 2. Since we are looking for positive solutions and the above hypotheses concern the positive semiaxis, without any loss of generality we may assume that

$$
\begin{equation*}
f(z, x)=0 \text { for almost all } z \in \Omega \text { and all } x \leqslant 0 . \tag{3}
\end{equation*}
$$

From hypotheses $H(f),(i i),(i i i)$ it follows that

$$
\lim _{x \rightarrow+\infty} \frac{f(z, x)}{x^{p-1}}=+\infty \text { uniformly for almost all } z \in \Omega
$$

Hence, for almost all $z \in \Omega$, the perturbation $f(z, \cdot)$ is $(p-1)$-superlinear near $+\infty$. However, this superlinearity of $f(z, \cdot)$ is not expressed by using the well-known AR-condition. We recall that the AR-condition (unilateral version due to (3)) says that there exist $q>p$ and $M>0$ such that

$$
\begin{align*}
& 0<q F(z, x) \leqslant f(z, x) x \text { for almost all } z \in \Omega \text { and all } x \geqslant M,  \tag{4a}\\
& 0<\underset{\Omega}{\operatorname{ess} \inf } F(\cdot, M) . \tag{4b}
\end{align*}
$$

Integrating (4a) and using (4b), we obtain the following weaker condition

$$
\begin{aligned}
& c_{4} x^{q} \leqslant F(z, x) \text { for almost all } z \in \Omega \text { all } x \geqslant M, \text { and some } c_{4}>0, \\
\Rightarrow \quad & c_{4} x^{q-1} \leqslant f(z, x) \text { for almost all } z \in \Omega \text { and all } x \geqslant M .
\end{aligned}
$$

So, the AR-condition dictates at least $(q-1)$-polynomial growth for $f(z, \cdot)$. Here, we replace the ARcondition with hypothesis $H(f)(i i i)$ which is less restrictive and permits superlinear nonlinearities with "slower" growth near $+\infty$. For example, the function

$$
f(x)=x^{p-1} \ln (1+x) \text { for all } x \geqslant 0 .
$$

(for the sake of simplicity we have dropped the $z$-dependence) satisfies hypotheses $H(f)$, but fails to satisfy the AR-condition.

We introduce the following sets:

$$
\begin{aligned}
& \mathcal{L}=\left\{\lambda>0: \text { problem }\left(P_{\lambda}\right) \text { has a positive solution }\right\}, \\
& S_{\lambda}=\text { the set of positive solutions of }\left(P_{\lambda}\right) .
\end{aligned}
$$

Also we set

$$
\lambda^{*}=\sup \mathcal{L} .
$$

## 3. Some auxiliary Robin problems

Let $\eta>0$. First, we examine the following auxiliary Robin problem

$$
\left\{\begin{array}{l}
-\Delta_{p} u(z)-\Delta_{q} u(z)+\xi(z) u(z)^{p-1}=\eta \text { in } \Omega,  \tag{6}\\
\frac{\partial u}{\partial n_{p q}}+\beta(z) u^{p-1}=0 \text { on } \partial \Omega, u>0 .
\end{array}\right\}
$$

Proposition 7. If hypotheses $H(\xi), H(\beta), H_{0}$ hold, then for every $\eta>0$ problem (6) has a unique solution $\tilde{u}_{\eta} \in D_{+}$, the mapping $\eta \mapsto \tilde{u}_{\eta}$ is strictly increasing (that is, $\eta<\eta^{\prime} \Rightarrow \tilde{u}_{\eta^{\prime}}-\tilde{u}_{\eta} \in \operatorname{int} \hat{C}_{+}$) and

$$
\tilde{u}_{\eta} \rightarrow 0 \text { in } C^{1}(\bar{\Omega}) \text { as } \eta \rightarrow 0^{+} .
$$

Proof. Consider the map $V: W^{1, p}(\Omega) \rightarrow W^{1, p}(\Omega)^{*}$ defined by

$$
\begin{align*}
& \langle V(u), h\rangle=\left\langle A_{p}(u), h\right\rangle+\left\langle A_{q}(u), h\right\rangle+\int_{\Omega} \xi(z)|u|^{p-2} u h d z+\int_{\partial \Omega} \beta(z)|u|^{p-2} u h d \sigma  \tag{7}\\
& \quad \text { for all } u, h \in W^{1, p}(\Omega) .
\end{align*}
$$

Evidently, $V(\cdot)$ is continuous, strictly monotone (hence maximal monotone, too) and coercive (see (1)). Therefore $V(\cdot)$ is surjective (see Gasinski \& Papageorgiou [3, Corollary 3.2.31, p. 319]). So, we can find $\tilde{u}_{\eta} \in W^{1, p}(\Omega), \tilde{u}_{\eta} \neq 0$ such that

$$
V\left(\tilde{u}_{\eta}\right)=\eta .
$$

The strict monotonicity of $V(\cdot)$ implies that $\tilde{u}_{\eta}$ is unique. We have

$$
\begin{equation*}
\left\langle V\left(\tilde{u}_{\eta}\right), h\right\rangle=\eta \int_{\Omega} h d z \text { for all } h \in W^{1, p}(\Omega) . \tag{8}
\end{equation*}
$$

In (8) we choose $h=-\tilde{u}_{\eta}^{-} \in W^{1, p}(\Omega)$. Then

$$
\begin{aligned}
& c_{2}\left\|\tilde{u}_{\eta}^{-}\right\|^{p} \leqslant 0(\text { see }(1)), \\
\Rightarrow \quad & \tilde{u}_{\eta} \geqslant 0, \tilde{u}_{\eta} \neq 0 .
\end{aligned}
$$

From (8) we have

$$
\left\{\begin{array}{l}
-\Delta_{p} \tilde{u}_{\eta}(z)-\Delta_{q} \tilde{u}_{\eta}(z)+\xi(z) \tilde{u}_{\eta}(z)^{p-1}=\eta \text { for almost all } z \in \Omega,  \tag{9}\\
\frac{\partial \tilde{u}_{\eta}}{\partial n_{p q}}+\beta(z) \tilde{u}_{\eta}^{p-1}=0 \text { on } \partial \Omega .
\end{array}\right\}
$$

From (9) and Proposition 7 of Papageorgiou \& Rădulescu [16] we deduce that

$$
\tilde{u}_{\eta} \in L^{\infty}(\Omega) .
$$

Then the nonlinear regularity theory of Lieberman [17] implies that

$$
\tilde{u}_{\eta} \in C_{+} \backslash\{0\} .
$$

From (9) we have

$$
\begin{aligned}
& \Delta_{p} \tilde{u}_{\eta}(z)+\Delta_{q} \tilde{u}_{\eta}(z) \leqslant\|\xi\|_{\infty} \tilde{u}_{\eta}(z)^{p-1} \text { for almost all } z \in \Omega, \\
\Rightarrow \quad & \tilde{u}_{\eta} \in D_{+}(\text {see Pucci \& Serrin }[19, \text { pp. 111, 120] }) .
\end{aligned}
$$

Suppose that $0<\eta_{1}<\eta_{2}$ and let $\tilde{u}_{\eta_{1}}, \tilde{\eta}_{\eta_{2}} \in D_{+}$be the corresponding solutions of problem (6). We have

$$
\begin{aligned}
& -\Delta_{p} \tilde{u}_{\eta_{1}}-\Delta_{q} \tilde{u}_{\eta_{1}}+\xi(z) \tilde{u}_{\eta_{1}}^{p-1}=\eta_{1}<\eta_{2}=-\Delta_{p} \tilde{u}_{\eta_{2}}-\Delta_{q} \tilde{u}_{\eta_{2}}+\xi(z) \tilde{u}_{\eta_{2}} \\
& \text { for almost all } z \in \Omega, \\
\Rightarrow & \tilde{u}_{\eta_{2}}-\tilde{u}_{\eta_{1}} \in \operatorname{int} \hat{C}_{+} \text {(see Proposition 4), } \\
\Rightarrow & \eta \mapsto \tilde{u}_{\eta} \text { is strictly increasing from }(0,+\infty) \text { into } C^{1}(\bar{\Omega}) .
\end{aligned}
$$

Finally, let $\eta_{n} \rightarrow 0^{+}$and let $\tilde{u}_{n}=\tilde{u}_{\eta_{n}} \in D_{+}$be the corresponding solutions of (6). As before, invoking Proposition 7 of Papageorgiou \& Rădulescu [16], we can find $c_{5}>0$ such that

$$
\left\|\tilde{u}_{n}\right\|_{\infty} \leqslant c_{5} \text { for all } n \in \mathbb{N}
$$

Then from Lieberman [17] we infer that there exist $\alpha \in(0,1)$ and $c_{6}>0$ such that

$$
\tilde{u}_{n} \in C^{1, \alpha}(\bar{\Omega}),\left\|\tilde{u}_{n}\right\|_{C^{1, \alpha}(\bar{\Omega})} \leqslant c_{6} \text { for all } n \in \mathbb{N} .
$$

Exploiting the compact embedding of $C^{1, \alpha}(\bar{\Omega})$ into $C^{1}(\bar{\Omega})$, the monotonicity of the sequence $\left\{\tilde{u}_{n}\right\}_{n \geqslant 1} \subseteq$ $D_{+}$and the fact that for $\eta=0, u \equiv 0$ is the only solution of (6) we obtain

$$
\tilde{u}_{n} \rightarrow 0 \text { in } C^{1}(\bar{\Omega}) .
$$

The proof is now complete.
Using Proposition 7, we see that we can find $\eta_{0}>0$ such that

$$
\begin{equation*}
\eta \leqslant \tilde{u}_{\eta}(z)^{-\gamma} \text { for all } z \in \bar{\Omega}, \quad 0<\eta \leqslant \eta_{0} . \tag{10}
\end{equation*}
$$

We consider the following purely singular problem

$$
\left\{\begin{array}{c}
-\Delta_{p} u(z)-\Delta_{q} u(z)+\xi(z) u(z)^{p-1}=u(z)^{-\gamma} \text { in } \Omega,  \tag{11}\\
\frac{\partial u}{\partial n_{p q}}+\beta(z) u^{p-1}=0 \text { on } \partial \Omega, u>0,0<\gamma<1 .
\end{array}\right\}
$$

In the first place, by a solution of (11) we understand a weak solution, that is, a function $u \in W^{1, p}(\Omega)$ such that

$$
\begin{aligned}
& u^{-\gamma} h \in L^{1}(\Omega) \text { and }\left\langle A_{p}(u), h\right\rangle+\left\langle A_{q}(u), h\right\rangle+\int_{\Omega} \xi(z) u^{p-1} h d z+\int_{\partial \Omega} \beta(z) u^{p-1} h d \sigma \\
& =\int_{\Omega} u^{-\gamma} h d z \text { for all } h \in W^{1, p}(\Omega)
\end{aligned}
$$

In fact, using the nonlinear regularity theory, we will be able to establish more regularity for the solution of (11), which in fact, is a strong solution (that is, the equation can be interpreted pointwise almost everywhere on $\Omega$ ).

Proposition 8. If hypotheses $H(\xi), H(\beta), H_{0}$ hold, then problem (11) admits a unique solution $v \in D_{+}$.

Proof. Let $\eta \in\left(0, \eta_{0}\right]$ (see (10)) and recall that $\tilde{u}_{\eta} \in D_{+}$. So $m_{\eta}=\min _{\bar{\Omega}} \tilde{u}_{\eta}>0$ and

$$
\begin{align*}
& \eta \leqslant \tilde{u}_{\eta}^{-\gamma} \leqslant m_{\eta}^{-\gamma}(\text { see }(10)) \\
\Rightarrow & \tilde{u}_{\eta}^{-\gamma} \in L^{\infty}(\Omega) \tag{12}
\end{align*}
$$

We consider the following truncation of the reaction in problem (11):

$$
k(z, x)= \begin{cases}\tilde{u}_{\eta}(z)^{-\gamma} & \text { if } x \leqslant \tilde{u}_{\eta}(z)  \tag{13}\\ x^{-\gamma} & \text { if } \tilde{u}_{\eta}(z)<x\end{cases}
$$

This is a Carathéodory function. We set $K(z, x)=\int_{0}^{x} k(z, s) d s$ and consider the $C^{1}$-functional $\Psi$ : $W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\Psi(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} K(z, u) d z \text { for all } u \in W^{1, p}(\Omega)
$$

From (12) and (13), we see that $\Psi(\cdot)$ is coercive. Also the Sobolev embedding theorem and the compactness of the trace map, imply that $\Psi(\cdot)$ is sequentially weakly lower semicontinuous. So, we can find $v \in W^{1, p}(\Omega)$ such that

$$
\begin{gather*}
\Psi(v)=\inf \left\{\Psi(u): u \in W^{1, p}(\Omega)\right\} \\
\Rightarrow \quad \Psi^{\prime}(v)=0 \\
\Rightarrow\left\langle A_{p}(v), h\right\rangle+\left\langle A_{q}(v), h\right\rangle+\int_{\Omega} \xi(z)|v|^{p-2} v h d z+\int_{\partial \Omega} \beta(z)|v|^{p-2} v h d \sigma= \\
\quad \int_{\Omega} k(z, v) h d z \text { for all } h \in W^{1, p}(\Omega) \tag{14}
\end{gather*}
$$

In (14) we choose $\left(\tilde{u}_{\eta}-v\right)^{+} \in W^{1, p}(\Omega)$. Then

$$
\begin{align*}
& \left\langle A_{p}(v),\left(\tilde{u}_{\eta}-v\right)^{+}\right\rangle+\left\langle A_{q}(v),\left(\tilde{u}_{\eta}-v\right)^{+}\right\rangle+\int_{\Omega} \xi(z)|v|^{p-2} v\left(\tilde{u}_{\eta}-v\right)^{+} d z+ \\
& \int_{\partial \Omega} \beta(z)|v|^{p-2} v\left(\tilde{u}_{\eta}-v\right)^{+} d \sigma=\int_{\Omega} \tilde{u}_{\eta}^{-\gamma}\left(\tilde{u}_{\eta}-v\right)^{+} d z(\text { see }(13)) \\
\geqslant & \int_{\Omega} \eta\left(\tilde{u}_{\eta}-v\right)^{+} d z\left(\text { see }(10) \text { and recall that } 0<\eta \leqslant \eta_{0}\right) \\
= & \left\langle A_{p}\left(\tilde{u}_{\eta}\right),\left(\tilde{u}_{\eta}-v\right)^{+}\right\rangle+\left\langle A_{q}\left(\tilde{u}_{\eta}\right),\left(\tilde{u}_{\eta}-v\right)^{+}\right\rangle+\int_{\Omega} \xi(z) \tilde{u}_{\eta}^{p-1}\left(\tilde{u}_{\eta}-v\right)^{+} d z+ \\
& \int_{\partial \Omega} \beta(z) \tilde{u}_{\eta}^{p-1}\left(\tilde{u}_{\eta}-v\right)^{+} d \sigma(\text { see Proposition } 7) \\
\Rightarrow & \tilde{u}_{\eta} \leqslant v \tag{15}
\end{align*}
$$

Then from (13), (14), (15) we obtain

$$
\left\{\begin{array}{l}
-\Delta_{p} v(z)-\Delta_{q} v(z)+\xi(z) v(z)^{p-1}=v(z)^{-\gamma} \text { for almost all } z \in \Omega,  \tag{16}\\
\frac{\partial v}{\partial n_{p q}}+\beta(z) v^{p-1}=0 \text { on } \partial \Omega \\
\text { (see Papageorgiou \& Rădulescu [20]). }
\end{array}\right\}
$$

From (15) we have $v^{-\gamma} \leqslant \tilde{u}_{\eta}^{-\gamma} \in L^{\infty}(\Omega)$ (see (12)). So, from (16) and [16] we have $v \in L^{\infty}(\Omega)$. Then the nonlinear regularity theory of Lieberman [17] implies that $v \in C_{+}$. Hence it follows from (15) that

$$
v \in D_{+}
$$

Next, we show that this positive solution is unique. To this end, let $\hat{v} \in W^{1, p}(\Omega)$ be another positive solution of (11). Again we have $\hat{v} \in D_{+}$. Then

$$
\begin{aligned}
& \left\langle A_{p}(v),(\hat{v}-v)^{+}\right\rangle+\left\langle A_{q}(v),(\hat{v}-v)^{+}\right\rangle+\int_{\Omega} \xi(z) v^{p-1}(\hat{v}-v)^{+} d z+ \\
& \int_{\partial \Omega} \beta(z) v^{p-1}(\hat{v}-v)^{+} d \sigma \\
= & \int_{\Omega} v^{-\gamma}(\hat{v}-v)^{+} d z \\
\geqslant & \int_{\Omega} \hat{v}^{-\gamma}(\hat{v}-v)^{+} d z \\
= & \left\langle A_{p}(\hat{v}),(\hat{v}-v)^{+}\right\rangle+\left\langle A_{q}(\hat{v}),(\hat{v}-v)^{+}\right\rangle+\int_{\Omega} \xi(z) \hat{v}^{p-1}(\hat{v}-v)^{+} d z+ \\
& \int_{\hat{v}} \beta(z) \hat{v}^{p-1}(\hat{v}-v)^{+} d \sigma \\
\Rightarrow & v .
\end{aligned}
$$

Interchanging the roles of $v$ and $\hat{v}$ in the above argument, we obtain

$$
\begin{aligned}
& v \leqslant \hat{v} \\
\Rightarrow \quad & v=\hat{v}
\end{aligned}
$$

This proves the uniqueness of the positive solution of the purely singular problem (11).
Next, we consider the following nonlinear Robin problem

$$
\left\{\begin{array}{l}
-\Delta_{p} u(z)-\Delta_{q} u(z)+\xi(z) u(z)^{p-1}=v(z)^{-\gamma}+1 \text { in } \Omega  \tag{17}\\
\frac{\partial u}{\partial n_{p q}}+\beta(z) u^{p-1}=0 \text { on } \partial \Omega, u>0
\end{array}\right\}
$$

Proposition 9. If hypotheses $H(\xi), H(\beta), H_{0}$ hold, then problem (17) admits a unique solution $\bar{u} \in D_{+}$ and $v \leqslant \bar{u}$.

Proof. We know that $v^{-\gamma} \in L^{\infty}(\Omega)$ (see (12) and (15)). Then the existence and uniqueness of the solution $\bar{u} \in W^{1, p}(\Omega) \backslash\{0\}, \bar{u} \geqslant 0$ of (17) follow from the surjectivity and strict monotonicity of the map $V(\cdot)$ (see the proof of Proposition 7). The nonlinear regularity theory and the nonlinear Hopf's theorem imply that $\bar{u} \in D_{+}$.

Moreover, we have

$$
\begin{aligned}
& \left\langle A_{p}(\bar{u}),(v-\bar{u})^{+}\right\rangle+\left\langle A_{q}(\bar{u}),(v-\bar{u})^{+}\right\rangle+\int_{\Omega} \xi(z) \bar{u}^{p-1}(v-\bar{u})^{+} d z+ \\
& \int_{\partial \Omega} \beta(z) \bar{u}^{p-1}(v-\bar{u})^{+} d \sigma
\end{aligned}
$$

$$
\begin{aligned}
= & \int_{\Omega}\left[v^{-\gamma}+1\right](v-\bar{u})^{+} d z(\text { see }(17)) \\
\geqslant & \int_{\Omega} v^{-\gamma}(v-\bar{u})^{+} d z \\
= & \left\langle A_{p}(v),(v-\bar{u})^{+}\right\rangle+\left\langle A_{q}\left(v,(v-\bar{v})^{+}\right)\right\rangle+\int_{\Omega} \xi(z) v^{p-1}(v-\bar{v})^{+} d z+ \\
& \int_{\partial \Omega} \beta(z) v^{p-1}(v-\bar{v})^{+} d \sigma \\
\Rightarrow & v \leqslant \bar{u} .
\end{aligned}
$$

The proof is now complete.

## 4. Positive solutions

In this section we prove the bifurcation-type theorem described in the Introduction.
Proposition 10. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold, then $\mathcal{L} \neq \emptyset$ and $S_{\lambda} \subseteq D_{+}$.

Proof. Let $v \in D_{+}$be the unique positive solution of the auxiliary problem (11) (see Proposition 8) and $\bar{u} \in D_{+}$the unique solution of (17) (see Proposition 9). We know that $v \leqslant \bar{u}$ (see Proposition 9). Since $\bar{u} \in D_{+}$, hypothesis $H(f)(i)$ implies that

$$
0 \leqslant f(z, \bar{u}(z)) \leqslant c_{7} \text { for some } c_{7}>0 \text { and almost all } z \in \Omega .
$$

So, we can find $\lambda_{0}>0$ so small that

$$
\begin{equation*}
0 \leqslant \lambda f(z, \bar{u}(z)) \leqslant 1 \text { for almost all } z \in \Omega \text { and all } 0<\lambda \leqslant \lambda_{0} . \tag{18}
\end{equation*}
$$

We consider the following truncation of the reaction in problem $\left(P_{\lambda}\right)$

$$
\vartheta_{\lambda}(z, x)= \begin{cases}v(z)^{-\gamma}+\lambda f(z, v(z)) & \text { if } x<v(z)  \tag{19}\\ x^{-\gamma}+\lambda f(z, x) & \text { if } v(z) \leqslant x \leqslant \bar{u}(z) \\ \bar{u}(z)^{-\gamma}+\lambda f(z, \bar{u}(z)) & \text { if } \bar{u}(z)<x .\end{cases}
$$

This is a Carathéodory function. We set $\theta_{\lambda}(z, x)=\int_{0}^{x} \vartheta_{\lambda}(z, s) d s$ and consider the functional $\mu_{\lambda}$ : $W^{1, p}(\Omega) \rightarrow \mathbb{R}\left(\lambda \in\left(0, \lambda_{0}\right]\right)$ defined by

$$
\mu_{\lambda}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} \theta_{\lambda}(z, u) d z \text { for all } u \in W^{1, p}(\Omega) .
$$

Since $0 \leqslant \bar{u}^{-\gamma} \leqslant v^{-\gamma} \in L^{\infty}(\Omega)$, we see that $\mu_{\lambda} \in C^{1}\left(W^{1, p}(\Omega)\right)$. Also, it is clear from (1) and (19), that $\mu_{\lambda}(\cdot)$ is coercive. In addition, it is sequentially weakly lower semicontinuous. So, we can find $u_{\lambda} \in W^{1, p}(\Omega)$ such that

$$
\begin{align*}
& \mu_{\lambda}\left(u_{\lambda}\right)=\inf \left\{\mu_{\lambda}(u): u \in W^{1, p}(\Omega)\right\}, \\
& \Rightarrow \\
\Rightarrow & \mu_{\lambda}^{\prime}\left(u_{\lambda}\right)=0, \\
= & \left.A_{p}\left(u_{\lambda}\right), h\right\rangle+\left\langle A_{q}\left(u_{\lambda}\right), h\right\rangle+\int_{\Omega} \xi(z)\left|u_{\lambda}\right|^{p-2} u_{\lambda} h d z+\int_{\partial \Omega} \beta(z)\left|u_{\lambda}\right|^{p-2} u_{\lambda} h d \sigma  \tag{20}\\
= & \int_{\Omega} \vartheta_{\lambda}\left(z, u_{\lambda}\right) h d z \text { for all } h \in W^{1, p}(\Omega) .
\end{align*}
$$

In (20) first we choose $h=\left(u_{\lambda}-\bar{u}\right)^{+} \in W^{1, p}(\Omega)$. Then

$$
\begin{aligned}
& \left\langle A_{p}\left(u_{\lambda}\right),\left(u_{\lambda}-\bar{u}\right)^{+}\right\rangle+\left\langle A_{q}\left(u_{\lambda}\right),\left(u_{\lambda}-\bar{u}\right)^{+}\right\rangle+\int_{\Omega} \xi(z) u_{\lambda}^{p+}\left(u_{\lambda}-\bar{u}\right)^{+} d z+ \\
& \int_{\partial \Omega} \beta(z) u_{\lambda}^{p-1}\left(u_{\lambda}-\bar{u}\right) d \sigma \\
& \left.=\int_{\Omega}\left[\bar{u}^{-\gamma}+\lambda f(z, \bar{u})\right]\left(u_{\lambda}-\bar{u}\right)^{+} d z(\text { see }(19))\right) \\
& \leqslant \int_{\Omega}\left[\bar{u}^{-\gamma}+1\right]\left(u_{\lambda}-\bar{u}\right)^{+} d z(\text { see (18)) } \\
& \leqslant \int_{\Omega}\left[v^{-\gamma}+1\right]\left(u_{\lambda}-\bar{u}\right)^{+} d z(\text { since } v \leqslant \bar{u}) \\
& =\left\langle A_{p}(\bar{u}),\left(u_{\lambda}-\bar{u}\right)^{+}\right\rangle+\left\langle A_{q}(\bar{u}),\left(u_{\lambda}-\bar{u}\right)^{+}\right\rangle+\int_{\Omega} \xi(z) \bar{u}^{p-1}\left(u_{\lambda}-\bar{u}\right)^{+} d z \\
& +\int_{\partial \Omega} \beta(z) \bar{u}^{p-1}\left(u_{\lambda}-\bar{u}\right)^{+} d \sigma(\text { see Proposition } 9), \\
& \Rightarrow u_{\lambda} \leqslant \bar{u} .
\end{aligned}
$$

Next, in (20) we choose $h=\left(v-u_{\lambda}\right)^{+} \in W^{1, p}(\Omega)$. Then

$$
\begin{aligned}
& \left\langle A_{p}\left(u_{\lambda}\right),\left(v-u_{\lambda}\right)^{+}\right\rangle+\left\langle A_{q}\left(u_{\lambda}\right),\left(v-u_{\lambda}\right)^{+}\right\rangle+\int_{\Omega} \xi(z)\left|u_{\lambda}\right|^{p-2} u_{\lambda}\left(v-u_{\lambda}\right)^{+} d z+ \\
& \int_{\partial \Omega} \beta(z)\left|u_{\lambda}\right|^{p-2} u_{\lambda}\left(v-u_{\lambda}\right)^{+} d \sigma \\
& =\int_{\Omega}\left[v^{-\gamma}+\lambda f(z, v)\right]\left(v-u_{\lambda}\right)^{+} d z(\text { see }(19)) \\
& \geqslant \int_{\Omega} v^{-\gamma}\left(v-u_{\lambda}\right)^{+} d z(\text { since } f \geqslant 0) \\
& =\left\langle A_{p}(v),\left(v-u_{\lambda}\right)^{+}\right\rangle+\left\langle A_{q}(v),\left(v-u_{\lambda}\right)^{+}\right\rangle+\int_{\lambda} \xi(z) v^{p-1}\left(v-u_{\lambda}\right)^{+} d z \\
& +\int_{\partial \Omega} \beta(z) v^{p-1}\left(v-u_{\lambda}\right)^{+} d \sigma(\text { see Proposition } 8), \\
& \Rightarrow v \leqslant u_{\lambda} .
\end{aligned}
$$

So, we have proved that

$$
\begin{equation*}
u_{\lambda} \in[v, \bar{u}] . \tag{21}
\end{equation*}
$$

From (19), (20), (21) it follows that

$$
\left\{\begin{array}{l}
-\Delta_{p} u_{\lambda}(z)-\Delta_{q} u_{\lambda}(z)+\xi(z) u_{\lambda}(z)^{p-1}=u_{\lambda}(z)^{-\gamma}+\lambda f\left(z, u_{\lambda}(z)\right)  \tag{22}\\
\text { for almost all } z \in \Omega, \\
\frac{\partial u_{\lambda}}{\partial n_{p q}}+\beta(z) u_{\lambda}^{p-1}=0 \text { on } \partial \Omega,(\text { see }[20]) .
\end{array}\right\}
$$

By (22) and Proposition 7 of Papageorgiou \& Rădulescu [16], we have that $u_{\lambda} \in L^{\infty}(\Omega)$. So, the nonlinear regularity theory of Lieberman [17] implies that $u_{\lambda} \in D_{+}($see (21)). Therefore we have proved that

$$
\left(0, \lambda_{0}\right] \leqslant \mathcal{L} \neq \emptyset \text { and } S_{\lambda} \subseteq D_{+} .
$$

The proof is now complete.

Next, we establish a lower bound for the elements of $S_{\lambda}$.

Proposition 11. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold, $\lambda \in \mathcal{L}$ and $u \in S_{\lambda}$, then $v \leqslant u$.

Proof. From Proposition 10 we know that $u \in D_{+}$. Then Proposition 7 implies that for $\eta>0$ small enough, we have $\tilde{u}_{\eta} \leqslant u$. So, we can define the following Carathéodory function

$$
e(z, x)= \begin{cases}\tilde{u}_{\eta}(z)^{-\gamma} & \text { if } x<\tilde{u}_{\eta}(z)  \tag{23}\\ x^{-\gamma} & \text { if } \tilde{u}_{\eta}(z) \leqslant x \leqslant u(z) \\ u(z)^{-\gamma} & \text { if } u(z)<x\end{cases}
$$

We set $E(z, x)=\int_{0}^{x} e(z, s) d s$ and consider the functional $d: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
d(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} E(z, u) d z \text { for all } u \in W^{1, p}(\Omega)
$$

As before, we have $d \in C^{1}\left(W^{1, p}(\Omega)\right)$. Also, $d(\cdot)$ is coercive (see (23)) and weakly lower semicontinuous. Hence, we can find $\hat{v} \in W^{1, p}(\Omega)$ such that

$$
\begin{align*}
& d(\hat{u})=\inf \left\{d(u): u \in W^{1, p}(\Omega)\right\} \\
\Rightarrow & d^{\prime}(\hat{v})=0 \\
\Rightarrow \quad & \left\langle A_{p}(\hat{v}), h\right\rangle+\left\langle A_{q}(\hat{v}), h\right\rangle+\int_{\Omega} \xi(z)|\hat{v}|^{p-2} \hat{v} h d z+\int_{\partial \Omega} \beta(z)|\hat{v}|^{p-2} \hat{v} h d \sigma=  \tag{24}\\
& \int_{\Omega} e(z, \hat{v}) h d z \text { for all } h \in W_{1, p}(\Omega)
\end{align*}
$$

In (24) first we choose $h=(\hat{v}-u)^{+} \in W^{1, p}(\Omega)$. Exploiting the fact that $u \in S_{\lambda}$ and recalling that $f \geqslant 0$, we obtain $\hat{v} \leqslant u$. Next, in (24) we test with $h=\left(\tilde{u}_{\eta}-v\right)^{+} \in W^{1, p}(\Omega)$. Using (23), (10) and Proposition 7, we obtain $\tilde{u}_{\eta} \leqslant \hat{v}$. Therefore

$$
\begin{equation*}
\hat{v} \in\left[\tilde{u}_{\eta}, u\right] \tag{25}
\end{equation*}
$$

From (23), (24), (25) and Proposition 8, we conclude that

$$
\begin{aligned}
& \hat{v}=v \\
\Rightarrow \quad & v \leqslant u \text { for all } u \in S_{\lambda}
\end{aligned}
$$

The proof is now complete.

Now we can deduce a structural property of $\mathcal{L}$.

Proposition 12. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold, $\lambda \in \mathcal{L}, 0<\mu<\lambda$ and $u_{\lambda} \in S_{\lambda} \subseteq D_{+}$, then $\mu \in \mathcal{L}$ and we can find $u_{\mu} \in S_{\mu} \subseteq D_{+}$such that $u_{\lambda}-u_{\mu} \in \operatorname{int} \hat{C}_{+}$.

Proof. From Proposition 11 we know that $v \leqslant u_{\lambda}$. Therefore we can define the following Carathéodory function

$$
\hat{k}_{\mu}(z, x)= \begin{cases}v(z)^{-\gamma}+\mu f(z, v(z)) & \text { if } x<v(z)  \tag{26}\\ x^{-\gamma}+\mu f(z, x) & \text { if } v(z) \leqslant x \leqslant u_{\lambda}(z) \\ u_{\lambda}(z)^{-\gamma}+\mu f\left(z, u_{\lambda}(z)\right) & \text { if } u_{\lambda}(z)<x\end{cases}
$$

We set $\hat{K}_{\mu}(z, x)=\int_{0}^{x} \hat{k}_{\mu}(z, s) d s$ and consider the $C^{1}$-functional $\hat{\Psi}_{\mu}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\hat{\Psi}_{\mu}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} \hat{K}_{\mu}(z, u) d z \text { for all } u \in W^{1, p}(\Omega)
$$

Evidently, $\hat{\Psi}_{\mu}(\cdot)$ is coercive (see (26)) and sequentially weakly lower semicontinuous. So, we can find $u_{\mu} \in W^{1, p}(\Omega)$ such that

$$
\begin{align*}
& \hat{\Psi}_{\mu}\left(u_{\mu}\right)=\inf \left\{\hat{\Psi}_{\mu}(u): u \in W^{1, p}(\Omega)\right\}, \\
\Rightarrow & \hat{\Psi}_{\mu}^{\prime}\left(u_{\mu}\right)=0 \\
\Rightarrow & \left\langle A_{p}\left(u_{\mu}\right), h\right\rangle+\left\langle A_{q}\left(u_{\mu}\right), h\right\rangle+\int_{\Omega} \xi(z)\left|u_{\mu}\right|^{p-2} u_{\mu} h d z+\int_{\partial \Omega} \beta(z)\left|u_{\mu}\right|^{p-2} u_{\mu} h d \sigma \\
= & \int_{\Omega} \hat{k}_{\mu}\left(z, u_{\mu}\right) h d z \text { for all } h \in W^{1, p}(\Omega) . \tag{27}
\end{align*}
$$

In (27) first we choose $h=\left(u_{\mu}-u_{\lambda}\right)^{+} \in W^{1, p}(\Omega)$. Using (26), the fact that $\mu<\lambda$ and that $f \geqslant 0$ and recalling that $u_{\lambda} \in S_{\lambda}$, we conclude that $u_{\mu} \leqslant u_{\lambda}$. Next, in (27) we choose $h=\left(v-u_{\mu}\right)^{+} \in W^{1, p}(\Omega)$. From (26), the fact that $f \geqslant 0$ and Proposition 8 , we infer that $v \leqslant u_{\mu}$. Therefore we have proved that

$$
\begin{equation*}
u_{\mu} \in\left[v, u_{\lambda}\right] . \tag{28}
\end{equation*}
$$

From (26), (27), (28) it follows that

$$
u_{\mu} \in S_{\mu} \subseteq D_{+}(\text {see Proposition } 10)
$$

Let $\rho=\left\|u_{\lambda}\right\|_{\infty}$ and let $\hat{\xi}_{\rho}^{\lambda}>0$ be as postulated by hypothesis $H(f)(v)$. We have

$$
\begin{align*}
& -\Delta_{p} u_{\lambda}(z)-\Delta_{q} u_{\mu}(z)+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] u_{\mu}(z)^{p-1}-u_{\mu}(z)^{-\gamma} \\
= & \mu f\left(z, u_{\mu}(z)\right)+\hat{\xi}_{\rho}^{\lambda} u_{\mu}(z)^{p-1} \\
= & \lambda f\left(z, u_{\mu}(z)\right)+\hat{\xi}_{\rho}^{\lambda} u_{\mu}(z)^{p-1}-(\lambda-\mu) f\left(z, u_{\mu}(z)\right) \\
< & \lambda f\left(z, u_{\mu}(z)\right)+\hat{\xi}_{\rho}^{\lambda} u_{\lambda}(z)^{p-1}(\text { recall that } \lambda>\mu) \\
\leqslant & \lambda f\left(z, u_{\lambda}(z)\right)+\hat{\xi}_{\rho}^{\lambda} u_{\lambda}(z)^{p-1}(\text { see }(28) \text { and hypothesis } H(f)(v)) \\
= & -\Delta_{p} u_{\lambda}(z)-\Delta_{q} u_{\lambda}(z)+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] u_{\lambda}(z)^{p-1}-u_{\lambda}(z)^{-\lambda} \text { for almost all } z \in \Omega \tag{29}
\end{align*}
$$

(recall that $u_{\lambda} \in S_{\lambda}$ ).
We know that

$$
0 \leqslant u_{\mu}^{-\gamma}, u_{\lambda}^{-\gamma} \leqslant v^{-\gamma} \in L^{\infty}(\Omega)
$$

Also, from hypothesis $H(f)(i v)$ and since $u_{\mu} \in D_{+}$, we have

$$
0<c_{8} \leqslant(\lambda-\mu) f\left(z, u_{\mu}(z)\right) \text { for almost all } z \in \Omega
$$

Invoking Proposition 4, from (29) we conclude that

$$
u_{\lambda}-u_{\mu} \in \operatorname{int} \hat{C}_{+} .
$$

The proof is now complete.
Proposition 13. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold, then $\lambda^{*}<+\infty$.
Proof. On account of hypotheses $H(f)(i) \rightarrow(i v)$, we can find $\lambda_{0}>0$ so big that

$$
\begin{equation*}
x^{-\gamma}+\lambda_{0} f(z, x) \geqslant x^{p-1} \text { for almost all } z \in \Omega \text { and all } x \geqslant 0 . \tag{30}
\end{equation*}
$$

Let $\lambda>\lambda_{0}$ and suppose that $\lambda \in \mathcal{L}$. Then we can find $u_{\lambda} \in S_{\lambda} \subseteq D_{+}$(see Proposition 10). Then $m_{\lambda}=\min _{\bar{\Omega}} u_{\lambda}>0$. For $\delta \in(0,1)$ we set $m_{\lambda}^{\delta}=m_{\lambda}+\delta$ and for $\rho=\left\|u_{\lambda}\right\|_{\infty}$ let $\hat{\xi}_{\rho}^{\lambda}>0$ be as postulated by hypothesis $H(f)(v)$. We have

$$
\begin{align*}
& -\Delta_{p} m_{\lambda}^{\delta}-\Delta_{q} m_{\lambda}^{\delta}+\left[\xi(z)+\hat{\xi}_{\rho}\right]\left(m_{\lambda}^{\delta}\right)^{p-1}-\left(m_{\lambda}^{\delta}\right)^{-\gamma} \\
= & {\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] m_{\lambda}^{p-1}-m_{\lambda}^{-\gamma}+\chi(\delta) \text { with } \chi(\delta) \rightarrow 0^{+} \text {as } \delta \rightarrow 0^{+} } \\
< & \xi(z) m_{\lambda}^{p-1}+\left(1+\hat{\xi}_{\rho}^{\lambda}\right) m_{\lambda}^{p-1}-m_{\lambda}^{-\gamma}+\chi(\delta) \\
\leqslant & \lambda_{0} f\left(z, m_{\lambda}\right)+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] m_{\lambda}^{p-1}+\chi(\delta)(\text { see }(30)) \\
\leqslant & \lambda_{0} f\left(z, u_{\lambda}\right)+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] u_{\lambda}^{p-1}+\chi(\delta)(\text { see hypothesis } H(f)(v)) \\
= & \lambda f\left(z, u_{\lambda}\right)+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] u_{\lambda}^{p-1}-\left(\lambda-\lambda_{0}\right) f\left(z, u_{\lambda}\right)+\chi(\delta) \\
= & \lambda f\left(z, u_{\lambda}\right)+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] u_{\lambda}^{p-1} \text { for } \delta \in(0,1) \text { small } \\
& \left(\text { recall that } u_{\lambda} \in D_{+} \text {and see } H(f)(i v)\right) \\
= & -\Delta_{p} u_{\lambda}-\Delta_{q} u_{\lambda}+\left[\xi(z)+\hat{\xi}_{\rho}^{\lambda}\right] u_{\lambda}^{p-1}-u_{\lambda}^{-\gamma} . \tag{31}
\end{align*}
$$

Since $\left(\lambda-\lambda_{0}\right) f\left(z, u_{\lambda}\right)-\chi(\delta) \geqslant c_{9}>0$ for almost all $z \in \Omega$ and for $\delta \in(0,1)$ small (just recall that $u_{\lambda} \in D_{+}$and use hypothesis $\mathrm{H}(\mathrm{f})(\mathrm{iv})$ ), invoking Proposition 4, from (31) we infer that

$$
u_{\lambda}-m_{\lambda}^{\delta} \in \operatorname{int} \hat{C}_{+} \text {for all } \delta \in(0,1) \text { small enough. }
$$

However, this contradicts the definition of $m_{\lambda}$. It follows that $\lambda \notin \mathcal{L}$ and so $\lambda^{*} \leqslant \lambda_{0}<+\infty$.

Therefore we have

$$
\left(0, \lambda^{*}\right) \subseteq \mathcal{L} \subseteq\left(0, \lambda^{*}\right]
$$

Proposition 14. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold and $\lambda \in\left(0, \lambda^{*}\right)$, then problem $\left(P_{\lambda}\right)$ has at least two positive solutions

$$
u_{0}, \hat{u} \in D_{+}, u_{0} \neq \hat{u}
$$

Proof. Let $0<\mu<\lambda<\eta<\lambda^{*}$. According to Proposition 12, we can find $u_{\eta} \in S_{\eta} \subseteq D_{+}, u_{0} \in S_{\lambda} \subseteq D_{+}$ and $u_{\mu} \in S_{\mu} \subseteq D_{+}$such that

$$
\begin{align*}
& u_{\eta}-u_{0} \in \operatorname{int} \hat{C}_{+} \text {and } u_{0}-u_{\mu} \in \operatorname{int} \hat{C}_{+} \\
& \Rightarrow u_{0} \in \operatorname{int}_{C^{1}(\hat{\Omega})}\left[u_{\mu}, u_{\eta}\right] . \tag{32}
\end{align*}
$$

We introduce the following Carathéodory function

$$
\tilde{\tau}_{\lambda}(z, x)= \begin{cases}u_{\mu}(z)^{-\gamma}+\lambda f\left(z, u_{\mu}(z)\right) & \text { if } x<u_{\mu}(z)  \tag{33}\\ x^{-\gamma}+\lambda f(z, x) & \text { if } u_{\mu}(z) \leqslant x \leqslant u_{\eta}(z) \\ u_{\eta}(z)^{-\gamma}+\lambda f\left(z, u_{\eta}(z)\right) & \text { if } u_{\eta}(z)<x\end{cases}
$$

Set $\tilde{T}_{\lambda}(z, x)=\int_{0}^{x} \tilde{\tau}_{\lambda}(z, s) d s$ and consider the $C^{1}$-functional $\tilde{\Psi}_{\lambda}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\tilde{\Psi}_{\lambda}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\lambda} \tilde{T}_{\lambda}(z, u) d z \text { for all } u \in W^{1, p}(\Omega)
$$

Using (33) and the nonlinear regularity theory, we can easily check that

$$
\begin{equation*}
K_{\tilde{\Psi}_{\lambda}} \subseteq\left[u_{\mu}, u_{\eta}\right] \cap D_{+} \tag{34}
\end{equation*}
$$

Also, consider the Carathéodory function

$$
\tau_{\lambda}^{*}(z, x)= \begin{cases}u_{\mu}(z)^{-\gamma}+\lambda f\left(z, u_{\mu}(z)\right) & \text { if } x \leqslant u_{\mu}(z)  \tag{35}\\ x^{-\gamma}+\lambda f(z, x) & \text { if } u_{\mu}(z)<x\end{cases}
$$

We set $T_{\lambda}^{*}(z, x)=\int_{0}^{x} \tau_{\lambda}^{*}(z, s) d s$ and consider the $C^{1}$-functional $\Psi_{\lambda}^{*}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\Psi_{\lambda}^{*}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} T_{\lambda}^{*}(z, u) d z \text { for all } u \in W^{1, p}(\Omega) .
$$

For this functional using (35), we show that

$$
\begin{equation*}
K_{\Psi_{\lambda}^{*}} \subseteq\left[u_{\mu}\right) \cap D_{+} . \tag{36}
\end{equation*}
$$

From (33) and (35) we see that

$$
\begin{equation*}
\left.\tilde{\Psi}_{\lambda}\right|_{\left[u_{\mu}, u_{\eta}\right]}=\left.\Psi_{\lambda}^{*}\right|_{\left[u_{\mu}, u_{\eta}\right]} \text { and }\left.\tilde{\Psi}_{\lambda}^{\prime}\right|_{\left[u_{\mu}, u_{\eta}\right]}=\left.\left(\Psi_{\lambda}^{*}\right)^{\prime}\right|_{\left[u_{\mu}, u_{\lambda}\right]} \tag{37}
\end{equation*}
$$

From (34), (36), (37), it follows that without any loss of generality, we may assume that

$$
\begin{equation*}
K_{\Psi_{\lambda}^{*}} \cap\left[u_{\mu}, u_{\eta}\right]=\left\{u_{0}\right\} . \tag{38}
\end{equation*}
$$

Otherwise it is clear from (35) and (36) that we already have a second positive smooth solution for problem $\left(P_{\lambda}\right)$ and so we are done.

Note that $\tilde{\Psi}_{\lambda}(\cdot)$ is coercive (see (33)). Also, it is sequentially weakly lower semicontinuous. So, we can find $\hat{u}_{0} \in W^{1, p}(\Omega)$ such that

$$
\begin{align*}
& \tilde{\Psi}_{\lambda}\left(\hat{u}_{0}\right)=\inf \left\{\tilde{\Psi}_{\lambda}(u): u \in W^{1, p}(\Omega)\right\}, \\
& \Rightarrow \hat{u}_{0} \in K_{\tilde{\Psi}_{\lambda}} \\
& \Rightarrow \hat{u}_{0} \in K_{\Psi_{\lambda}^{*}} \cap\left[u_{\mu}, u_{\eta}\right](\text { see }(34),(37)),  \tag{39}\\
& \Rightarrow \hat{u}_{0}=u_{0} \in D_{+}(\text {see }(38)), \\
& \Rightarrow u_{0} \text { is a local } C^{1}(\bar{\Omega}) \text {-minimizer of } \Psi_{\lambda}^{*}(\text { see }(32)), \\
& \left.\Rightarrow u_{0} \text { is a local } W^{1, p}(\Omega) \text {-minimizer of } \Psi_{\lambda}^{*} \text { (see Proposition } 5\right) .
\end{align*}
$$

We assume that $K_{\Psi_{\lambda}^{*}}$ is finite. Otherwise on account of (35) and (36) we see that we already have an infinity of positive smooth solutions for problem $\left(P_{\lambda}\right)$ and so we are done. Then (39) implies that we can find $\rho \in(0,1)$ small such that

$$
\begin{equation*}
\Psi_{\lambda}^{*}\left(u_{0}\right)<\inf \left\{\Psi_{\lambda}^{*}(u):\left\|u-u_{0}\right\|=\rho\right\}=m_{\lambda}^{*} \tag{40}
\end{equation*}
$$

(see Papageorgiou, Rădulescu \& Repovš [21, Theorem 5.7.6, p. 367]).
On account of hypothesis $H(f)(i i)$ we have

$$
\begin{equation*}
\Psi_{\lambda}^{*}\left(t \hat{u}_{1}(p)\right) \rightarrow-\infty \text { as } t \rightarrow+\infty \tag{41}
\end{equation*}
$$

Claim 1. $\Psi_{\lambda}^{*}(\cdot)$ satisfies the $C$ - condition.
Let $\left\{u_{n}\right\}_{n \geqslant 1} \subseteq \mathrm{~W}^{1, p}(\Omega)$ be a sequence such that

$$
\begin{align*}
& \left|\Psi_{\lambda}^{*}\left(u_{n}\right)\right| \leqslant c_{10} \text { for some } c_{10}>0 \text { and all } n \in \mathbb{N},  \tag{42}\\
& \left(1+\left\|u_{n}\right\|\right)\left(\Psi_{\lambda}^{*}\right)^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } \mathrm{W}^{1, p}(\Omega)^{*} . \tag{43}
\end{align*}
$$

From (43) we have

$$
\begin{align*}
& \left.\left|\left\langle A_{p}\left(u_{n}\right), h\right\rangle+\left\langle A_{q}\left(u_{n}\right), h\right\rangle+\int_{\Omega} \xi(z)\right| u_{n}\right|^{p-2} u_{n} h d z+\int_{\partial \Omega} \beta(z)\left|u_{n}\right|^{p-2} u_{n} h d \sigma \\
& -\int_{\Omega} \tau_{\lambda}^{*}\left(z, u_{n}\right) h d z \left\lvert\, \leqslant \frac{\epsilon_{n}\|h\|}{1+\left\|u_{n}\right\|}\right. \text { for all } h \in W^{1, p}, \text { with } \epsilon_{n} \rightarrow 0^{+} . \tag{44}
\end{align*}
$$

Choosing $h=-u_{n}^{-} \in W^{1, p}(\Omega)$, we obtain

$$
\begin{align*}
& \gamma_{p}\left(u_{n}^{-}\right)+\left\|D u_{n}^{-}\right\|_{q}^{q} \leqslant c_{11}\left\|u_{n}^{-}\right\| \text {for some } c_{11}>0 \text { and all } n \in \mathbb{N}(\text { see (35)) } \\
\Rightarrow & \left.\left\{u_{n}^{-}\right\}_{n \geqslant 1} \subseteq W^{1, p}(\Omega) \text { is bounded (see (1) and recall that } 1<p\right) . \tag{45}
\end{align*}
$$

Next in (44) we choose $h=u_{n}^{+} \in W^{1, p}(\Omega)$. Then

$$
\begin{align*}
& -\gamma_{p}\left(u_{n}^{+}\right)-\left\|D u_{n}^{+}\right\|_{q}^{q}+\int_{\Omega} \tau_{\lambda}^{*}\left(z, u_{n}\right) u_{n}^{+} d z \leqslant \epsilon_{n} \text { for all } n \in \mathbb{N}, \\
& \Rightarrow-\gamma_{p}\left(u_{n}^{+}\right)-\left\|D u_{n}^{+}\right\|_{q}^{q}+\int_{\left\{u_{n} \leqslant u_{\mu}\right\}}\left[u_{\mu}^{-\gamma}+\lambda f\left(z, u_{\mu}\right)\right] u_{n}^{+} d z  \tag{46}\\
& +\int_{\left\{u_{\mu}<u_{n}\right\}}\left[u_{n}^{-\gamma}+\lambda f\left(z, u_{n}\right)\right] u_{n}^{+} d z \leqslant \epsilon_{n} \text { for all } n \in \mathbb{N}(\text { see (35)) }
\end{align*}
$$

On the other hand from (42) and (45), we have

$$
\begin{gather*}
\gamma_{p}\left(u_{n}^{+}\right)+\frac{p}{q}\left\|D u_{n}^{+}\right\|_{q}^{q}-\int_{\left\{u_{n} \leqslant u_{\mu}\right\}} p\left[u_{\mu}^{-\gamma}+\lambda f\left(z, u_{p}\right)\right] u_{n}^{+} d z \\
-\int_{\left\{u_{\mu}<u_{n}\right\}}\left[\frac{p}{1-\gamma}\left(u_{n}^{1-\gamma}-u_{\mu}^{1-\gamma}\right)+p\left(\lambda F\left(z, u_{n}\right)-\lambda F\left(z, u_{\mu}\right)\right)\right] d z \leqslant \epsilon_{n} \\
\text { for all } n \in \mathbb{N}(\text { see }(35)) . \\
\Rightarrow \gamma_{p}\left(u_{n}^{+}\right)+\frac{p}{q}\left\|D u_{n}^{+}\right\|_{p}^{p}-\int_{\left\{u_{n} \leqslant u_{\mu}\right\}} p\left[u_{\mu}^{-\gamma}+\lambda f\left(z, u_{\mu}\right)\right] u_{n}^{+} d z \\
-\int_{\left\{u_{p}<u_{n}\right\}}\left[\frac{p}{1-\gamma} u_{n}^{1-\gamma}+\lambda p F\left(z, u_{n}\right)\right] d z \leqslant c_{12} \text { for some } c_{12}>0 \text { and all } n \in \mathbb{N} . \tag{47}
\end{gather*}
$$

We add (46) and (47). Since $p>q$, we obtain

$$
\begin{align*}
& \lambda \int_{\left\{u_{\mu}<u_{n}\right\}}\left[f\left(z, u_{n}\right) u_{n}^{+}-p F\left(z, u_{n}\right)\right] d z \leqslant(p-1) \int_{\left\{u_{n} \leqslant u_{\mu}\right\}}\left[u_{\mu}^{-\gamma}+\lambda f\left(z, u_{\mu}\right)\right] u_{n}^{+} d z \\
& +\left(\frac{p}{1-\gamma}-1\right) \int_{\left\{u_{\mu}<u_{n}\right\}} u_{n}^{1-\gamma} d z \\
\Rightarrow & \lambda \int_{\Omega}\left[f\left(z, u_{n}^{+}\right) u_{n}^{+}-p F\left(z, u_{n}^{+}\right)\right] d z \leqslant c_{13}\left[\left\|u_{n}^{+}\right\|_{1}+1\right]  \tag{48}\\
& \text { for some } c_{13}>0, \text { all } n \in \mathbb{N} .
\end{align*}
$$

On account of hypotheses $H(f)(i),(i i i)$ we can find $\hat{\beta}_{1} \in\left(0, \hat{\beta}_{0}\right)$ and $c_{14}>0$ such that

$$
\begin{equation*}
\hat{\beta}_{1} x^{\tau}-c_{14} \leqslant f(z, x)-p F(z, x) \text { for almost all } z \in \Omega \text { and all } x \geqslant 0 . \tag{49}
\end{equation*}
$$

Using (49) in (48), we obtain

$$
\begin{gather*}
\left\|u_{n}^{+}\right\|_{\tau}^{\tau} \leqslant c_{15}\left[\left\|u_{n}^{+}\right\|_{\tau}+1\right] \text { for some } c_{15}>0 \text { and all } n \in \mathbb{N}, \\
\Rightarrow\left\{u_{n}^{+}\right\}_{n \geqslant 1} \leqslant L^{\tau}(\Omega) \text { is bounded. } \tag{50}
\end{gather*}
$$

First assume $N \neq p$. From hypothesis $H(f)(i i i)$ it is clear that we may assume without any loss of generality that $\tau<r<p^{*}$. Let $t \in(0,1)$ be such that

$$
\frac{1}{r}=\frac{1-t}{\tau}+\frac{t}{p^{*}}
$$

Then from the interpolation inequality (see Papageorgiou \& Winkert [22, Proposition 2.3.17, p. 116]), we have

$$
\begin{align*}
& \left\|u_{n}^{+}\right\|_{r} \leqslant\left\|u_{n}^{+}\right\|_{\tau}^{1-t}\left\|u_{n}^{+}\right\|_{p^{*}}^{t}, \\
\Rightarrow & \left\|u_{n}^{+}\right\|_{r}^{r} \leqslant c_{16}\left\|u_{n}^{+}\right\|^{t r} \text { for some } c_{16}>0 \text { and all } n \in \mathbb{N} \text { (see (50)). } \tag{51}
\end{align*}
$$

From hypothesis $H(f)(i)$ we have

$$
\begin{equation*}
f(z, x) x \leqslant c_{17}\left[1+x^{r}\right] \text { for all } z \in \Omega, \text { all } x \geqslant 0 \text { and some } c_{17}>0 . \tag{52}
\end{equation*}
$$

From (44) with $h=u_{n}^{+} \in W^{1, p}(\Omega)$, we obtain

$$
\begin{align*}
& \gamma_{p}\left(u_{n}^{+}\right)+\left\|D u_{n}^{+}\right\|_{q}^{q}-\int_{\Omega} \tau_{\lambda}^{*}\left(z, u_{n}\right) u_{n}^{+} d z \leqslant \epsilon_{n} \text { for all } n \in \mathbb{N}, \\
\Rightarrow & \gamma_{p}\left(u_{n}^{+}\right)+\left\|D u_{n}^{+}\right\|_{q}^{q} \leqslant \int_{\Omega}\left[\left(u_{n}^{+}\right)^{1-\gamma}+f\left(z, u_{n}^{+}\right) u_{n}^{+}\right] d z+c_{18} \\
& \text { for some } c_{18}>0 \text { and all } n \in \mathbb{N}(\text { see (35) ) } \\
\leqslant & c_{19}\left[1+\left\|u_{n}^{+}\right\|_{r}^{r}\right] \text { for some } c_{19}>0 \text { and all } n \in \mathbb{N} \text { (see (52)) } \\
\leqslant & c_{20}\left[1+\left\|u_{n}^{+}\right\|^{t r}\right] \text { for some } c_{20}>0 \text { and all } n \in \mathbb{N} \text { (see (51)). } \tag{53}
\end{align*}
$$

The hypothesis on $\tau($ see $H(f)(i i i))$ implies that $t r<p$. So, from (53) we infer that

$$
\begin{align*}
& \left\{u_{n}^{+}\right\}_{n \geqslant 1} \subseteq W^{1, p}(\Omega) \text { is bounded, } \\
& \Rightarrow\left\{u_{n}\right\}_{n \geqslant 1} \subseteq W^{1, p}(\Omega) \text { is bounded (see (45)). } \tag{54}
\end{align*}
$$

If $N=p$, then $p^{*}=+\infty$ and from the Sobolev embedding theorem, we know that $W^{1, p}(\Omega) \hookrightarrow L^{s}(\Omega)$ for all $1 \leqslant s<\infty$. Then in order for the previous argument to work, we replace $p^{*}=+\infty$ by $s>r>\tau$ and let $t \in(0,1)$ as before such that

$$
\begin{aligned}
& \frac{1}{r}=\frac{1-t}{\tau}+\frac{t}{s} \\
& \Rightarrow t r=\frac{s(r-\tau)}{s-\tau}
\end{aligned}
$$

Note that $\frac{s(r-\tau)}{s-\tau} \rightarrow r-\tau$ as $s \rightarrow+\infty$. But $r-\tau<p$ (see hypothesis $\mathrm{H}(\mathrm{f})(\mathrm{iii})$ ). We choose $s>r$ big so that $t r<p$. Then again we have (54).

Because of (54) and by passing to a subsequence if necessary, we may assume that

$$
\begin{equation*}
u_{n} \xrightarrow{w} u \text { in } W^{1, p}(\Omega) \text { and } u_{n} \rightarrow u \text { in } L^{r}(\Omega) \text { and } L^{p}(\partial \Omega) . \tag{55}
\end{equation*}
$$

In (44) we choose $h=u_{n}-u \in W^{1, p}(\Omega)$, pass to the limit as $n \rightarrow \infty$ and use (55). Then

$$
\begin{aligned}
& \quad \lim _{n \rightarrow \infty}\left[\left\langle A_{p}\left(u_{n}\right), u_{n}-u\right\rangle+\left\langle A_{q}\left(u_{n}\right), u_{n}-u\right\rangle\right]=0, \\
& \Rightarrow \limsup _{n \rightarrow \infty}\left[\left\langle A_{p}\left(u_{n}\right), u_{n}-u\right\rangle+\left\langle A_{q}(u), u_{n}-u\right\rangle\right] \leqslant 0 \\
& \text { (since } A_{q}(\cdot) \text { is monotone) } \\
& \Rightarrow \limsup _{n \rightarrow \infty}\left\langle A_{p}\left(u_{n}\right), u_{n}-u\right\rangle \leqslant 0, \\
& \Rightarrow u_{n} \rightarrow u \text { in } W^{1, p}(\Omega) \text { (see Proposition 1). }
\end{aligned}
$$

Therefore $\Psi_{\lambda}^{*}(\cdot)$ satisfies the C-condition. This proves the Claim.
Then (40), (41) and the Claim permit the use of the mountain pass theorem. So, we can find $\hat{u} \in W^{1, p}(\Omega)$ such that

$$
\hat{u} \in K_{\Psi_{\lambda}^{*}} \leqslant\left[u_{\mu}\right) \cap D_{+}(\operatorname{see}(36)), m_{\lambda}^{*} \leqslant \Psi_{\lambda}^{*}(\hat{u})(\text { see }(40)) .
$$

Therefore $\hat{u} \in D_{+}$is a second positive solution of problem $\left(P_{\lambda}\right)\left(\lambda \in\left(0, \lambda^{*}\right)\right)$ distinct from $u_{0} \in D_{+}$.

Next, we examine what can be said in the critical parameter $\lambda^{*}$.

Proposition 15. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold, then $\lambda^{*} \in \mathcal{L}$.

Proof. Let $\left\{\lambda_{n}\right\}_{n \geqslant 1} \subseteq\left(0, \lambda^{*}\right)$ be such that $\lambda_{n}<\lambda^{*}$. We can find $u_{n} \in S_{\lambda_{n}} \subseteq D_{+}$for all $n \in \mathbb{N}$.
We consider the following Carathéodory function

$$
\mu_{n}(z, x)= \begin{cases}v(z)^{-\gamma}+\lambda_{n} f(z, v(z)) & \text { if } x \leqslant v(z)  \tag{56}\\ x^{-\gamma}+\lambda_{n} f(z, x) & \text { if } v(z)<x\end{cases}
$$

We set $M_{n}(z, x)=\int_{0}^{x} \mu_{n}(z, x) d s$ and consider the $C^{1}$-functional $j_{n}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
j_{n}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} M_{n}(z, u) d z \text { for all } u \in W^{1, p}(\Omega) .
$$

Also, we consider the following truncation of $\mu_{n}(z, \cdot)$

$$
\hat{\mu}_{n}(z, x)= \begin{cases}\mu_{n}(z, x) & \text { if } x \leqslant u_{n+1}(z)  \tag{57}\\ \mu_{n}\left(z, u_{n+1}(z)\right) & \text { if } u_{n+1}(z)<x\end{cases}
$$

(recall that $v \leqslant u_{n+1}$ for all $n \in \mathbb{N}$, see Proposition 11). This is a Carathéodory function. We set $\hat{M}_{n}(z, x)=\int_{0}^{x} \hat{\mu}_{n}(z, s) d s$ and consider the $C^{1}$-functional $\hat{J}_{n}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\hat{J}_{n}(u)=\frac{1}{p} \gamma_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\int_{\Omega} \hat{M}_{n}(z, u) d z \text { for all } u \in W^{1, p}(\Omega) .
$$

From (1), (56) and (57), it is clear that $\hat{J}_{n}(\cdot)$ is coercive. Also, it is sequentially weakly lower semicontinuous. So, we can find $\hat{u}_{n} \in W^{1, p}(\Omega)$ such that

$$
\begin{equation*}
\hat{J}_{n}\left(\hat{u}_{n}\right)=\inf \left\{\hat{J}_{n}(u): u \in W^{1, p}(\Omega)\right\} . \tag{58}
\end{equation*}
$$

Then we have

$$
\begin{align*}
\hat{J}_{n}\left(\hat{u}_{n}\right) \leqslant & \hat{J}_{n}(v) \\
\leqslant & \frac{1}{p} \gamma_{p}(v)+\frac{1}{q}\|D v\|_{q}^{q}-\frac{1}{1-\gamma} \int_{\Omega} v^{1-\gamma} d z \\
& (\text { see }(56),(57) \text { and recall that } f \geqslant 0) \\
\leqslant & \left\langle A_{p}(v), v\right\rangle+\left\langle A_{q}(v), v\right\rangle-\int_{\Omega} v^{1-\gamma} d z=0  \tag{59}\\
& (\text { see Proposition } 8) .
\end{align*}
$$

From (58) we have

$$
\begin{equation*}
\hat{u}_{n} \in K_{\hat{J}_{n}} \subseteq\left[v, u_{n+1}\right] \cap D_{+} \text {for all } n \in \mathbb{N}(\text { see (57)) } \tag{60}
\end{equation*}
$$

Similarly, using (56) we obtain

$$
\begin{equation*}
K_{j_{n}} \subseteq[v) \cap D_{+} . \tag{61}
\end{equation*}
$$

Note that

$$
\left.J_{n}\right|_{\left[v, u_{n+1}\right]}=\left.\hat{J}_{n}\right|_{\left[v, u_{n+1}\right]} \text { and }\left.J_{n}^{\prime}\right|_{\left[v, u_{n+1}\right]}=\left.\hat{J}_{n}^{\prime}\right|_{\left[v, u_{n+1}\right]}(\text { see }(56),(57)) .
$$

Then from (59), (60), (61), we have

$$
\begin{equation*}
J_{n}\left(\hat{u}_{n}\right) \leqslant 0 \text { for all } n \in \mathbb{N} \tag{62}
\end{equation*}
$$

$$
\begin{align*}
& \left\langle A_{p}\left(\hat{u}_{n}\right), h\right\rangle+\left\langle A_{q}\left(\hat{u}_{n}\right), h\right\rangle+\int_{\Omega} \xi(z) \hat{u}_{n}^{p-1} h d z+\int_{\partial \Omega} \beta(z) \hat{u}_{n}^{p-1} h d \sigma=\int_{\Omega} \mu_{n}\left(z, \hat{u}_{n}\right) h d z  \tag{63}\\
& \quad \text { for all } h \in W^{1, p}(\Omega), \text { all } n \in \mathbb{N} .
\end{align*}
$$

Using (62), (63) and reasoning as in the Claim in the proof of Proposition 14, we show that

$$
\left\{\hat{u}_{n}\right\}_{n \geqslant 1} \subseteq W^{1, p}(\Omega) \text { is bounded. }
$$

So, we may assume that

$$
\begin{equation*}
\hat{u}_{n} \xrightarrow{w} \hat{u}_{*} \text { in } W^{1, p}(\Omega) \text { and } \hat{u}_{n} \rightarrow \hat{u}_{*} \text { in } L^{r}(\Omega) \text { and } L^{p}(\partial \Omega) . \tag{64}
\end{equation*}
$$

In (63) we choose $h=\hat{u}_{n}-\hat{u}_{*} \in W^{1, p}(\Omega)$, pass to the limit as $n \rightarrow \infty$ and use (64). Then as before (see the proof of Proposition 14), we obtain

$$
\begin{equation*}
\hat{u}_{n} \rightarrow \hat{u}_{*} \text { in } W^{1, p}(\Omega) \tag{65}
\end{equation*}
$$

In (63) we pass to the limit as $n \rightarrow \infty$ and use (65). Then

$$
\begin{aligned}
& \left\langle A_{p}\left(\hat{u}_{*}\right), h\right\rangle+\left\langle A_{q}\left(\hat{u}_{*}\right), h\right\rangle+\int_{\Omega} \xi(z) \hat{u}_{*}^{p-1} h d z+\int_{\partial \Omega} \beta(z) \hat{u}_{*}^{p-1} h d z \\
& \quad=\int_{\Omega}\left[\hat{u}_{*}^{-\gamma}+\lambda^{*} f\left(z, \hat{u}_{*}\right)\right] h d z \text { for all } h \in W^{1, p}(\Omega)(\text { see (56), (61)) }, \\
& \quad \Rightarrow \hat{u}_{*} \in S_{\lambda^{*}} \subseteq D_{+} \text {and so } \lambda^{*} \in \mathcal{L} .
\end{aligned}
$$

The proof is now complete.
From this proposition it follows that

$$
\mathcal{L}=(0, \lambda *] .
$$

The next bifurcation-type theorem summarizes our findings and provides a complete description of the dependence of the set of positive solutions of problem $\left(P_{\lambda}\right)$ on the parameter $\lambda>0$.

Theorem 16. If hypotheses $H(\xi), H(\beta), H_{0}, H(f)$ hold, then there exists $\lambda^{*}>0$ such that
(a) for all $\lambda \in\left(0, \lambda^{*}\right)$ problem $\left(P_{\lambda}\right)$ has at least two positive solutions

$$
u_{0}, \hat{u} \in D_{+}, u_{0} \neq \hat{u} ;
$$

(b) for $\lambda=\lambda^{*}$ problem $\left(P_{\lambda}\right)$ has at least one positive solution $\hat{u}_{*} \in D_{+}$;
(c) for all $\lambda>\lambda^{*}$ problem $\left(P_{\lambda}\right)$ does not have any positive solutions.

## Acknowledgments

This research was supported by the Slovenian Research Agency grants P1-0292, J1-8131, J1-7025, N1-0064, and N1-0083.

## References

[1] L. Cherfils, Y. Ilyasov, On the stationary solutions of generalized reaction-diffusion equations with $p$ \& $q$ Laplacian, Commun. Pure Appl. Anal. 4 (2005) 9-22.
[2] V.V. Zhikov, Averaging of functionals of the calculus of variations and elasticity theory, Math. USSR-Izv. 29 (1987) 33-66.
[3] L. Gasinski, N.S. Papageorgiou, Nonlinear Analysis, Chapman \& Hall/CRC, Boca Raton, FL, 2006.
[4] J. Giacomoni, J. Schindler, P. Takač, Sobolev versus Hölder local minimizers and existence of multiple solutions for a singular quasilinear equation, Ann. Sc. Norm Super. Pisa Ser. V 6 (2007) 117-158.
[5] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Positive solutions for nonlinear parametric singular Dirichlet problems, Bull. Math. Sci. 9 (2) (2019) 1950011, 21.
[6] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Pairs of positive solutions for resonant singular equations with the p-Laplacian, Electron. J. Differ. Equ. (2017) 13, Paper No. 249.
[7] N.S. Papageorgiou, G. Smyrlis, A bifurcation-type theorem for singular nonlinear elliptic equations, Methods Appl. Anal. 22 (2015) 147-170.
[8] N.S. Papageorgiou, P. Winkert, Singular p-Laplacian equations with superlinear perturbation, J. Differential Equations 266 (2019) 1462-1487.
[9] K. Perera, Z. Zhang, Multiple positive solutions of singular $p$-Laplacian problems by variational methods, Bound. Value Probl. 2005 (2005) 3.
[10] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Double-phase problems with reaction of arbitrary growth, Z. Angew. Math. Phys. 69 (4) (2018) 21, Art. 108.
[11] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Double-phase problems and a discontinuity property of the spectrum, Proc. Amer. Math. Soc. 147 (7) (2019) 2899-2910.
[12] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Existence and multiplicity of solutions for double-phase Robin problems, Bull. Lond. Math. Soc. 52 (2020) 546-560.
[13] A. Bahrouni, V.D. Rădulescu, D.D. Repovš, Double-phase transonic flow problems with variable growth: nonlinear patterns and stationary waves, Nonlinearity 32 (7) (2019) 2481-2495.
[14] M. Ghergu, V.D. Rădulescu, Singular Elliptic Problems. Bifurcation and Asymptotic Analysis, in: Oxford Lecture Series in Mathematics and its Applications, vol. 37, The Clarendon Press, Oxford University Press, Oxford, 2008.
[15] N.S. Papageorgiou, V.D. Rădulescu, Positive solutions for nonlinear nonhomogeneous parametric Robin problems, Forum Math. 30 (2018) 553-580.
[16] N.S. Papageorgiou, V.D. Rădulescu, Nonlinear nonhomogeneous Robin problems with superlinear reaction, Adv. Nonlinear Stud. 16 (2016) 737-764.
[17] G. Lieberman, The natural generalization of the natural conditions of Ladyzhenskaya and Uraltseva for elliptic equations, Comm. Partial Differential Equations 16 (1991) 311-361.
[18] L. Gasinski, N.S. Papageorgiou, Exercises in Analysis. Part 2: Nonlinear Analysis, Springer, Cham, 2016.
[19] P. Pucci, J. Serrin, The Maximum Principle, Birkhäuser, Basel, 2007.
[20] N.S. Papageorgiou, V.D. Rădulescu, Multiple solutions with precise sign for nonlinear parametre Robin problems, J. Differential Equations 254 (2014) 393-430.
[21] N.S. Papageorgiou, V.D. Rădulescu, D.D. Repovš, Modern Nonlinear Analysis - Theory and Methods, Springer, Cham, 2019.
[22] N.S. Papageorgiou, P. Winkert, Applied Nonlinear Functional Analysis, de Gruyter, Berlin, 2018.


[^0]:    * Corresponding author at: Department of Mathematics, University of Craiova, 200585 Craiova, Romania.

    E-mail addresses: npapg@math.ntua.gr (N.S. Papageorgiou), radulescu@inf.ucv.ro (V.D. Rădulescu), dusan.repovs@guest.arnes.si (D.D. Repovš).

