# An infinity of nodal solutions for superlinear Robin problems with an indefinite and unbounded potential 

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#### Abstract

We consider a semilinear Robin problem driven by the Laplacian plus an indefinite and unbounded potential using a suitable version of the symmetric mountain pass theorem, we show that the problem has an infinity of nodal solutions whose energy level diverges to $+\infty$. © 2017 Elsevier Masson SAS. All rights reserved.


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## 1. Introduction

Let $\Omega \subseteq \mathbb{R}^{N}$ be a bounded domain with a $C^{2}$-boundary $\partial \Omega$. In this paper, we study the following semilinear Robin problem

$$
\left\{\begin{array}{ll}
-\Delta u(z)+\xi(z) u(z)=f(z, u(z)) & \text { in } \Omega  \tag{1}\\
\frac{\partial u}{\partial n}+\beta(z) u=0 & \text { on } \partial \Omega
\end{array}\right\}
$$

In this problem, the potential function $\xi \in L^{s}(\Omega)$ with $s>N$ and is indefinite (that is, $\xi(\cdot)$ is sign changing). The reaction term $f(z, x)$ is a Carathéodory function (that is, for all $x \in \mathbb{R} z \rightarrow f(z, x)$ is measurable and for almost all $z \in \Omega x \rightarrow f(z, x)$ is continuous), which is superlinear in the $x \in \mathbb{R}$ variable, but without satisfying the usual in such cases Ambrosetti-Rabinowitz condition (AR-condition for short). In addition, for almost all $z \in \Omega f(z, \cdot)$ satisfies a one-sided Lipschitz condition and it is odd. In the boundary condition $\frac{\partial u}{\partial n}$ denotes the usual normal derivative defined by extension of the map

$$
C^{1}(\bar{\Omega}) \ni u \rightarrow \frac{\partial u}{\partial n}=(D u, n)_{\mathbb{R}^{N}}
$$

with $n(\cdot)$ being the outward unit normal on $\partial \Omega$. The boundary coefficient $\beta \in W^{1, \infty}(\partial \Omega)$ with $\beta(z) \geqslant 0$ for all $z \in \partial \Omega$. When $\beta \equiv 0$, then we recover the Neumann problem.

We are looking for the existence of multiple nodal (that is, sign changing) solutions for problem (1). Using a version of the symmetric mountain pass theorem due to Qian and Li [13, Theorem 4.2], we show the existence of a sequence of distinct nodal solutions with energies diverging to $+\infty$.

In the past, an infinity of nodal solutions for superlinear Dirichlet problems with $\xi \equiv 0$, we proved by Qian and $\mathrm{Li}[13$, Theorem 5.4] using AR-condition and with more restrictive conditions on the reaction term $f$. Subsequently, Qian [12, Theorem 1.1] produced an infinity of nodal solutions for a superlinear Neumann problem with $\xi \equiv a \in(0,+\infty)$. So, in Qian [12] the differential operator (right-hand side of the equation), is coercive and this simplifies the arguments considerably. Qian [12] did not use the AR-condition and instead employed a condition which was first introduced by Jeanjean [4]. This condition is global in nature and for this reason not entirely satisfactory. For Robin problems, there is the work of Qian and $\mathrm{Li}[14]$, who assume that $\xi \equiv 0$ and $f \in C(\bar{\Omega} \times \mathbb{R})$ satisfies the Jeanjean condition. They produce an infinity of distinct solutions, but they do not show that these solutions are nodal (see [14, Theorem 1.3]).

Problems with indefinite linear part (that is, the potential function $\xi(\cdot)$ is indefinite), were investigated by Zhang and Liu [18], Qin, Tang and Zhang [15], Zhang, Tang and Zhang [19]. All the aforementioned works deal with Dirichlet problems and use a nonquadraticity condition analogous to the one employed by Costa and Magalhaes [2]. They produce infinitely many nontrivial solutions, but the not show that they are nodal. Multiple nodal solutions for problems with indefinite linear part, were produced by Papageorgiou and Papalini [7] (Dirichlet problems), Papageorgiou and Rădulescu [8]
(Neumann problems) and Papageorgiou and Rădulescu [10] (Robin problems). None of the above works produces infinitely many nodal solutions. Finally we mention the very recent paper of Papageorgiou and Rădulescu [11], who produce a sequence of nodal solutions for nonlinear Robin problems but under different conditions and using different tools.

## 2. Mathematical background

Let $X$ be a Banach space and $X^{*}$ its topological dual. By $\langle\cdot, \cdot\rangle$ we denote the duality brackets for the pair $\left(X^{*}, X\right)$. Given $\varphi \in C^{1}(X, \mathbb{R})$, we say that $\varphi$ satisfies the "Cerami condition" (the "C-condition" for short), if the following property holds
"Every sequence $\left\{u_{n}\right\}_{n \geqslant 1} \subseteq X$ such that $\left\{\varphi\left(u_{n}\right)\right\}_{n \geqslant 1} \subseteq \mathbb{R}$ is bounded and

$$
\left(1+\left\|u_{n}\right\|\right) \varphi^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } X^{*}
$$

admits a strongly convergent subsequence".
This is a compactness-type condition on $\varphi$, more general than the usual Palais-Smale condition. Nevertheless, it leads to the same deformation theorem from which one can derive the minimax theory of the critical values of $\varphi$.

The following spaces will be important in our analysis:

- The Sobolev space $H^{1}(\Omega)$;
- The Banach space $C^{1}(\bar{\Omega})$;
- The "boundary" Lebesgue spaces $L^{q}(\partial \Omega), 1 \leqslant q \leqslant \infty$.

The Sobolev space $H^{1}(\Omega)$ is a Hilbert space with inner product

$$
(u, h)_{H^{1}(\Omega)}=\int_{\Omega} u h d z+\int_{\Omega}(D u, D h)_{\mathbb{R}^{N}} d z \text { for all } u, h \in H^{1}(\Omega)
$$

and corresponding norm

$$
\|u\|=\left[\|u\|_{2}^{2}+\|D u\|_{2}^{2}\right]^{1 / 2} \text { for all } u \in H^{1}(\Omega)
$$

The Banach space $C^{1}(\bar{\Omega})$ is an order Banach space with positive (order) cone given by

$$
C_{+}=\left\{u \in C^{1}(\bar{\Omega}: u(z) \geqslant 0 \text { for all } z \in \bar{\Omega})\right\}
$$

This cone has a nonempty interior containing

$$
D_{+}=\left\{u \in C_{+}: u(z)>0 \text { for all } z \in \bar{\Omega}\right\}
$$

On $\partial \Omega$ we consider the ( $N-1$ )-dimensional Hausdorff (surface) measure $\sigma(\cdot)$. Using this measure, we can define in the usual way the Lebesgue spaces $L^{q}(\partial \Omega) 1 \leqslant q \leqslant \infty$. According to the theory of Sobolev spaces, there exists a unique continuous linear map $\gamma_{0}: H^{1}(\Omega) \rightarrow L^{2}(\partial \Omega)$ known as the "trace map", which satisfies

$$
\gamma_{0}(u)=\left.u\right|_{\partial \Omega} \text { for all } u \in H^{1}(\Omega) \cap C(\bar{\Omega})
$$

So, the trace map assigns "boundary values" to all Sobolev functions. This map is compact into $L^{q}(\partial \Omega)$ for all $q \in\left[1, \frac{2 N-2}{N-2}\right)$ if $N \geqslant 3$ and into $L^{q}(\partial \Omega)$ for all $q \geqslant 1$ if $N=1,2$. In addition we have

$$
\operatorname{ker} \gamma_{0}=H_{0}^{1}(\Omega) \text { and } i m \gamma_{0}=H^{\frac{1}{2}, 2}(\partial \Omega)
$$

In the sequel, for the sake of notational economy, we drop the use of the map $\gamma_{0}$. All restrictions of Sobolev functions on $\partial \Omega$ are understood in the sense of traces.

Next we consider the following linear eigenvalue problem:

$$
\left\{\begin{array}{ll}
-\Delta u(z)+\xi(z) u(z)=\hat{\lambda} u(z) & \text { in } \Omega  \tag{2}\\
\frac{\partial u}{\partial n}+\beta(z) u=0 & \text { on } \partial \Omega
\end{array}\right\}
$$

We impose the following conditions on the data of this eigenvalue problem

- $\xi \in L^{\frac{N}{2}}(\Omega)$ if $N \geqslant 3, \xi \in L^{q}(\Omega)$ with $q>1$ if $N=2$ and $\xi \in L^{1}(\Omega)$ if $N=1$.
- $\beta \in W^{1, \infty}(\partial \Omega)$ with $\beta(z) \geqslant 0$ for all $z \in \partial \Omega$.

Let $\vartheta: H^{1}(\Omega) \rightarrow \mathbb{R}$ be the $C^{1}$-functions defined by

$$
\vartheta(u)=\|D u\|_{2}^{2}+\int_{\Omega} \xi(z) u^{2} d z+\int_{\partial \Omega} \beta(z) u^{2} d \sigma \text { for all } u \in H^{1}(\Omega)
$$

We know (see [10]) that there exists $\mu>0$ such that

$$
\begin{equation*}
\vartheta(u)+\mu\|u\|_{2}^{2} \geqslant c_{0}\|u\|^{2} \text { for all } u \in H^{1}(\Omega), \text { some } c_{0}>0 \tag{3}
\end{equation*}
$$

Using (3) and the spectral theorem for compact self-adjoint operators on a Hilbert space, we define the spectrum of (2) consisting of a sequence $\left\{\hat{\lambda}_{k}\right\}_{k \geqslant 1} \subseteq \mathbb{R}$ such that $\hat{\lambda}_{k} \rightarrow+\infty$. By $E\left(\hat{\lambda}_{k}\right) k \in \mathbb{N}$ we denote the eigenspace corresponding to the eigenvalue $\hat{\lambda}_{k}$. We know that each $E\left(\hat{\lambda}_{k}\right)$ is finite dimensional and we have the following orthogonal direct sum decomposition

$$
H^{1}(\Omega)=\overline{{\underset{k}{ }}_{\oplus} E\left(\hat{\lambda}_{k}\right)}
$$

We know that

- $\hat{\lambda}_{1}$ is simple (that is, $\left.\operatorname{dim} E\left(\hat{\lambda}_{1}\right)=1\right)$.
- $\hat{\lambda}_{1}=\inf \left[\frac{\vartheta(u)}{\|u\|_{2}^{2}}: u \in H^{1}(\Omega), u \neq 0\right]$.
- $\hat{\lambda}_{m}=\inf \left[\frac{\vartheta(u)}{\|u\|_{2}^{2}}: u \in \overline{\bigoplus_{k}^{\oplus} E\left(\hat{\lambda}_{k}\right)}, u \neq 0\right]$
$=\sup \left[\frac{\vartheta(u)}{\|u\|_{2}^{2}}: u \in \underset{\mathrm{k}=1}{\oplus} E\left(\hat{\lambda}_{k}\right), u \neq 0\right] m \geqslant 2$.
The infimum in (4) is realized on $E\left(\hat{\lambda}_{1}\right)$. Both the infimum and supremum in (5) are realized on $E\left(\hat{\lambda}_{m}\right)$. Evidently the elements of $E\left(\hat{\lambda}_{1}\right)$ do not change sign, while the elements of $E\left(\hat{\lambda}_{m}\right) m \geqslant 2$ are nodal (that is sign changing).

In what follows $A: H^{1}(\Omega) \rightarrow H^{1}(\Omega)^{*}$ is the bounded linear operator defined by

$$
\langle A(u), h\rangle=\int_{\Omega}(D u, D h)_{\mathbb{R}^{N}} d z \text { for all } u, h \in H^{1}(\Omega)
$$

Also, by $|\cdot|_{N}$ we denote the Lebesgue measure on $\mathbb{R}^{N}$ and $2^{*}= \begin{cases}\frac{2 N}{N-2} & \text { if } N \geqslant 3, \\ +\infty & \text { if } N=1,2 .\end{cases}$

## 3. A sequence of nodal solutions

Our hypotheses on the data of (1) are the following:
$H(\xi): \xi \in L^{s}(\Omega), s>N$ and $\xi^{+} \in L^{\infty}(\Omega)$.
$H(\beta): \beta \in W^{1, \infty}(\partial \Omega)$ with $\beta(z) \geqslant 0$ for all $z \in \partial \Omega$.
$H(f): f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a Carathéodory function such that for almost all $z \in \Omega f(z, \cdot)$ is odd and
(i) $|f(z, x)| \leqslant a(z)\left(1+|x|^{r-1}\right)$ for almost all $z \in \Omega$, all $x \in \mathbb{R}$, with $a \in L^{\infty}(\Omega), r \in$ $\left(2,2^{*}\right)$;
(ii) if $F(z, x)=\int_{0}^{x} f(z, s) d s$, then $\lim _{x \rightarrow \pm \infty} \frac{F(z, x)}{x^{2}}=+\infty$ uniformly for almost all $z \in \Omega$;
(iii) if $e(z, x)=f(z, x) x-2 F(z, x)$, then there exists $d \in L^{1}(\Omega)$ such that

$$
e(z, x) \leqslant e(z, y)+d(z) \text { for almost all } z \in \Omega, \text { all } 0 \leqslant x \leqslant y \text { or } y \leqslant x \leqslant 0
$$

(iv) there exists $\hat{\eta}>0$ such that for almost all $z \in \Omega$ the function

$$
x \rightarrow f(z, x)+\hat{\eta} x
$$

is increasing on $\mathbb{R}$;
(v) there exist $\hat{c}_{0}, \hat{c}_{1}>0$ such that

$$
-\hat{c}_{0} \leqslant \liminf _{x \rightarrow 0} \frac{f(z, x)}{x} \leqslant \limsup _{x \rightarrow 0} \frac{f(z, x)}{x} \leqslant \hat{c}_{1} \text { uniformly for almost all } z \in \Omega
$$

Remark 1. Hypothesis $H(f)(i i)$ implies that the primitive $F(z, \cdot)$ is superquadratic near $+\infty$. Hypotheses $H(f)(i i),(i i i)$ imply that

$$
\lim _{x \rightarrow \pm \infty} \frac{f(z, x)}{x}=+\infty \text { uniformly for almost all } z \in \Omega
$$

So, the reaction term $f(z, \cdot)$ is superlinear. However, this superlinearity is not expressed via the classical AR-condition, which says that there exist $q>2$ and $M>0$ such that

$$
\begin{align*}
& 0<q F(z, x) \leqslant f(z, x) x \text { for almost all } z \in \Omega, \text { all }|x| \geqslant M  \tag{6a}\\
& 0<e s s \inf _{\Omega} F(\cdot, \pm M) \tag{6b}
\end{align*}
$$

(see Ambrosetti and Rabinowitz [1] and Mugnai [6]). Integrating (6a) and using (6b), we obtain the weaker condition

$$
c_{1}|x|^{q} \leqslant F(z, x) \text { for almost all } z \in \Omega, \text { all }|x| \geqslant M, \text { some } c_{1}>0
$$

This means that under the AR-condition $f(z, \cdot)$ has at least ( $q-1$ )-polynomial growth near $\pm \infty$. The Jeanjean condition used in some works mentioned in the Introduction, says that there exist $\eta \geqslant 1$ and $s \in[0,1]$ such that

$$
e(z, s x) \leqslant \eta e(z, x) \text { for almost all } z \in \Omega, \text { all } x \in \mathbb{R}
$$

We mention the global nature of this condition. This is a feature which we would like to avoid. Here instead of the AR-condition and the Jeanjean condition we employ a quasimonotonicity condition on $e(z, \cdot)$ (see hypothesis $H(f)(i i i)$ ). This condition is a slightly more general version of a condition used by Li and Yang [5]. It is satisfied if there exists $M \geqslant 0$ such that

$$
e(z, \cdot) \text { is nondecreasing on } x \geqslant M \text { and nonincreasing on } x \leqslant-M \text {. }
$$

In turn, this is implied by the following condition

$$
\begin{aligned}
& x \rightarrow \frac{f(z, x)}{x} \text { is nondecreasing on } x \geqslant M, \\
& x \rightarrow \frac{f(z, x)}{x} \text { is nonincreasing on } x \leqslant-M
\end{aligned}
$$

We stress the local character of the last two conditions. Hypothesis $H(f)(i v)$ is a one-sided Lipschitz condition. Finally hypothesis $H(f)(v)$ implies that for almost all $z \in \Omega, f(z, \cdot)$ is linear near zero.

Let $\varphi: H^{1}(\Omega) \rightarrow \mathbb{R}$ be the energy (Euler) functional for problem (1) defined by

$$
\varphi(u)=\frac{1}{2} \vartheta(u)-\int_{\Omega} F(z, u) d z \text { for all } u \in H^{1}(\Omega)
$$

Evidently $\varphi \in C^{1}\left(H^{1}(\Omega)\right)$.
Proposition 1. If hypotheses $H(\xi), H(\beta), H(f)$ hold, then the functional $\varphi$ satisfies the C-condition.

Proof. We consider a sequence $\left\{u_{n}\right\}_{n \geqslant 1} \subseteq H^{1}(\Omega)$ such that

$$
\begin{align*}
& \left|\varphi\left(u_{n}\right)\right| \leqslant M_{1} \text { for some } M_{1}>0, \text { all } n \in \mathbb{N}  \tag{7}\\
& (1+\|u\|) \varphi^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } H^{1}(\Omega)^{*} \text { as } n \rightarrow \infty \tag{8}
\end{align*}
$$

From (8) we have

$$
\begin{equation*}
\left|\left\langle A\left(u_{n}\right), h\right\rangle+\int_{\Omega} \xi(z) u_{n} h d z+\int_{\partial \Omega} \beta(z) u_{n} h d \sigma-\int_{\Omega} f\left(z, u_{n}\right) h d z\right| \leqslant \frac{\epsilon_{n}\|h\|}{1+\left\|u_{n}\right\|} \tag{9}
\end{equation*}
$$

for all $h \in H^{1}(\Omega)$, with $\epsilon_{n} \rightarrow 0^{+}$
In (9) we choose $h=u_{n} \in H^{1}(\Omega)$. Then

$$
\begin{equation*}
\vartheta\left(u_{n}\right)+\int_{\Omega} f\left(z, u_{n}\right) u_{n} d z \leqslant \epsilon_{n} \text { for all } n \in \mathbb{N} \tag{10}
\end{equation*}
$$

From (7), we have

$$
\begin{equation*}
-\vartheta\left(u_{n}\right)-\int_{\Omega} 2 F\left(z, u_{n}\right) d z \leqslant 2 M_{1} \text { for all } n \in \mathbb{N} \tag{11}
\end{equation*}
$$

Adding (10) and (11), we obtain

$$
\begin{equation*}
\int_{\Omega} e\left(z, u_{n}\right) d z \leqslant M_{2} \text { for some } M_{2}>0, \text { all } n \in \mathbb{N} \text {. } \tag{12}
\end{equation*}
$$

Claim 1. $\left\{u_{n}\right\}_{n \geqslant 1} \subseteq H^{1}(\Omega)$ is bounded.

We argue by contradiction. So, suppose that the Claim is not true. By passing to a suitable subsequence if necessary, we may assume that

$$
\begin{equation*}
\left\|u_{n}\right\| \rightarrow \infty \tag{13}
\end{equation*}
$$

Let $y_{n}=\frac{u_{n}}{\left\|u_{n}\right\|} n \in \mathbb{N}$. Then $\left\|y_{n}\right\|=1$ for all $n \in \mathbb{N}$ and so we may assume that

$$
\begin{equation*}
y_{n} \xrightarrow{x} y \text { in } H^{1}(\Omega) \text { and } y_{n} \rightarrow y \text { in } L^{r}(\Omega) \text { and } L^{2}(\partial \Omega) \tag{14}
\end{equation*}
$$

(note that we can always assume $r \geqslant \frac{2 s}{s-1}$, see hypothesis $H(f)(i)$ ).
First assume that $y \neq 0$ and let $\Omega_{0}=\{z \in \Omega: y(z) \neq 0\}$. We have $\left|\Omega_{0}\right|_{N}>0$ and

$$
\left|u_{n}(z)\right| \rightarrow+\infty \text { for almost all } z \in \Omega_{0}
$$

Then hypothesis $H(f)(i i)$ implies that

$$
\begin{equation*}
\frac{F\left(z, u_{n}(z)\right)}{\left\|u_{n}\right\|^{2}}=\frac{F\left(z, u_{n}(z)\right)}{u_{n}(z)^{2}} y_{n}(z)^{2} \rightarrow+\infty \text { for almost all } z \in \Omega_{0} \text { as } n \rightarrow \infty \tag{15}
\end{equation*}
$$

Using (15) and Fatou's lemma (it can be used on account of hypothesis $H(f)(i i i)$ ), we have

$$
\begin{equation*}
\frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega_{0}} F\left(z, u_{n}\right) d z \rightarrow+\infty \text { as } n \rightarrow \infty \tag{16}
\end{equation*}
$$

Hypothesis $h(f)(i i)$ implies that we can find $M_{3}>0$ such that

$$
\begin{equation*}
F(z, x) \geqslant 0 \text { for almost all } z \in \Omega, \text { all }|x| \geqslant M_{3} . \tag{17}
\end{equation*}
$$

We have

$$
\begin{align*}
\frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega} F\left(z, u_{n}\right) d z= & \frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega_{0}} F\left(z, u_{n}\right) d z+\frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega_{0}^{c} \cap\left\{\left|u_{n}\right| \geqslant M_{3}\right\}} F\left(z, u_{n}\right) d z+ \\
& \frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega_{0}^{c} \cap\left\{\left|u_{n}\right|<M_{3}\right\}} F\left(z, u_{n}\right) d z \\
\geqslant & \frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega_{0}} F\left(z, u_{n}\right) d z-\frac{c_{2}}{\left\|u_{n}\right\|^{2}} \text { for some } c_{2}>0, \text { all } n \in \mathbb{N} \\
& (\operatorname{see}(17 \text { and use hypothesis }) H(f)(i)) \\
\Rightarrow & \lim _{n \rightarrow \infty} \frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega} F\left(z, u_{n}\right) d z=+\infty(\text { see }(16)) \tag{18}
\end{align*}
$$

On the other hand from (7) we have

$$
\begin{align*}
& \frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega} F\left(z, u_{n}\right) d z \leqslant \frac{M_{1}}{\left\|u_{n}\right\|^{2}}+\frac{1}{2} \vartheta\left(y_{n}\right) \text { for all } n \in \mathbb{N} \\
\Rightarrow & \frac{1}{\left\|u_{n}\right\|^{2}} \int_{\Omega} F\left(z, u_{n}\right) d z \leqslant M_{4} \text { for some } M_{4}>0, \text { all } n \in \mathbb{N} \tag{19}
\end{align*}
$$

(see hypotheses $H(\xi), H(\beta)$ and recall that $\left\|y_{n}\right\|=1, n \in \mathbb{N}$ )
Comparing (16) and (19), we reach a contradiction.
Next suppose that $y=0$. Given $\tau>0$, let

$$
v_{n}=(2 \tau)^{1 / 2} y_{n} \text { for all } n \in \mathbb{N}
$$

We have

$$
\left.v_{n} \rightarrow 0 \text { in } L^{r}(\Omega) \text { and in } L^{2}(\partial \Omega) \text { (see (14) and recall that } \mathrm{y}=0\right) .
$$

It follows that

$$
\begin{equation*}
\int_{\Omega} F\left(z, v_{n}\right) d z \rightarrow 0 \text { as } n \rightarrow \infty \tag{20}
\end{equation*}
$$

From (13) we see that we can find $n_{0} \in \mathbb{N}$ such that

$$
\begin{equation*}
0<(2 \tau)^{1 / 2} \frac{1}{\left\|u_{n}\right\|} \leqslant 1 \text { for all } n \geqslant n_{0} \tag{21}
\end{equation*}
$$

Choose $t_{n} \in[0,1]$ such that

$$
\begin{equation*}
\varphi\left(t_{n} u_{n}\right)=\max [\varphi(t u): 0 \leqslant t \leqslant 1] \text { for all } n \in \mathbb{N} . \tag{22}
\end{equation*}
$$

Taking into account (21), we have

$$
\begin{align*}
\varphi\left(t_{n} u_{n}\right) & \geqslant \varphi\left(v_{n}\right) \\
& =\tau \vartheta\left(y_{n}\right)-\int_{\Omega} F\left(z, u_{n}\right) d z \\
& \geqslant \tau\left[c_{0}-\mu\left\|y_{n}\right\|_{2}^{2}\right]-\int_{\Omega} F\left(z, v_{n}\right) d z \text { for all } n \geqslant n_{0}(\text { see }(3)) . \tag{23}
\end{align*}
$$

Passing to the limit as $n \rightarrow \infty$ in (23) and using (14) and (20) and recalling that $y=0$, we obtain

$$
\liminf _{n \rightarrow \infty} \varphi\left(t_{n} u_{n}\right) \geqslant \tau c_{0}
$$

But $\tau>0$ is arbitrary. So, it follows that

$$
\begin{equation*}
\varphi\left(t_{n} u_{n}\right) \rightarrow+\infty \text { as } n \rightarrow \infty \tag{24}
\end{equation*}
$$

We have

$$
\varphi(0)=0 \text { and } \varphi\left(u_{n}\right) \leqslant M_{1} \text { for all } n \in \mathbb{N}(\text { see }(7)) .
$$

The from (24) we infer that we can find $n_{2} \in \mathbb{N}$ such that

$$
\begin{equation*}
t_{n} \in(0,1) \text { for all } n \geqslant n_{2} \tag{25}
\end{equation*}
$$

From (22) and (25) it follows that

$$
\begin{align*}
& \left.\frac{d}{d t} \varphi\left(t u_{n}\right)\right|_{t=t_{n}}=0 \text { for all } n \geqslant n_{2} \\
\Rightarrow & \left\langle\varphi^{\prime}\left(t_{n} u_{n}\right), t_{n} u_{n}\right\rangle=0 \text { for all } n \geqslant n_{2}(\text { bu the chain rule }), \\
\Rightarrow & \vartheta\left(t_{n} u_{n}\right)=\int_{\Omega} f\left(z, t_{n} u_{n}\right)\left(t_{n} u_{n}\right) d z \text { for all } n \geqslant n_{2} \tag{26}
\end{align*}
$$

Hypothesis $H(f)(i i i)$ and (25) imply that

$$
\begin{aligned}
& \int_{\Omega} e\left(z, t u_{n}\right) d z \\
= & \int_{\Omega}\left[e\left(z, t u_{n}^{+}\right)+e\left(z,-t_{n} u_{n}^{-}\right)\right] d z \text { (note that } e(z, 0)=0 \text { for almost all } z \in \Omega \text { ) } \\
\leqslant & \int_{\Omega}\left[e\left(z, u_{n}^{+}\right)+e\left(z,-u_{n}^{-}\right)\right] d z+\|d\|_{1} \text { (see hypothesis } H(f)(i i i) \text { ) } \\
= & \int_{\Omega} e\left(z, u_{n}\right) d z+\|d\|_{1} \text { for all } n \geqslant n_{2} .
\end{aligned}
$$

Using this inequality in (26), we obtain

$$
\begin{align*}
2 \varphi\left(t_{n} u_{n}\right) & \leqslant \int_{\Omega} e\left(z, u_{n}\right) d z+\|d\|_{1} \\
& \leqslant M_{2}+\|d\|_{1}=M_{5} \text { for all } n \geqslant n_{2}(\text { see }(12)) \tag{27}
\end{align*}
$$

Comparing (24) and (27), we have a contradiction.
This proves the Claim.

On account of the Claim, we may assume that

$$
\begin{equation*}
u_{n} \xrightarrow{w} u \text { and } u_{n} \rightarrow u \text { in } L^{r}(\Omega) \text { and in } L^{2}(\partial \Omega) . \tag{28}
\end{equation*}
$$

In (9) we choose $h=-u \in H^{1}(\Omega)$, pass to the limit as $n \rightarrow \infty$ and use (28) (recall $r \geqslant \frac{2 s}{s-1}$ ). We obtain

$$
\begin{aligned}
& \lim _{n \rightarrow \infty}\left\langle A\left(u_{n}\right), u_{n}-u\right\rangle=0 \\
\Rightarrow & \|D u\|_{2} \rightarrow\|D u\|_{2}, \\
\Rightarrow & u_{n} \rightarrow u \text { in } H^{1}(\Omega) \text { (by the Kadec-Klee property for Hilbert spaces, see (28)), } \\
\Rightarrow & \varphi \text { satisfies the C-condition. }
\end{aligned}
$$

For every $m \in \mathbb{N}$, we define

$$
Y_{m}=\underset{\mathrm{k}=1}{\oplus} E\left(\hat{\lambda}_{k}\right) \text { and } V_{m}=\overline{\oplus_{k \geqslant m}^{\oplus} E\left(\hat{\lambda}_{k}\right)}
$$

Let

$$
\begin{equation*}
\beta_{m}=\sup \left[\|u\|_{r}: u \in V_{m},\|u\|=1\right] \tag{29}
\end{equation*}
$$

As in the proof of Lemma 3.8 of Willem [17, p. 60], we show that

$$
\begin{equation*}
\beta_{m} \rightarrow 0^{+} \text {as } m \rightarrow+\infty \tag{30}
\end{equation*}
$$

Proposition 2. If hypotheses $H(\xi), H(\beta), H(f)$ hold, then we can find $\left\{l_{m}\right\}_{m \in \mathbb{N}} \subseteq$ $(0,+\infty)$ such that

$$
\gamma_{m}=\inf \left[\varphi(u): u \in V_{m},\|u\|=l_{m}\right] \rightarrow+\infty \text { as } m \rightarrow \infty .
$$

Proof. Let $u \in V_{m}$. We have

$$
\begin{align*}
\varphi(u)= & \frac{1}{2} \vartheta(u)-\int_{\Omega} F(z, u) d z \\
= & \frac{1}{2} \vartheta(u)+\frac{\mu}{2}\|u\|_{2}^{2}-\frac{\mu}{2}\|u\|_{2}^{2}-\int_{\Omega} F(z, u) d z(\text { with } \mu>0 \text { as in (3)) } \\
\geqslant & \frac{c_{0}}{2}\|u\|^{2}-\frac{\mu}{2}\|u\|_{2}^{2}-c_{3}\|u\|_{r}^{r}-c_{3} \text { for some } c_{3}>0  \tag{31}\\
& (\text { see }(3 \text { and hypothesis H(f)(i)) }
\end{align*}
$$

Recall that $r>2$ (see hypothesis $H(f)(i))$. So, we can find $c_{4}>0$ such that

$$
\begin{equation*}
\|u\|_{2} \leqslant c_{4}\|u\|_{r} \text { for all } u \in H^{1}(\Omega) \tag{32}
\end{equation*}
$$

Using (32) and (31) we obtain

$$
\varphi(u) \geqslant \frac{c_{0}}{2}\|u\|^{2}-c_{5}\left(\|u\|_{r}^{r}+\|u\|_{r}^{2}\right)-c_{3} \text { for some } c_{5}>0, \text { all } u \in V_{m}
$$

Suppose that $\|u\| \geqslant 1$. Then using once more the fact that $r>2$, we obtain

$$
\begin{equation*}
\varphi(u) \geqslant \frac{c_{0}}{2}\|u\|^{2}-c_{6}\|u\|_{r}^{r}-c_{3} \text { with } c_{6}=2 c_{5}>0, \text { all } u \in V_{m},\|u\| \geqslant 1 \tag{33}
\end{equation*}
$$

From (29) we have

$$
\beta_{m}\|u\| \geqslant\|u\|_{r} \text { for all } u \in V_{m}
$$

Using this inequality in (33), we obtain

$$
\begin{equation*}
\varphi(u) \geqslant \frac{c_{0}}{2}\|u\|^{2}-c_{6} \beta_{m}^{r}\|u\|^{r}-c_{3} \text { for all } u \in V_{m},\|u\| \geqslant 1 . \tag{34}
\end{equation*}
$$

Let $l_{m}=\left(\frac{c_{6}}{c_{0}} r \beta_{m}^{r}\right)^{\frac{1}{2-r}}$, we have

$$
l_{m} \rightarrow+\infty \text { as } m \rightarrow+\infty(\text { see }(30) \text { and recall that } r>2)
$$

Hence we may assume that $l_{m} \geqslant 1$ for all $m \in \mathbb{N}$. Then from (34) we see that for all $u \in V_{m}$ with $\|u\|=l_{m}$, we have

$$
\begin{aligned}
\varphi(u) & \geqslant \frac{c_{0}}{2}\left(\frac{c_{6}}{c_{0}} r \beta_{m}^{r}\right)^{\frac{2}{2-r}}-c_{6} \beta_{m}^{r}\left(\frac{c_{6}}{c_{0}} r \beta_{m}^{r}\right)^{\frac{r}{2-r}} \\
& =\left[\frac{c_{0}}{2}-c_{6} \beta_{m}^{r} \frac{c_{0}}{c_{6} r \beta_{m}^{r}}\right]\left(\frac{c_{6}}{c_{0}} r \beta_{m}^{r}\right)^{\frac{2}{2-r}} \\
& =c_{0}\left[\frac{1}{2}-\frac{1}{r}\right]\left(\frac{c_{6}}{c_{0}} r \beta_{m}^{r}\right)^{\frac{2}{2-r}}, \\
& \Rightarrow l_{m} \rightarrow+\infty \text { as } m \rightarrow \infty(\text { see }(30) \text { and recall that } r>2)
\end{aligned}
$$

Proposition 3. If hypotheses $H(\xi), H(\beta), H(f)$ hold, then we can find $\left\{\rho_{m}\right\}_{m \in \mathbb{N}} \subseteq(0, \infty)$, $\rho_{0}>l_{m}>0$ for all $m \in \mathbb{N}$ such that

$$
\Im_{m}=\sup \left[\varphi(u): u \in Y_{m},\|u\|=\rho_{m}\right] \leqslant 0 \text { for all } m \in \mathbb{N} .
$$

Proof. Let $u \in Y_{m}$. We have

$$
\begin{align*}
\varphi(u) & =\frac{1}{2} \vartheta(u)-\int_{\Omega} F(z, u) d z \\
& \leqslant \frac{1}{2}\|D u\|_{2}^{2}+\frac{1}{2} \int_{\Omega} \xi^{+}(z) u^{2} d z+\frac{1}{2} \int_{\partial \Omega} \beta(z) u^{2} d \sigma-\int_{\Omega} F(z, u) d z \tag{35}
\end{align*}
$$

Hypotheses $H(f)(i),(i i)$ imply that given any $\eta>0$, we can find $c_{7}=c_{7}(\eta)>0$ such that

$$
F(z, x) \geqslant \eta x^{2}-c_{7} \text { for almost all } z \in \Omega, \text { all } x \in \mathbb{R}
$$

Using this unilateral growth estimate and hypothesis $H(\xi)$ in (35), we obtain

$$
\varphi(u) \leqslant c_{8}\|u\|^{2}-\eta\|u\|_{2}^{2}+c_{7}|\Omega|_{N} \text { for some } c_{8}>0, \text { all } u \in Y_{m}
$$

But $Y_{m}$ is finite dimensional. So, all norms are equivalent. Hence we can find $c_{9}>0$ such that

$$
\begin{align*}
\varphi(u) & \leqslant c_{8}\|u\|^{2}-\eta c_{9}\|u\|^{2}+c_{7}|\Omega|_{N} \\
& =\left[c_{8}-\eta c_{9}\right]\|u\|^{2}+c_{7}|\Omega|_{N} \text { for all } u \in Y_{m} \tag{36}
\end{align*}
$$

Recall that $\eta>0$ is arbitrary. So, we choose $\eta>\frac{c_{8}}{c_{9}}$. Then from (36) it is clear that we can find $\rho_{m}>l_{m} m \in \mathbb{N}$ such that

$$
\begin{aligned}
& \varphi(u) \leqslant 0 \text { for all } u \in Y_{m},\|u\|=\rho_{m}, \\
\Rightarrow & \Im_{m} \leqslant 0 \text { for all } m \in \mathbb{N} .
\end{aligned}
$$

Proposition 4. If hypotheses $H(\xi), H(\beta), H(f)$ hold and $u \in H^{1}(\Omega)$ is a solution of (1), then $u \in C^{1, \alpha}(\bar{\Omega})$ with $\alpha=1-\frac{N}{s}>0$ (see hypothesis $H(\xi)$ ).

Proof. Hypotheses $H(f)(i),(v)$ imply that

$$
\begin{equation*}
|f(z, x)| \leqslant c_{10}\left(|x|+|x|^{r-1}\right) \text { for almost all } z \in \Omega, \text { all } x \in \mathbb{R}, \text { some } c_{10}>0 \tag{37}
\end{equation*}
$$

By hypothesis we have

$$
-\Delta u(z)+\xi(z) u(z)=f(z, u(z)) \text { for almost all } z \in \Omega
$$

(see also Papageorgiou and Rădulescu [9]),

$$
\Rightarrow-\Delta u(z)=\left[\frac{f(z, u(z))}{u(z)}-\xi(z)\right] u(z) \text { for almost all } z \in \Omega
$$

Note that $f(z, 0)=0$ for almost all $z \in \Omega$ (see hypothesis $H(f)(v))$ and let

$$
\hat{a}(z)= \begin{cases}\frac{f(z, u(z))}{u(z)} & \text { if } u(z) \neq 0 \\ 0 & \text { if } u(z)=0\end{cases}
$$

Then

$$
\begin{aligned}
|\hat{a}(z)| & \leqslant \frac{|f(z, u(z))|}{|u(z)|}+|\xi(z)| \\
& \leqslant c_{10}\left(1+|u(z)|^{r-2}\right)+|\xi(z)| \text { for almost all } z \in \Omega(\text { see }(37))
\end{aligned}
$$

Note that $|u(\cdot)|^{r-2} \in L^{\frac{2^{*}}{r-2}}(\Omega)$ (recall that $u \in H^{1}(\Omega)$ and use the Sobolev embedding theorem) and observe that $\frac{2^{*}}{r-2}>\frac{N}{2}$ (recall that $r<2^{*}$ ). Therefore

$$
\hat{a} \in L^{q}(\Omega) \text { with } q>\frac{N}{2}(\text { see hypothesis } H(\xi))
$$

Then Lemma 5.1 of Wang [16] implies that

$$
u \in L^{\infty}(\Omega)
$$

Using this fact and hypotheses $H(f)(i)$ and $H(\xi)$, we have

$$
f(\cdot, u(\cdot))-\xi(\cdot) u(\cdot) \in L^{s}(\Omega) \text { with } s>N
$$

So, the Calderon-Zygmund estimates (see Wang [16, Lemma 5.2]), we have

$$
u \in W^{2, s}(\Omega)
$$

The Sobolev embedding theorem implies that

$$
u \in C^{1, \alpha}(\bar{\Omega}) \text { with } \alpha=1-\frac{N}{s}>0
$$

Proposition 5. If hypotheses $H(\xi), H(\beta), H(f)$ hold and $u, v \in H^{1}(\Omega)$ are distinct solutions of (1) such that $v \leqslant u$, then $u-v \in D_{+}$.

Proof. From Proposition 4, we know that

$$
u, v \in C^{1}(\bar{\Omega})
$$

Let $\hat{\eta}>0$ be as in hypothesis $H(f)(i v)$. Then

$$
\begin{aligned}
& -\Delta v(z)+(\xi(z)+\hat{\eta}) v(z) \\
= & f(z, v(z))+\hat{\eta} v(z) \\
\leqslant & f(z, u(z))+\hat{\eta} u(z) \text { (see hypothesis } H(f)(i v) \text { and recall that } v \leqslant u \text { ) } \\
= & -\Delta u(z)+(\xi(z)+\hat{\eta}) u(z) \text { for almost all } z \in \Omega,
\end{aligned}
$$

$$
\begin{aligned}
\Rightarrow & \Delta(u-v)(z) \leqslant\left[\left\|\xi^{+}\right\|_{\infty}+\hat{\eta}\right](u-v)(z) \text { for almost all } z \in \Omega \\
& (\text { see hypothesis } H(\xi)), \\
\Rightarrow & u-v \in D_{+}
\end{aligned}
$$

(by the strong maximum principle, see Gasinski and Papageorgiou [3, p. 738]).
Corollary 6. If hypotheses $H(\xi), H(\beta), H(f)$ hold and $u \in H^{1}(\Omega), u \neq 0, u \geqslant 0$ is a solution of (1), then $u \in D_{+}$.

From Proposition 5, Corollary 6 and Proposition 5.4 of Qian and Li [13] (see also the proof of Theorem 2 in Papageorgiou and Papalini [7]), we obtain the following result.

Proposition 7. If hypotheses $H(\xi), H(\beta), H(f)$ hold, then $C_{+}$is an admissible invariant set of $\varphi$.

We make a final observation before formulating our multiplicity theorem.
Proposition 8. If hypotheses $H(\xi), H(\beta), H(f)$ and $l_{m}>0 m \in \mathbb{N}$ is as in Proposition 2, then $V_{m} \cap \partial B_{l_{m}} \cap C_{+}=\emptyset$ for all $m \geqslant 2$ (hence $\partial B_{l_{m}}=\left\{u \in H^{1}(\Omega):\|u\|=l_{m}\right\}$ ).

Proof. Let $\hat{u}_{1}$ be the positive, $L^{2}$-normalized (that is, $\left\|\hat{u}_{1}\right\|_{2}=1$ ) eigenfunction corresponding to $\hat{\lambda}_{1}$. The regularity theory (see [16]) and the strong maximum principle (see [3]), imply that $\hat{u}_{1} \in D_{+}$.

For $u \in C_{+} \backslash\{0\}$ we have

$$
\int_{\Omega} u \hat{u}_{1} d z>0
$$

On the other hand for every $u \in V_{m}$ with $m \geqslant 2$, we have

$$
\int_{\Omega} u \hat{u}_{1} d z=0\left(\text { since } V_{m}^{1} \supseteq E\left(\hat{\lambda}_{1}\right)\right) .
$$

Therefore $V_{m} \cap \partial B_{l_{m}} \cap C_{+}=\emptyset$.
All these auxiliary results permit the use of Theorem 4.2 of Qian and Li [13] (the symmetric mountain pass theorem). So, we have the following multiplicity theorem.

Theorem 9. If hypotheses $H(\xi), H(\beta), H(f)$ hold, the problem (1) admits a sequence $\left\{u_{n}\right\}_{n \geqslant 1} \subseteq C^{1}(\bar{\Omega})$ of distinct nodal solutions such that $\varphi\left(u_{n}\right) \rightarrow+\infty$.

## Conflict of interest statement

There is no conflict of interest.

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