# Nonlinear eigenvalue problems for the $(p, q)$-Laplacian 

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A B S T R A C T
We consider a parametric ( $p, q$ )-equations with sign-changing reaction and Robin boundary condition. We show that for all values of the parameter $\lambda$ bigger than a certain value which we determine precisely, the problem has at least three nontrivial solutions all with sign information and ordered. For the particular case of $(p, 2)$-equations we produce a second nodal solution, for a total of four nontrivial solutions. Under symmetry conditions, we show the existence of infinitely many nodal solutions. The same results are also valid for the Dirichlet problem.
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## 1. Introduction

We study the following parametric $(p, q)$-equation with Robin boundary condition

$$
\left\{\begin{array}{l}
-\Delta_{p} u(z)-\Delta_{q} u(z)+\xi(z)|u(z)|^{p-2} u(z)=\lambda f(z, u(z)) \text { in } \Omega \\
\frac{\partial u}{\partial n_{p q}}+\beta(z)|u|^{p-2} u=0 \text { on } \partial \Omega, 1<q<p, \lambda>0
\end{array}\right.
$$

In this problem, $\Omega \subseteq \mathbb{R}^{N}$ is a bounded domain with a $C^{2}$-boundary $\partial \Omega$. For $1<r<$ $+\infty$ we denote by $\Delta_{r}$ the $r$-Laplace differential operator defined by

$$
\Delta_{r} u=\operatorname{div}\left(|D u|^{r-2} D u\right) \text { for all } u \in W^{1, r}(\Omega)
$$

In problem $\left(P_{\lambda}\right)$, in the left-hand side we have the sum of two such operators. So, the differential operator in $\left(P_{\lambda}\right)$ is not homogeneous. There is also a potential term $\xi(z)|u|^{p-2} u$ with $\xi \geq 0$. The reaction (right-hand side of $\left(P_{\lambda}\right)$ ) is parametric with $\lambda>0$ being the parameter and $f(z, x)$ is a Carathéodory function (that is, for all $x \in \mathbb{R}$, $z \mapsto f(z, x)$ is measurable and for a.a. $z \in \Omega, x \mapsto f(z, x)$ is continuous).

In contrast to most similar works in the literature, $f(z, \cdot)$ can be sign-changing. In the boundary condition $\frac{\partial u}{\partial n_{p q}}$ denotes the conormal derivative corresponding to the differential operator $u \mapsto-\Delta_{p} u-\Delta_{q} u$ (the ( $p, q$ )-Laplacian). We interpret this directional derivative using the nonlinear Green's identity (see [21, p. 35]). We know that if $u \in C^{1}(\bar{\Omega})$, then

$$
\frac{\partial u}{\partial n_{p q}}=\left(|D u|^{p-2}+|D u|^{q-2}\right) \frac{\partial u}{\partial n}
$$

with $n(\cdot)$ being the outward unit normal.
So, problem $\left(P_{\lambda}\right)$ is a kind of a nonlinear eigenvalue problem for the Robin $(p, q)$ Laplacian plus a potential term. We want to find those parameter values for which problem $\left(P_{\lambda}\right)$ has solutions and provide sign information for all of them. Our work here complements those of Gasiński \& Papageorgiou [8], Li \& Yang [12], Papageorgiou \& Rădulescu [15], Papageorgiou, Rădulescu \& Repovš [20]. In these works the reaction $f(z, \cdot)$ is $(p-1)$-superlinear as $x \rightarrow \pm \infty$ and they focus only on the existence of positive solutions. In addition, Gasiński \& Papageorgiou [8] and Li \& Yang [12] deal with equations driven by the Dirichlet $p$-Laplacian only. Related to our work, is also the last part in the paper of Gasiński \& Papageorgiou [6], who consider equations driven by the Dirichlet $p$-Laplacian and a sign-changing reaction satisfying more restrictive conditions. They prove a bifurcation type result describing the changes in the set of positive solutions as the parameter $\lambda$ moves on $\stackrel{\circ}{\mathbb{R}}_{+}=(0,+\infty)$. We also mention the recent work of Papageorgiou \& Zhang [23], on positive solutions of resonant $(p, q)$-equations.

Under minimal conditions of $f(z, \cdot)$, we show that for all $\lambda>0$ problem $\left(P_{\lambda}\right)$ has constant sign smooth solutions. If the parameter $\lambda>0$ is restricted to be big enough
(we determine the lower bound of the values of $\lambda$ using the data of the problem), then we can show the existence of a smooth nodal solution. Under a symmetry condition of $f(z, \cdot)$, we show the existence of a sequence of nodal solutions. When $q=2$ (case of $(p, 2)$-equations), then we are able to show the existence of a second nodal solution. Our tools are variational from critical point theory, combined with truncation and comparison techniques and critical groups.

The double-phase problem $\left(P_{\lambda}\right)$ is motivated by numerous models arising in mathematical physics. For instance, we can refer to the following Born-Infeld equation [1] that appears in electromagnetism:

$$
-\operatorname{div}\left(\frac{\nabla u}{\left(1-2|\nabla u|^{2}\right)^{1 / 2}}\right)=h(u) \text { in } \Omega .
$$

Indeed, by the Taylor formula, we have

$$
(1-x)^{-1 / 2}=1+\frac{x}{2}+\frac{3}{2 \cdot 2^{2}} x^{2}+\frac{5!!}{3!\cdot 2^{3}} x^{3}+\cdots+\frac{(2 n-3)!!}{(n-1)!2^{n-1}} x^{n-1}+\cdots \text { for }|x|<1
$$

Taking $x=2|\nabla u|^{2}$ and adopting the first order approximation, we obtain problem $\left(P_{\lambda}\right)$ for $p=4$ and $q=2$. Furthermore, the $n$-th order approximation problem is driven by the multi-phase differential operator

$$
-\Delta u-\Delta_{4} u-\frac{3}{2} \Delta_{6} u-\cdots-\frac{(2 n-3)!!}{(n-1)!} \Delta_{2 n} u
$$

Our work here appears to be the first one on nonlinear eigenvalue problems driven by the $(p, q)$-Laplacian with Robin boundary condition. Our hypotheses on the reaction are minimal, very general, and they include the case of sign-changing forcing term. Moreover, we provide sign information for all solutions produced.

## 2. Background material and hypotheses

The main spaces in the analysis of problem $\left(P_{\lambda}\right)$, are the Sobolev space $W^{1, p}(\Omega)$, the Banach space $C^{1}(\bar{\Omega})$ and the "boundary" Lebesgue spaces $L^{s}(\partial \Omega), 1 \leq s \leq+\infty$.

By $\|\cdot\|$ we denote the norm of the Sobolev space $W^{1, p}(\Omega)$. We have

$$
\|u\|=\left(\|u\|_{p}^{p}+\|D u\|_{p}^{p}\right)^{1 / p}
$$

The space $C^{1}(\bar{\Omega})$ is an ordered Banach space with positive (order) cone

$$
C_{+}=\left\{u \in C^{1}(\bar{\Omega}): u(z) \geq 0 \text { for all } z \in \bar{\Omega}\right\}
$$

This cone has a nonempty interior given by

$$
\operatorname{int} C_{+}=\left\{u \in C_{+}: u(z)>0 \text { for all } z \in \bar{\Omega}\right\} .
$$

On $\partial \Omega$ we consider the ( $N-1$ )-dimensional Hausdorff (surface) measure $\sigma(\cdot)$. Using this measure, we can define in the usual way the "boundary" Lebesgue spaces $L^{s}(\partial \Omega)$ $(1 \leq s \leq+\infty)$. From the theory of Sobolev spaces, we know that there exists a unique continuous linear map $\gamma_{0}: W^{1, p}(\Omega) \mapsto L^{p}(\partial \Omega)$, known as the "trace map". We know that if $u \in W^{1, p}(\Omega) \cap C(\bar{\Omega})$, then $\gamma_{0}(u)=\left.u\right|_{\partial \Omega}$. So, the trace map extends to all Sobolev functions the notion of boundary values. We know that $\gamma_{0}(\cdot)$ is compact into $L^{s}(\partial \Omega)$, for all $1 \leq s<\frac{(N-1) p}{N-p}$ if $p<N$ and into $L^{s}(\partial \Omega)$ for all $1 \leq s<+\infty$ if $N \leq p$. Moreover, we have

$$
\operatorname{im} \gamma_{0}=W^{\frac{1}{p^{\prime}}, p}(\partial \Omega) \quad\left(\frac{1}{p^{\prime}}+\frac{1}{p}=1\right) \text { and } \operatorname{ker} \gamma_{0}=W_{0}^{1, p}(\Omega)
$$

In what follows for the sake of notational economy, we drop the use of the trace map $\gamma_{0}(\cdot)$. All restrictions of the Sobolev functions on $\partial \Omega$ are understood in the sense of traces.

If we consider the $q$-Laplace differential operator with Neumann boundary condition, then $\hat{\lambda}_{1}(q)=0$ is the first eigenvalue with corresponding eigenspace $\mathbb{R}$ (the constant functions). The positive $L^{q}(\Omega)$-normalized principal eigenfunction is $\hat{u}_{1}(q)=\frac{1}{|\Omega|_{N}}$ with $|\cdot|_{N}$ being the Lebesgue measure on $\mathbb{R}^{N}$. By $\hat{\lambda}_{2}(q)$ we denote the first positive eigenvalue. We have the following variational characterization of $\hat{\lambda}_{2}(q)$ (see Cuesta, de Figueiredo \& Gossez [3] (Dirichlet problems), Mugnai \& Papageorgiou [14], Neumann problems with indefinite potential). We set $\partial B_{1}^{L^{q}}=\left\{u \in L^{q}(\Omega):\|u\|_{q}=1\right\}, M=W^{1, p}(\Omega) \cap \partial B_{1}^{L^{q}}$ and $\Gamma=\left\{\gamma \in C([-1,1], M): \gamma(-1)=-\hat{u}_{1}(q), \gamma(1)=\hat{u}_{1}(q)\right\}$.

Proposition 2.1. $\hat{\lambda}_{2}(q)=\inf _{\gamma \in \Gamma} \max _{-1 \leq t \leq 1}\|D \gamma(t)\|_{q}^{q}$.
If $q=2$, then we know that $-\Delta$ with Neumann boundary condition has a sequence of distinct eigenvalues $\left\{\hat{\lambda}_{m}(2)\right\}_{m \in \mathbb{N}}$ which satisfy $\hat{\lambda}_{m}(2) \rightarrow+\infty$ as $m \rightarrow \infty$ and describe completely the spectrum of the operator. Of course $\hat{\lambda}_{1}(2)=0$. There is a corresponding sequence $\left\{\hat{u}_{n}\right\}_{n \in \mathbb{N}} \subseteq H_{0}^{1}(\Omega)$ of eigenfunctions which are an orthonormal basis for $H^{1}(\Omega)$. By $E\left(\hat{\lambda}_{m}(2)\right)$ we denote the eigenspace corresponding to the eigenvalue $\hat{\lambda}_{m}(2)$. These items have the following properties:
(a) $E\left(\hat{\lambda}_{m}(2)\right)(m \in \mathbb{N})$ is finite dimensional and $E\left(\hat{\lambda}_{m}(2)\right) \subseteq C^{1}(\bar{\Omega})$ (see Brezis [2])
(b) Each eigenspace has the so-called "Unique Continuation Property" (UCP for short), which means that if $u \in E\left(\hat{\lambda}_{m}(2)\right)$ vanishes on $A$ with $|A|_{N}>0$, then $u \equiv 0$.
(c) $H^{1}(\Omega)=\underset{m \geq 1}{\oplus} E\left(\hat{\lambda}_{m}(2)\right)$ (orthogonal direct sum decomposition) and

$$
\begin{equation*}
\hat{\lambda}_{1}(2)=\inf \left\{\frac{\|D u\|_{2}^{2}}{\|u\|_{2}^{2}}: u \in H^{1}(\Omega), u \neq 0\right\}=0 \tag{1}
\end{equation*}
$$

$$
\begin{equation*}
\hat{\lambda}_{n}(2)=\sup \left\{\frac{\|D u\|_{2}^{2}}{\|u\|_{2}^{2}}: u \in \bar{H}_{n}, u \neq 0\right\}=\inf \left\{\frac{\|D u\|_{2}^{2}}{\|u\|_{2}^{2}}: u \in \hat{H}_{n}, u \neq 0\right\} \tag{2}
\end{equation*}
$$

where $\bar{H}_{n}=\underset{m=1}{\oplus} E\left(\hat{\lambda}_{m}(2)\right), \hat{H}_{n}=\overline{\underset{m \geq n}{\oplus} E\left(\hat{\lambda}_{m}(2)\right)}, n \in \mathbb{N}$ (see Papageorgiou \& Rădulescu [18]).

The infimum in (1) is clearly attained on $\mathbb{R}$ (the eigenspace of $\hat{\lambda}_{1}(2)=0$ ), while both the supremum and infimum in (2) are realized on $E\left(\hat{\lambda}_{m}(2)\right)$. All eigenvalues $\hat{\lambda}_{m}(2)$ ( $m \geq 2$ ) have nodal eigenfunctions.

Using the orthogonality of the eigenspaces, the UCP and (1), (2) we have the following Lemma (see Papageorgiou \& Winkert [22]).

## Lemma 2.2.

(a) If $m \in \mathbb{N}, \vartheta \in L^{\infty}(\Omega), \vartheta(z) \geq \hat{\lambda}_{m}(2)$ for a.a. $z \in \Omega$ and the inequality is strict on $a$ set $A$ with $|A|_{N}>0$, then

$$
C_{1}\|u\|_{H^{1}(\Omega)}^{2} \leq \int_{\Omega} \vartheta(z) u^{2} d z-\|D u\|_{2}^{2}
$$

for some $C_{1}>0$, all $u \in \bar{H}_{m}$.
(b) If $m \in \mathbb{N}, \vartheta \in L^{\infty}(\Omega), \vartheta \leq \hat{\lambda}_{m}(2)$ for a.a. $z \in \Omega$ and the inequality is strict on a set $A$ with $|A|_{N}>0$, then

$$
C_{2}\|u\|_{H^{1}(\Omega)}^{2} \leq\|D u\|_{2}^{2}-\int_{\Omega} \vartheta(z) u^{2} d z
$$

for some $C_{2}>0$ all $u \in \hat{H}_{m}$.

Our hypotheses on the potential function $\xi(\cdot)$ and the boundary coefficient $\beta(\cdot)$ are the following:
$\mathbf{H}_{\mathbf{0}}: \xi \in L^{\infty}(\Omega), \beta \in C^{0, \alpha}(\partial \Omega)$ with $0<\alpha<1, \xi(z) \geq 0$ for a.a. $z \in \Omega, \beta(z) \geq 0$ for all $z \in \partial \Omega$ and $\xi \neq 0$ or $\beta \neq 0$.

If $k_{p}: W^{1, p}(\Omega) \mapsto \mathbb{R}$ is the $C^{1}$-functional defined by

$$
k_{p}(u)=\|D u\|_{p}^{p}+\int_{\Omega} \xi(z)|u|^{p} d z+\int_{\partial \Omega} \beta(z)|u|^{p} d \sigma
$$

then using the hypotheses $\mathbf{H}_{\mathbf{0}}$, Lemma 4.11 of Mugnai \& Papageorgiou [14] and Proposition 2.4 of Gasiński \& Papageorgiou [6], we have

$$
\begin{equation*}
C_{0}\|u\|^{p} \leq k_{p}(u) \text { for some } C_{0}>0, \text { all } W^{1, p}(\Omega) \tag{3}
\end{equation*}
$$

In particular, the nonlinear eigenvalue problem:

$$
\left\{\begin{array}{l}
-\Delta_{p} u+\xi(z)|u|^{p-2} u=\tilde{\lambda}|u|^{p-2} u \text { in } \Omega \\
\frac{\partial u}{\partial n_{p}}+\beta(z)|u|^{p-2} u=0 \text { on } \partial \Omega
\end{array}\right.
$$

has a positive smallest eigenvalue $\tilde{\lambda}_{1}(p)$ which is isolated, simple and

$$
\begin{gathered}
\tilde{\lambda}_{1}(p)=\inf \left\{\frac{k_{p}(u)}{\|u\|_{p}^{p}}: u \in W^{1, p}(\Omega), u \neq 0\right\}>0 \\
\text { (see Papageorgiou \& Rădulescu [18]). }
\end{gathered}
$$

If $u, v: \Omega \mapsto \mathbb{R}$ are measurable functions such that $v(z) \leq u(z)$ for a.a. $z \in \Omega$, then we introduce the following order interval in $W^{1, p}(\Omega)$

$$
[v, u]=\left\{h \in W^{1, p}(\Omega): v(z) \leq h(z) \leq u(z) \text { for a.a. } z \in \Omega\right\}
$$

By int $C_{C^{1}(\bar{\Omega})}[v, u]$ we denote the interior of $[v, u] \cap C^{1}(\bar{\Omega})$ in $C^{1}(\bar{\Omega})$.
If $u \in W^{1, p}(\Omega)$, we set $u^{ \pm}=\max \{ \pm u, 0\}$. We know that $u^{ \pm} \in W^{1, p}(\Omega), u=u^{+}-u^{-}$, $|u|=u^{+}+u^{-}$.

Given $r \in(1,+\infty)$, we denote by $A_{r}: W^{1, r}(\Omega) \rightarrow W^{1, r}(\Omega)^{*}$ the nonlinear operator defined by

$$
\left\langle A_{r}(u), h\right\rangle=\int_{\Omega}|D u|^{r-2}(D u, D h)_{\mathbb{R}^{N}} d z \text { for all } u, h \in W^{1, r}(\Omega)
$$

This operator is continuous, monotone (hence maximal monotone) and of type $(S)_{+}$, that is,

$$
\begin{aligned}
& \text { "if } u_{n} \xrightarrow{w} u \text { in } W^{1, r}(\Omega) \text { and } \limsup _{n \rightarrow \infty}\left\langle A_{r}\left(u_{n}\right), u_{n}-u\right\rangle \leq 0, \\
& \text { then } u_{n} \rightarrow u \text { in } W^{1, r}(\Omega) . "
\end{aligned}
$$

This property is a consequence of the Kadec-Klee property (also known as the RadonRiesz property) of uniformly convex spaces. This property says that if $X$ is uniformly convex and $x_{n} \xrightarrow{w} x,\left\|x_{n}\right\| \rightarrow\|x\|$, then $x_{n} \rightarrow x$.

Let $X$ be a Banach space, $\varphi \in C^{1}(X)$ and $c \in \mathbb{R}$. We define

$$
K_{\varphi}=\left\{u \in X: \varphi^{\prime}(u)=0\right\}, K_{\varphi}^{c}=\left\{u \in K_{\varphi}: \varphi(u)=c\right\}, \varphi^{c}=\{u \in X: \varphi(u) \leq c\}
$$

We say that $\varphi(\cdot)$ satisfies the "PS-condition", if:

$$
\begin{aligned}
& \left\{\varphi\left(u_{n}\right)\right\}_{n \in \mathbb{N}} \subseteq \mathbb{R} \text { is bounded } \\
& \text { and } \varphi^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } X^{*} \text { as } n \rightarrow \infty
\end{aligned}
$$

admits a strongly convergent subsequence".

Finally let $Y_{2} \subseteq Y_{1} \subseteq X$. By $H_{k}\left(Y_{1}, Y_{2}\right)\left(k \in \mathbb{N}_{0}\right)$, we denote the $k^{t h}$-relative singular homology, group with integer coefficients. If $u \in K_{\varphi}$ is isolated, then the critical groups of $\varphi$ at $u$ are defined by

$$
C_{k}(\varphi, u)=H_{k}\left(\varphi^{c} \cap U, \varphi^{c} \cap U \backslash\{u\}\right)
$$

with $c=\varphi(u), k \in \mathbb{N}_{0}$ and $U$ an open neighborhood of $u$ such that $\varphi^{c} \cap K_{\varphi} \cap U=\{u\}$. The excision property of singular homology implies that this definition is independent of the isolating neighborhood $U$. Suppose that $\varphi$ satisfies the PS-condition and that $K_{\varphi}$ is finite. Then the critical groups of $\varphi$ at infinity are defined by

$$
C_{k}(\varphi, \infty)=H_{k}\left(X, \varphi^{c}\right)
$$

for all $k \in \mathbb{N}_{0}$ with $c<\inf \varphi\left(K_{\varphi}\right)$. The Second Deformation Theorem (see [21, p. 386]), implies that this definition is independent of the choice of $c<\inf \varphi\left(K_{\varphi}\right)$. We define

$$
\begin{aligned}
& M(t, u)=\sum_{k \geq 0} \operatorname{rank} C_{k}(\varphi, u) t^{k} \text { for all } t \in \mathbb{R}, \text { all } u \in K_{\varphi} \\
& P(t, \infty)=\sum_{k \geq 0} \operatorname{rank} C_{k}(\varphi, \infty) t^{k} \text { for all } t \in \mathbb{R}
\end{aligned}
$$

The Morse relation says that

$$
\sum_{u \in K_{\varphi}} M(t, u)=P(t, \infty)+(1+t) Q(t) \text { for all } t \in \mathbb{R}
$$

with $Q(t)=\sum_{k \geq 0} \beta_{k} t^{k}$ a formal series in $t$ with nonnegative integer coefficients.
Next we introduce the hypotheses on $f(z, x)$ :
$\mathbf{H}_{\mathbf{1}}: f: \Omega \times \mathbb{R} \mapsto \mathbb{R}$ is a Carathéodory function such that $f(z, 0)=0$ for a.a. $z \in \Omega$ and
(i) $|f(z, x)| \leq a(z)\left(1+|x|^{r-1}\right)$ for a.a. $z \in \Omega$, all $x \in \mathbb{R}$, with $a \in L^{\infty}(\Omega)$,

$$
p<r<p^{*}= \begin{cases}\frac{N p}{N-p}, & \text { if } p<N \\ +\infty, & \text { if } N \leq p\end{cases}
$$

(ii) $\limsup _{x \rightarrow+\infty} \frac{f(z, x)}{|x|^{p-2} x} \leq 0$ uniformly for a.a. $z \in \Omega$;
(iii) there exists $\vartheta>0$ such that

$$
\vartheta \leq \liminf _{x \rightarrow 0} \frac{f(z, x)}{|x|^{q-2} x} \text { uniformly for a.a. } z \in \Omega .
$$

Remark 2.3. Evidently the hypotheses on $f$ are very general and include also functions which may change sign as $x \rightarrow \pm \infty$. Note that near zero $f(z, x) x \geq 0$ for a.a. $z \in \Omega$.

Let $F(z, x)=\int_{0}^{x} f(z, s) d s$ (the primitive of $\left.f(z, \cdot)\right)$. We introduce the $C^{1}$-functionals $\varphi_{\lambda}, \varphi_{\lambda}^{ \pm}: W^{1, p}(\Omega) \mapsto \mathbb{R}$ defined by

$$
\begin{aligned}
& \varphi_{\lambda}(u)=\frac{1}{p} k_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\lambda \int_{\Omega} F(z, u) d z, \\
& \varphi_{\lambda}^{ \pm}(u)=\frac{1}{p} k_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\lambda \int_{\Omega} F\left(z, \pm u^{ \pm}\right) d z \text { for all } u \in W^{1, p}(\Omega) .
\end{aligned}
$$

## 3. Constant sign solutions

First we show that $\left(P_{\lambda}\right)$ has constant sign solutions for all $\lambda>0$.
Proposition 3.1. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}$ hold, then for every $\lambda>0$ problem $\left(P_{\lambda}\right)$ has at least two constant sign solutions $u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+}$.

Proof. First we show the existence of a positive solution. On account of hypotheses $\mathbf{H}_{\mathbf{1}}(i)$, (ii) given $\varepsilon>0$, we can find $C_{\varepsilon}>0$ such that

$$
\begin{equation*}
F(z, x) \leq \frac{\varepsilon}{p}|x|^{p}+C_{\varepsilon} \text { for a.a. } z \in \Omega, \text { all } x \in \mathbb{R} \tag{4}
\end{equation*}
$$

Then for all $u \in W^{1, p}(\Omega)$ we have

$$
\begin{aligned}
\varphi_{\lambda}^{+}(u) & \geq \frac{1}{p}\left(k_{p}(u)-\lambda \varepsilon\|u\|_{p}^{p}\right)-C_{3} \text { for some } C_{3}=C_{3}(\varepsilon)>0(\text { see }(4)) \\
& \geq \frac{1}{p}\left(C_{0}-\lambda \varepsilon\right)\|u\|^{p}-C_{3}(\text { see }(3))
\end{aligned}
$$

Choosing $\varepsilon \in\left(0, \frac{C_{0}}{\lambda}\right)$, we see that

$$
\varphi_{\lambda}^{+}(\cdot) \text { is coercive. }
$$

The Sobolev embedding theorem and the compactness of the trace map, imply that $\varphi_{\lambda}^{+}$ is sequentially weakly lower semicontinuous. Thus, by the Weierstrass-Tonelli theorem, we can find $u_{\lambda} \in W^{1, p}(\Omega)$ such that

$$
\begin{equation*}
\varphi_{\lambda}^{+}\left(u_{\lambda}\right)=\inf \left\{\varphi_{\lambda}^{+}(u): u \in W^{1, p}(\Omega)\right\} \tag{5}
\end{equation*}
$$

Hypothesis $\mathbf{H}_{\mathbf{1}}(i i i)$ implies that given $\varepsilon \in(0, \vartheta)$, we can find $\delta=\delta(\varepsilon) \in(0,1)$ such that

$$
\begin{equation*}
F(z, x) \geq \frac{1}{q}(\vartheta-\varepsilon)|x|^{q} \text { for a.a. } z \in \Omega, \text { all }|x| \leq \delta \tag{6}
\end{equation*}
$$

Let $\eta \in(0, \delta)$. Then

$$
\begin{align*}
\varphi_{\lambda}^{+}(\eta) & \leq \frac{\eta^{p}}{p}\left(\int_{\Omega} \xi(z) d z+\int_{\partial \Omega} \beta(z) d \sigma\right)-\frac{\eta^{q}}{q} \lambda(\vartheta-\varepsilon)(\text { see }(6)) \\
& =C_{4} \eta^{p}-C_{5} \eta^{q} \text { for some } C_{4}, C_{5}>0 \tag{7}
\end{align*}
$$

Since $q<p$, choosing $\eta \in(0, \delta)$ even smaller if necessary we have

$$
\begin{aligned}
& \varphi_{\lambda}^{+}(\eta)<0(\text { see }(7)) \\
\Rightarrow & \varphi_{\lambda}^{+}\left(u_{\lambda}\right)<0=\varphi_{\lambda}^{+}(0)(\text { see }(5)) \\
\Rightarrow & u_{\lambda} \neq 0
\end{aligned}
$$

From (5) we have

$$
\begin{align*}
&\left(\varphi_{\lambda}^{+}\right)^{\prime}\left(u_{\lambda}\right)=0 \\
& \Rightarrow\left\langle k_{p}^{\prime}\left(u_{\lambda}\right), h\right\rangle+\left\langle A_{q}\left(u_{\lambda}\right), h\right\rangle=\lambda \int_{\Omega} f\left(z, u_{\lambda}^{+}\right) h d z \tag{8}
\end{align*}
$$

for all $h \in W^{1, p}(\Omega)$.
In (8) we choose $h=-u_{\lambda}^{-} \in W^{1, p}(\Omega)$ and obtain

$$
\begin{aligned}
& k_{p}\left(u_{\lambda}^{-}\right) \leq 0 \\
\Rightarrow & C_{0}\left\|u_{\lambda}^{-}\right\|^{p} \leq 0(\text { see }(3)), \\
\Rightarrow & u_{\lambda} \geq 0, u_{\lambda} \neq 0
\end{aligned}
$$

Therefore $u_{\lambda}$ is a positive solution of $\left(P_{\lambda}\right)$. Proposition 2.10 of Papageorgiou \& Rădulescu [17], implies that $u_{\lambda} \in L^{\infty}(\Omega)$. Then using the nonlinear regularity theory of Lieberman [11], we have $u_{\lambda} \in C_{+} \backslash\{0\}$. Let $\rho=\left\|u_{\lambda}\right\|_{\infty}$. Hypotheses $\mathbf{H}_{\mathbf{1}}(i)$, (iii) imply that we can find $\hat{\xi}_{\rho}>0$ such that

$$
f(z, x) x+\hat{\xi}_{\rho}|x|^{p} \geq 0 \text { for a.a. } z \in \Omega, \text { all }|x| \leq \rho .
$$

We have

$$
\Delta_{p} u_{\lambda}+\Delta_{q} u_{\lambda} \leq\left(\|\xi\|_{\infty}+\lambda \hat{\xi}_{\rho}\right) u_{\lambda}^{p-1} \text { in } \Omega
$$

Then the maximum principle of Pucci \& Serrin [24, pp. 111, 120], implies that $u_{\lambda} \in$ $\operatorname{int} C_{+}$.

Similarly working with $\varphi_{\lambda}^{-}$, we produce a negative solution $v_{\lambda} \in-\operatorname{int} C_{+}$.
In fact we can show the existence of a smallest positive solution and of a biggest negative solution. We will need these extremal constant sign solutions in order to produce a nodal one (see Section 4).

To produce these extremal constant sign solutions, we need to do some preparatory work. Hypotheses $\mathbf{H}_{\mathbf{1}}(i)$, (iii) imply that given $\varepsilon \in(0, \vartheta)$, we can find $C_{6}=C_{6}(\varepsilon)>0$ such that

$$
\begin{equation*}
f(z, x) x \geq(\vartheta-\varepsilon)|x|^{q}-C_{6}|x|^{r} \text { for a.a. } z \in \Omega, \text { all } x \in \mathbb{R} . \tag{9}
\end{equation*}
$$

This unilateral growth condition on $f(z, \cdot)$ leads to the following auxiliary Robin problem

$$
\left\{\begin{array}{l}
-\Delta_{p} u-\Delta_{q} u+\xi(z)|u|^{p-2} u=\lambda\left((\vartheta-\varepsilon)|u|^{q-2} u-C_{6}|u|^{r-2} u\right) \text { in } \Omega, \\
\frac{\partial u}{\partial n_{p q}}+\beta(z)|u|^{p-2} u=0 \text { on } \partial \Omega, \lambda>0, \varepsilon \in(0, \vartheta) .
\end{array}\right.
$$

Proposition 3.2. If hypotheses $\mathbf{H}_{\mathbf{0}}$ hold, then for every $\lambda>0$ problem ( $10_{\lambda}$ ) has a unique positive solution $\bar{u}_{\lambda} \in \operatorname{int} C_{+}$and since the equation is odd $\bar{v}_{\lambda}=-\bar{u}_{\lambda} \in-\operatorname{int} C_{+}$is the unique negative solution of problem ( $10_{\lambda}$ ).

Proof. First we show the existence of a positive solution.
So, we consider the $C^{1}$-functional $\psi_{\lambda}^{+}: W^{1, p}(\Omega) \mapsto \mathbb{R}$ defined by

$$
\psi_{\lambda}^{+}(u)=\frac{1}{p} k_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\frac{\lambda(\vartheta-\varepsilon)}{q}\left\|u^{+}\right\|_{q}^{q}+\frac{\lambda C_{6}}{r}\left\|u^{+}\right\|_{r}^{r}
$$

for all $u \in W^{1, p}(\Omega)$.
Since $q<p<r$, it is clear that

$$
\psi_{\lambda}^{+} \text {is coercive. }
$$

Also, it is sequentially weakly lower semicontinuous. So, we can find $\bar{u}_{\lambda} \in W^{1, p}(\Omega)$ such that

$$
\begin{equation*}
\psi_{\lambda}^{+}\left(\bar{u}_{\lambda}\right)=\inf \left\{\psi_{\lambda}^{+}(u): u \in W^{1, p}(\Omega)\right\} \tag{10}
\end{equation*}
$$

Since $\varepsilon \in(0, \vartheta)$ and $q<p<r$, we see that for $\eta \in(0,1)$ small we have

$$
\begin{aligned}
& \psi_{\lambda}^{+}(\eta)<0 \\
\Rightarrow & \psi_{\lambda}^{+}\left(\bar{u}_{\lambda}\right)<0=\psi_{\lambda}^{+}(0)(\text { see }(10)), \\
\Rightarrow & \bar{u}_{\lambda} \neq 0
\end{aligned}
$$

From (10) we have

$$
\begin{gather*}
\left(\psi_{\lambda}^{+}\right)^{\prime}\left(\bar{u}_{\lambda}\right)=0 \\
\Rightarrow\left\langle k_{p}^{\prime}\left(\bar{u}_{\lambda}\right), h\right\rangle+\left\langle A_{q}\left(\bar{u}_{\lambda}\right), h\right\rangle=\lambda \int_{\Omega}\left((\vartheta-\varepsilon)\left|\bar{u}_{\lambda}\right|^{q-2} \bar{u}_{\lambda}-C_{6}\left|\bar{u}_{\lambda}\right|^{r-2} \bar{u}_{\lambda}\right) h d z \tag{11}
\end{gather*}
$$

for all $h \in W^{1, p}(\Omega)$.
In (11) we use the test function $h=-\bar{u}_{\lambda}^{-} \in W^{1, p}(\Omega)$ and using (3) we obtain that $\bar{u}_{\lambda} \geq 0, \bar{u}_{\lambda} \neq 0$. This implies that $\bar{u}_{\lambda}$ is a positive solution of $\left(10_{\lambda}\right)$. As before the nonlinear regularity theory and the nonlinear maximum principle imply that $\bar{u}_{\lambda} \in \operatorname{int} C_{+}$.

In what follows, $\hat{k}_{p}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ is the $C^{1}$-functional defined by

$$
\hat{k}_{p}(u)=\|D u\|_{p}^{p}+\int_{\Omega} \xi(z)|u|^{p} d z \text { for all } u \in W^{1, p}(\Omega)
$$

Next, we show the uniqueness of this positive solution. To this end, we introduce the integral functional $j: L^{1}(\Omega) \mapsto \overline{\mathbb{R}}=\mathbb{R} \cup\{+\infty\}$ defined by

$$
j(u)= \begin{cases}\frac{1}{p} \hat{k}_{p}\left(u^{1 / q}\right)+\frac{1}{q}\left\|D u^{1 / q}\right\|_{q}^{q}, & \text { if } u \geq 0, u^{1 / q} \in W^{1, p}(\Omega) \\ +\infty, & \text { otherwise }\end{cases}
$$

Let dom $j=\left\{u \in L^{1}(\Omega): j(u)<+\infty\right\}$ (the effective domain of $j(\cdot)$ ). We introduce function $G_{0}: \mathbb{R}_{+} \mapsto \mathbb{R}_{+}$defined by

$$
G_{0}(t)=\frac{1}{p} t^{p}+\frac{1}{q} t^{q} \text { for all } t \geq 0
$$

Evidently $G_{0}(\cdot)$ is increasing and $t \mapsto G_{0}\left(t^{1 / q}\right)$ is convex. We set $G(y)=G_{0}(|y|)$ for all $y \in \mathbb{R}^{N}$. Clearly $G(\cdot)$ is convex. So, if $u_{1}, u_{2} \in \operatorname{dom} j$ and $u=\left(t u_{1}+(1-t) u_{2}\right)^{1 / q}$, $t \in[0,1]$, then from Diaz \& Saa [4, Lemma 1], we have

$$
\begin{aligned}
& |D u| \leq\left(t\left|D u_{1}^{1 / q}\right|^{q}+(1-t)\left|D u_{2}^{1 / q}\right|^{q}\right)^{1 / q} \\
\Rightarrow & G_{0}(|D u|) \leq G_{0}\left(\left(t\left|D u_{1}^{1 / q}\right|^{q}+(1-t)\left|D u_{2}^{1 / q}\right|^{q}\right)^{1 / q}\right)
\end{aligned}
$$

(since $G_{0}(\cdot)$ is increasing),

$$
\leq t G_{0}\left(\left|D u_{1}^{1 / q}\right|\right)+(1-t) G_{0}\left(\left|D u_{2}^{1 / q}\right|\right)
$$

(since $t \mapsto G_{0}\left(t^{1 / q}\right)$ is convex),
$\Rightarrow G(D u) \leq t G\left(D u_{1}^{1 / q}\right)+(1-t) G\left(D u_{2}^{1 / q}\right)$,
$\Rightarrow j(\cdot)$ is convex (recall that $q<p$ and see hypotheses $\mathbf{H}_{\mathbf{0}}$ ).
Also, by Fatou's lemma we see that $j(\cdot)$ is lower semicontinuous.
Suppose $\tilde{u}_{\lambda}$ is another positive solution of problem $\left(10_{\lambda}\right)$. Again we have $\tilde{u}_{\lambda} \in \operatorname{int} C_{+}$. Hence using Proposition 4.1.22 of Papageorgiou, Rădulescu \& Repovš [21, p. 274], we have

$$
\frac{\bar{u}_{\lambda}}{\tilde{u}_{\lambda}} \in L^{\infty}(\Omega) \text { and } \frac{\tilde{u}_{\lambda}}{\bar{u}_{\lambda}} \in L^{\infty}(\Omega)
$$

Let $h=\bar{u}_{\lambda}^{q}-\tilde{u}_{\lambda}^{q}$. Then for $|t|<1$ small we have

$$
\bar{u}_{\lambda}^{q}+t h \in \operatorname{dom} j, \tilde{u}_{\lambda}^{q}+t h \in \operatorname{dom} j .
$$

Then we can calculate the Gâteaux (directional) derivative of $j(\cdot)$ at $\bar{u}_{\lambda}^{q}$ and at $\tilde{u}_{\lambda}^{q}$ in the direction $h$. In fact, using the chain rule and reasoning as in Gasiński and Papageorgiou [7, p. 492], we have

$$
\begin{aligned}
j^{\prime}\left(\bar{u}_{\lambda}^{q}\right)(h) & =\frac{1}{q}\left[\left\langle A_{p}\left(\bar{u}_{\lambda}\right), \frac{h}{\bar{u}_{\lambda}^{q-1}}\right\rangle+\left\langle A_{q}\left(\bar{u}_{\lambda}\right), \frac{h}{\bar{u}_{\lambda}^{q-1}}\right\rangle+\int_{\Omega} \frac{\xi(z) \bar{u}_{\lambda}^{p-1}}{\bar{u}_{\lambda}^{q-1}} h d z\right] \\
& =\frac{1}{q} \int_{\Omega} \frac{-\Delta_{p} \bar{u}_{\lambda}-\Delta_{q} \bar{u}_{\lambda}+\xi(z) \bar{u}_{\lambda}^{p-1}}{\bar{u}_{\lambda}^{q-1}} h d z \\
& =\frac{1}{q} \int_{\Omega} \lambda\left((\vartheta-\varepsilon)-C_{6} \bar{u}_{\lambda}^{r-q}\right) h d z
\end{aligned}
$$

(using Green's identity, see [21, p. 35]).
Similarly we have

$$
j^{\prime}\left(\tilde{u}_{\lambda}^{q}\right)(h)=\frac{1}{q} \int_{\Omega} \lambda\left((\vartheta-\varepsilon)-C_{6} \tilde{u}_{\lambda}^{r-q}\right) h d z
$$

The convexity of $j(\cdot)$ implies the monotonicity of $j^{\prime}(\cdot)$. Hence

$$
\begin{aligned}
0 & \leq \lambda C_{6} \int_{\Omega}\left(\tilde{u}_{\lambda}^{r-q}-\bar{u}_{\lambda}^{r-q}\right)\left(\bar{u}_{\lambda}^{q}-\tilde{u}_{\lambda}^{q}\right) d z, \\
\Rightarrow \bar{u}_{\lambda} & =\tilde{u}_{\lambda}
\end{aligned}
$$

This proves the uniqueness of the positive solution $\bar{u}_{\lambda} \in \operatorname{int} C_{+}$of problem (10 ${ }_{\lambda}$ ).
Since the equation is odd, then $\bar{v}_{\lambda}=-\bar{u}_{\lambda} \in-\operatorname{int} C_{+}$is the unique negative solution of problem $\left(10_{\lambda}\right)$.

We introduce the following two sets

$$
\begin{aligned}
& S_{\lambda}^{+}=\text {set of positive solutions of }\left(P_{\lambda}\right) \\
& S_{\lambda}^{-}=\text {set of negative solutions of }\left(P_{\lambda}\right)
\end{aligned}
$$

From Proposition 3.1 and its proof, we know that for all $\lambda>0$, we have

$$
\emptyset \neq S_{\lambda}^{+} \subseteq \operatorname{int} C_{+} \text {and } \emptyset \neq S_{\lambda}^{-} \subseteq-\operatorname{int} C_{+} .
$$

The solutions of $\left(10_{\lambda}\right)$ produced in Proposition 3.2 provide bounds for the two solution sets $S_{\lambda}^{+}, S_{\lambda}^{-}$.

Proposition 3.3. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}$ hold and $\lambda>0$, then $\bar{u}_{\lambda} \leq u$ for all $u \in S_{\lambda}^{+}$and $v \leq \bar{v}_{\lambda}$ for all $v \in S_{\lambda}^{-}$.

Proof. Let $u \in S_{\lambda}^{+} \subseteq \operatorname{int} C_{+}$. We consider the Carathéodory function $l(z, x)$ defined by

$$
l(z, x)= \begin{cases}(\vartheta-\varepsilon)\left(x^{+}\right)^{q-1}-C_{6}\left(x^{+}\right)^{r-1}, & \text { if } x \leq u(z)  \tag{12}\\ (\vartheta-\varepsilon) u(z)^{q-1}-C_{6} u(z)^{r-1}, & \text { if } u(z)<x\end{cases}
$$

(recall that $\varepsilon \in(0, \vartheta)$ ). We set $L(z, x)=\int_{0}^{x} l(z, s) d s$ and consider the $C^{1}$-functional $\hat{\psi}_{\lambda}: W^{1, p}(\Omega) \rightarrow \mathbb{R}$ defined by

$$
\hat{\psi}_{\lambda}(u)=\frac{1}{p} k_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\lambda \int_{\Omega} L(z, u) d z \text { for all } u \in W^{1, p}(\Omega)
$$

From (3) and (12), it is clear that $\hat{\psi}(\cdot)$ is coercive. Also it is sequentially weakly lower semicontinuous. So, we can find $\tilde{u}_{\lambda} \in W^{1, p}(\Omega)$ such that

$$
\begin{equation*}
\hat{\psi}_{\lambda}\left(\tilde{u}_{\lambda}\right)=\inf \left\{\hat{\psi}_{\lambda}(u): u \in W^{1, p}(\Omega)\right\} \tag{13}
\end{equation*}
$$

Since $q<p<r$, for $\eta \in(0,1)$ small we will have

$$
\begin{aligned}
& \hat{\psi}_{\lambda}(\eta)<0 \\
\Rightarrow & \hat{\psi}_{\lambda}\left(\tilde{u}_{\lambda}\right)<0=\hat{\psi}_{\lambda}(0) \\
\Rightarrow & \tilde{u}_{\lambda} \neq 0
\end{aligned}
$$

From (13) we have

$$
\begin{gather*}
\hat{\psi}_{\lambda}^{\prime}\left(\tilde{u}_{\lambda}\right)=0 \\
\Rightarrow\left\langle k_{p}^{\prime}\left(\tilde{u}_{\lambda}\right), h\right\rangle+\left\langle A_{q}\left(\tilde{u}_{\lambda}\right), h\right\rangle=\int_{\Omega} \lambda l(z, u) h d z \tag{14}
\end{gather*}
$$

for all $h \in W^{1, p}(\Omega)$.
Using $h=-\tilde{u}_{\lambda}^{-}$we obtain $\tilde{u}_{\lambda} \geq 0, \tilde{u}_{\lambda} \neq 0$. If we use $h=\left(\tilde{u}_{\lambda}-u\right)^{+} \in W^{1, p}(\Omega)$, then

$$
\begin{aligned}
&\left\langle k_{p}^{\prime}\left(\tilde{u}_{\lambda}\right),\left(\tilde{u}_{\lambda}-u\right)^{+}\right\rangle+\left\langle A_{q}\left(\tilde{u}_{\lambda}\right),\left(\tilde{u}_{\lambda}-u\right)^{+}\right\rangle \\
&= \lambda \int_{\Omega}\left[(\vartheta-\varepsilon) u^{q-1}-C_{6} u^{r-1}\right]\left(\tilde{u}_{\lambda}-u\right)^{+} d z \\
& \leq \lambda \int_{\Omega} f(z, u)\left(\tilde{u}_{\lambda}-u\right)^{+} d z(\text { see }(9)) \\
&=\left\langle k_{p}^{\prime}(u),\left(\tilde{u}_{\lambda}-u\right)^{+}\right\rangle+\left\langle A_{q}(u),\left(\tilde{u}_{\lambda}-u\right)^{+}\right\rangle\left(\text {since } u \in S_{\lambda}^{+}\right), \\
& \Rightarrow \tilde{u}_{\lambda} \leq u
\end{aligned}
$$

So, we have proved that

$$
\begin{equation*}
\tilde{u}_{\lambda} \in[0, u], \tilde{u}_{\lambda} \neq 0 \tag{15}
\end{equation*}
$$

From (15), (12) and (14) we see that

$$
\begin{aligned}
& \tilde{u}_{\lambda} \text { is a positive solution of problem }\left(10_{\lambda}\right), \\
\Rightarrow & \tilde{u}_{\lambda}=\bar{u}_{\lambda} \in \operatorname{int} C_{+}(\text {see Proposition 3.2). }
\end{aligned}
$$

Similarly we show that $v \leq \bar{v}_{\lambda}$ for all $v \in S_{\lambda}^{-} \subseteq-\operatorname{int} C_{+}$.
From Papageorgiou, Rădulescu \& Repovš [19] (see the proof of Proposition 7), we know that $S_{\lambda}^{+}$is downward directed (that is, if $u_{1}, u_{2} \in S_{\lambda}^{+}$, then there exists $u \in S_{\lambda}^{+}$ such that $u \leq u_{1}, u \leq u_{2}$ ) while $S_{\lambda}^{-}$is upward directed (that is, if $v_{1}, v_{2} \in S_{\lambda}^{-}$, then there exists $v \in S_{\lambda}^{-}$such that $v_{1} \leq v, v_{2} \leq v$ ). In the next proposition we establish the existence of extremal constant sign solutions.

Proposition 3.4. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}$ hold and $\lambda>0$, then problem $\left(P_{\lambda}\right)$ has a smallest positive solution $u_{\lambda}^{*} \in S_{\lambda}^{+} \subseteq \operatorname{int} C_{+}\left(\right.$that is, $u_{\lambda}^{*} \leq u$ for all $\left.u \in S_{\lambda}^{+}\right)$and a biggest negative solution $v_{\lambda}^{*} \in S_{\lambda}^{-} \subseteq-\operatorname{int} C_{+}\left(\right.$that is, $v \leq v_{\lambda}^{*}$ for all $\left.S_{\lambda}^{-}\right)$.

Proof. Using Lemma 3.10 of $\mathrm{Hu} \&$ Papageorgiou [9], we can find a decreasing sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq S_{\lambda}^{+}$such that

$$
\inf _{n \in \mathbb{N}} u_{n}=\inf S_{\lambda}^{+}
$$

We have

$$
\begin{align*}
& \left\langle k_{p}^{\prime}\left(u_{n}\right), h\right\rangle+\left\langle A_{q}\left(u_{n}\right), h\right\rangle=\lambda \int_{\Omega} f\left(z, u_{n}\right) h d z  \tag{16}\\
& \text { for all } h \in W^{1, p}(\Omega), \text { all } n \in \mathbb{N} \\
& \bar{u}_{\lambda} \leq u_{n} \leq u_{1} \text { for all } n \in \mathbb{N} \text { (see Proposition 3.3). } \tag{17}
\end{align*}
$$

In (16) we use the test function $h=u_{n} \in W^{1, p}(\Omega)$. Then we have

$$
\begin{aligned}
& C_{0}\left\|u_{n}\right\|^{p} \leq k_{p}\left(u_{n}\right) \leq C_{7} \text { for some } C_{7}>0, \text { all } n \in \mathbb{N} \\
& \left(\text { see }(17) \text { and hypothesis } \mathbf{H}_{\mathbf{1}}(i)\right), \\
\Rightarrow & \left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq W^{1, p}(\Omega) \text { is bounded. }
\end{aligned}
$$

So, we may assume that

$$
\begin{equation*}
u_{n} \xrightarrow{w} u_{\lambda}^{*} \text { in } W^{1, p}(\Omega) \text { and } u_{n} \rightarrow u_{\lambda}^{*} \text { in } L^{p}(\Omega) \text { and in } L^{p}(\partial \Omega) \tag{18}
\end{equation*}
$$

In (16) we use $h=u_{n}-u_{\lambda}^{*} \in W^{1, p}(\Omega)$, pass to the limit as $n \rightarrow \infty$ and use (18). We obtain

$$
\begin{align*}
& \lim _{n \rightarrow \infty}\left(\left\langle A_{p}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle+\left\langle A_{q}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle\right)=0 \\
\Rightarrow & \limsup \left(\left\langle A_{p}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle+\left\langle A_{q}\left(u_{\lambda}^{*}\right), u_{n}-u_{\lambda}^{*}\right\rangle\right) \leq 0 \\
& \left(\text { since } A_{q}(\cdot) \text { is monotone }\right), \\
\Rightarrow & \limsup _{n \rightarrow \infty}\left\langle A_{p}\left(u_{n}\right), u_{n}-u_{\lambda}^{*}\right\rangle \leq 0(\text { see }(18)), \\
\Rightarrow & \left.u_{n} \rightarrow u_{\lambda}^{*} \text { in } W^{1, p}(\Omega) \text { (by the }(S)_{+} \text {-property of } A_{p}(\cdot)\right) \tag{19}
\end{align*}
$$

If in (16) we pass to the limit as $n \rightarrow \infty$ and use (19), then

$$
\begin{aligned}
& \left\langle k_{p}^{\prime}\left(u_{\lambda}^{*}\right), h\right\rangle+\left\langle A_{q}\left(u_{\lambda}^{*}\right), h\right\rangle=\lambda \int_{\Omega} f\left(z, u_{\lambda}^{*}\right) h d z \text { for all } h \in W^{1, p}(\Omega) \\
& \bar{u}_{\lambda} \leq u_{\lambda}^{*}
\end{aligned}
$$

It follows that $u_{\lambda}^{*} \in S_{\lambda}^{+} \subseteq \operatorname{int} C_{+}$and $u_{\lambda}^{*}=\inf S_{\lambda}^{+}$.
Similarly we produce maximal negative solution $v_{\lambda}^{*} \in S_{\lambda}^{-} \subseteq-\operatorname{int} C_{+}$. In this case we can find an increasing sequence $\left\{v_{n}\right\}_{n \in \mathbb{N}} \subseteq S_{\lambda}^{-}$such that $\sup _{n \in \mathbb{N}} v_{n}=\sup S_{\lambda}^{-}$.

In the next section we use these extremal constant sign solutions in order to produce a nodal one.

## 4. Nodal solutions

To produce a nodal (sign-changing) solution, we look for nontrivial solutions of problem $\left(P_{\lambda}\right)$ in the order interval $\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right]$ distinct from $u_{\lambda}^{*}$ and $v_{\lambda}^{*}$. On account of the extremality of $u_{\lambda}^{*}$ and $v_{\lambda}^{*}$, any such solution is necessarily nodal. To limit ourselves on the order interval $\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right]$, we use truncations techniques. For this method to lead to the desired nodal solution, we need to restrict the parameter $\lambda>0$.

Let $u_{\lambda}^{*} \in \operatorname{int} C_{+}$and $v_{\lambda}^{*} \in-\operatorname{int} C_{+}$be the two extremal constant sign solutions produced in Proposition 3.4. We introduce the following truncation of $f(z, \cdot)$

$$
\hat{f}(z, x)= \begin{cases}f\left(z, v_{\lambda}^{*}(z)\right), & \text { if } x<v_{\lambda}^{*}(z)  \tag{20}\\ f(z, x), & \text { if } v_{\lambda}^{*}(z) \leq x \leq u_{\lambda}^{*}(z) \\ f\left(z, u_{\lambda}^{*}(z)\right), & \text { if } u_{\lambda}^{*}(z)<x\end{cases}
$$

This is a Carathéodory function. We also consider the positive and negative truncations of $f(z, \cdot)$, namely the Carathéodory functions

$$
\begin{equation*}
\hat{f}_{ \pm}(z, x)=\hat{f}\left(z, \pm x^{ \pm}\right) \tag{21}
\end{equation*}
$$

We set $\hat{F}(z, x)=\int_{0}^{x} \hat{f}(z, s) d s$ and $\hat{F}_{ \pm}(z, x)=\int_{0}^{x} \hat{f}_{ \pm}(z, s) d s$ and introduce the $C^{1}$ functionals $\hat{\varphi}_{\lambda}, \hat{\varphi}_{\lambda}^{ \pm}: W^{1, p}(\Omega) \mapsto \mathbb{R}$ defined by

$$
\begin{aligned}
& \hat{\varphi}_{\lambda}(u)=\frac{1}{p} k_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\lambda \int_{\Omega} \hat{F}(z, u) d z \\
& \hat{\varphi}_{\lambda}^{ \pm}(u)=\frac{1}{p} k_{p}(u)+\frac{1}{q}\|D u\|_{q}^{q}-\lambda \int_{\Omega} \hat{F}_{ \pm}(z, u) d z \text { for all } u \in W^{1, p}(\Omega)
\end{aligned}
$$

From (20), (21) and the extremality of $u_{\lambda}^{*}, v_{\lambda}^{*}$, we obtain easily the following result.
Proposition 4.1. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}$ hold and $\lambda>0$, then $K_{\hat{\varphi}_{\lambda}} \subseteq\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega})$, $K_{\hat{\varphi}_{\lambda}^{+}}=\left\{0, u_{\lambda}^{*}\right\}, K_{\hat{\varphi}_{\lambda}^{-}}=\left\{0, v_{\lambda}^{*}\right\}$.

Now we are ready to prove the existence of a nodal solution.
Proposition 4.2. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}$ hold and $\lambda>\frac{\hat{\lambda}_{2}(q)}{\vartheta}+1$, then problem $\left(P_{\lambda}\right)$ has a nodal solution

$$
y_{\lambda} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega})
$$

Proof. First we show that $u_{\lambda}^{*}$ and $v_{\lambda}^{*}$ are local minimizers of $\hat{\varphi}_{\lambda}(\cdot)$.

From (20) and (21) it is clear that $\hat{\varphi}_{\lambda}^{+}(\cdot)$ is coercive. Also, it is sequentially weakly lower semicontinuous. Hence we can find $\tilde{u}_{\lambda}^{*} \in W^{1, p}(\Omega)$ such that

$$
\begin{gathered}
\hat{\varphi}_{\lambda}^{+}\left(\tilde{u}_{\lambda}^{*}\right)=\inf \left\{\hat{\varphi}_{\lambda}^{+}(u): u \in W^{1, p}(\Omega)\right\}<0=\hat{\varphi}_{\lambda}^{+}(0) \\
\text { (see the proof of Proposition 3.1), } \\
\Rightarrow \tilde{u}_{\lambda}^{*} \neq 0
\end{gathered}
$$

It follows that $\tilde{u}_{\lambda}^{*} \in K_{\hat{\varphi}_{\lambda}^{+}} \backslash\{0\}$ and so using Proposition 4.1 we infer that

$$
\begin{equation*}
\tilde{u}_{\lambda}^{*}=u_{\lambda}^{*} \in \operatorname{int} C_{+} . \tag{22}
\end{equation*}
$$

From (20) and (21), we see that

$$
\left.\hat{\varphi}_{\lambda}^{+}\right|_{C_{+}}=\left.\hat{\varphi}_{\lambda}\right|_{C_{+}} .
$$

But then (22) implies that

$$
\begin{align*}
& u_{\lambda}^{*} \text { is a local } C^{1}(\bar{\Omega}) \text {-minimizer of } \hat{\varphi}_{\lambda}(\cdot) \\
\Rightarrow & u_{\lambda}^{*} \text { is a local } W^{1, p}(\Omega) \text {-minimizer of } \hat{\varphi}_{\lambda}(\cdot)  \tag{23}\\
& (\text { see Papageorgiou \& Rădulescu [17, Proposition 2.12]). }
\end{align*}
$$

Similarly, using the functional $\hat{\varphi}_{\lambda}^{-}$, we show

$$
\begin{equation*}
v_{\lambda}^{*} \text { is a local } W^{1, p}(\Omega) \text {-minimizer of } \hat{\varphi}_{\lambda}(\cdot) \text {. } \tag{24}
\end{equation*}
$$

We may assume that $\hat{\varphi}_{\lambda}\left(v_{\lambda}^{*}\right) \leq \hat{\varphi}_{\lambda}\left(u_{\lambda}^{*}\right)$. The reasoning is the same if the opposite inequality holds, using (24) instead of (23).

From Proposition 4.1, we see that we may assume that

$$
\begin{equation*}
K_{\hat{\varphi}_{\lambda}} \text { is finite. } \tag{25}
\end{equation*}
$$

Otherwise we already have a sequence of distinct smooth nodal solutions so we are done.

From (23), (25) and Theorem 5.7.6 of Papageorgiou, Rădulescu \& Repovš [21, p. 449], we can find $\rho \in(0,1)$ small such that

$$
\begin{equation*}
\hat{\varphi}_{\lambda}\left(v_{\lambda}^{*}\right) \leq \hat{\varphi}_{\lambda}\left(u_{\lambda}^{*}\right)<\inf \left\{\hat{\varphi}_{\lambda}(u):\left\|u-u_{\lambda}^{*}\right\|=\rho\right\}=\hat{m}_{\lambda},\left\|v_{\lambda}^{*}-u_{\lambda}^{*}\right\|>\rho . \tag{26}
\end{equation*}
$$

From [21] it follows that $\hat{\varphi}_{\lambda}(\cdot)$ is coercive. Hence by Proposition 5.1.15 of [21, p. 369] we obtain that

$$
\begin{equation*}
\hat{\varphi}_{\lambda}(\cdot) \text { satisfies the PS-condition. } \tag{27}
\end{equation*}
$$

Then on account of (26) and (27), we see that we can apply the Mountain Pass Theorem and produce $y_{\lambda} \in W^{1, p}(\Omega)$ such that

$$
\begin{aligned}
& y_{k} \in K_{\hat{\varphi}_{\lambda}} \subseteq\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega})\left(\text { see Proposition 4.1) and } \hat{m}_{\lambda} \leq \hat{\varphi}_{\lambda}\left(y_{\lambda}\right)(\text { see }(26)),\right. \\
\Rightarrow & y_{\lambda} \notin\left\{u_{\lambda}^{*}, v_{\lambda}^{*}\right\} .
\end{aligned}
$$

So, if we can show that $y_{\lambda} \neq 0$, then we can conclude that $y_{\lambda} \in C^{1}(\bar{\Omega})$ is a nodal solution of problem $\left(P_{\lambda}\right)$.

From the Mountain Pass Theorem, we know that

$$
\begin{equation*}
\hat{\varphi}_{\lambda}\left(y_{\lambda}\right)=\inf _{\gamma \in \Gamma} \max _{-1 \leq t \leq 1} \hat{\varphi}_{\lambda}(\gamma(t)) \tag{28}
\end{equation*}
$$

with $\Gamma=\left\{\gamma \in C\left([-1,1], W^{1, p}(\Omega)\right): \gamma(-1)=v_{\lambda}^{*}, \gamma(1)=u_{\lambda}^{*}\right\}$.
Let $\partial B_{1}^{L^{q}}, M$ be the manifolds from Proposition 2.1 and $M_{c}=M \cap C^{1}(\bar{\Omega})$. We introduce the following two sets of paths:

$$
\begin{aligned}
& \hat{\Gamma}=\left\{\hat{\gamma} \in C([-1,1], M): \hat{\gamma}(-1)=-\hat{u}_{1}(q), \hat{\gamma}(1)=\hat{u}_{1}(q)\right\} \\
& \hat{\Gamma}_{c}=\left\{\hat{\gamma} \in C\left([-1,1], M_{c}\right): \hat{\gamma}(-1)=-\hat{u}_{1}(q), \hat{\gamma}(1)=\hat{u}_{1}(q)\right\} .
\end{aligned}
$$

Claim: $\hat{\Gamma}_{c}$ is dense in $\hat{\Gamma}$.
Let $\hat{\gamma} \in \hat{\Gamma}$ and $\varepsilon \in(0,1)$. We introduce the multifunction $\hat{H}_{\varepsilon}:[-1,1] \mapsto 2^{C^{1}(\bar{\Omega})}$ defined by

$$
\hat{H}_{\varepsilon}(t)= \begin{cases}\left\{u \in C^{1}(\bar{\Omega}):\|u-\hat{\gamma}(t)\|<\varepsilon\right\}, & \text { if } t \in(-1,1) \\ \left\{ \pm \hat{u}_{1}(q)\right\}, & \text { if } t= \pm 1\end{cases}
$$

Evidently $\hat{H}_{\varepsilon}(\cdot)$ has nonempty convex values. Moreover, for $t \in(-1,1), \hat{H}_{\varepsilon}(t)$ is open, while $\hat{H}_{\varepsilon}( \pm 1)$ are singletons. In addition the continuity of $\hat{\gamma}(\cdot)$ implies the lower semicontinuity of the multifunction $\hat{H}_{\varepsilon}(\cdot)$ (see Proposition 2.6 of Hu \& Papageorgiou [9, p. 37]). Therefore we can use Theorem 3.1'I' of Michael [13] and have a continuous map $\hat{\gamma}_{\varepsilon}:[-1,1] \mapsto C^{1}(\bar{\Omega})$ such that $\hat{\gamma}_{\varepsilon}(t) \in \hat{H}_{\varepsilon}(t)$ for all $t \in[-1,1]$.

Now let $\varepsilon_{n}=\frac{1}{n}$ and $\hat{\gamma}_{n}=\hat{\gamma}_{\varepsilon_{n}} n \in \mathbb{N}$ as above. We have

$$
\begin{equation*}
\left\|\hat{\gamma}_{n}(t)-\hat{\gamma}(t)\right\|<\frac{1}{n} \text { for all } t \in[-1,1] \tag{29}
\end{equation*}
$$

Recall that $\hat{\gamma}(t) \in \partial B_{1}^{L^{q}}$ for all $t \in[-1,1]$. So, from (29) we see that we may assume that $\hat{\gamma}_{n}(t) \neq 0$ for all $t \in[-1,1]$, all $n \in \mathbb{N}$. We set

$$
\tilde{\gamma}_{n}(t)=\frac{\hat{\gamma}_{n}(t)}{\left\|\hat{\gamma}_{n}(t)\right\|_{q}} \text { for all } t \in[-1,1], \text { all } n \in \mathbb{N}
$$

We see that $\tilde{\gamma}_{n} \in C\left([-1,1], M_{c}\right), \tilde{\gamma}_{n}( \pm 1)= \pm \hat{u}_{1}(q)$ for all $n \in \mathbb{N}$.

Also we have

$$
\begin{align*}
\left\|\tilde{\gamma}_{n}(t)-\hat{\gamma}(t)\right\| \leq & \left\|\tilde{\gamma}_{n}(t)-\hat{\gamma}_{n}(t)\right\|+\left\|\hat{\gamma}_{n}(t)-\hat{\gamma}(t)\right\| \\
\leq & \frac{\left|1-\left\|\hat{\gamma}_{n}(t)\right\|_{q}\right|}{\left\|\hat{\gamma}_{n}(t)\right\|_{q}}\left\|\hat{\gamma}_{n}(t)\right\|+\frac{1}{n}  \tag{30}\\
& \quad \text { for all } t \in[-1,1], \text { all } n \in \mathbb{N}
\end{align*}
$$

Note that

$$
\begin{aligned}
& \max _{-1 \leq t \leq 1}\left|1-\left\|\hat{\gamma}_{n}(t)\right\|_{q}\right| \\
= & \left.\max _{-1 \leq t \leq 1}\left|\|\hat{\gamma}(t)\|_{q}-\left\|\hat{\gamma}_{n}(t)\right\|_{q}\right| \text { (recall that } \hat{\gamma} \in \hat{\Gamma}\right) \\
\leq & \max _{-1 \leq t \leq 1}\left\|\hat{\gamma}(t)-\hat{\gamma}_{n}(t)\right\|_{q} \\
\leq & C_{8} \max _{-1 \leq t \leq 1}\left\|\hat{\gamma}(t)-\hat{\gamma}_{n}(t)\right\| \text { for some } C_{8}>0\left(\text { since } W^{1, p}(\Omega) \hookrightarrow L^{q}(\Omega)\right) . \\
\leq & C_{8} \frac{1}{n} \text { for all } n \in \mathbb{N}(\text { see }(29)) .
\end{aligned}
$$

Let $m^{*}=\max _{-1 \leq t \leq 1}\|\hat{\gamma}(t)\|$ and $m_{n}^{*}=\max _{-1 \leq t \leq 1}\left\|\hat{\gamma}_{n}(t)\right\|$. We know that

$$
\begin{aligned}
& \left\|\hat{\gamma}_{n}(t)\right\| \leq \frac{1}{n}+\|\hat{\gamma}(t)\| \\
& \text { for all } t \in[-1,1], \text { all } n \in \mathbb{N}(\text { see }(29)) \\
\Rightarrow & m_{n}^{*} \leq \frac{1}{n}+m^{*} \\
\Rightarrow & \sup _{n \in \mathbb{N}} m_{n}^{*} \leq 1+m^{*}
\end{aligned}
$$

Also we have $\|\hat{\gamma}(t)\|_{q}=1$ (since $\hat{\gamma} \in \hat{\Gamma}$ ) and from (29) and since $W^{1, p}(\Omega) \hookrightarrow L^{q}(\Omega)$, we have

$$
\begin{aligned}
& \left\|\hat{\gamma}_{n}(t)-\hat{\gamma}(t)\right\|_{q} \leq \frac{C_{9}}{n} \text { for some } C_{9}>0, \text { all } n \in \mathbb{N}, \\
\Rightarrow & 1 \leq \frac{C_{9}}{n}+\left\|\hat{\gamma}_{n}(t)\right\| .
\end{aligned}
$$

So, if $m_{*}^{n}=\min _{-1 \leq t \leq 1}\left\|\hat{\gamma}_{n}(t)\right\|$, then $1 \leq \inf _{n \in \mathbb{N}} m_{*}^{n}$. Returning to (30), we have

$$
\begin{aligned}
& \left\|\tilde{\gamma}_{n}(t)-\hat{\gamma}(t)\right\| \leq \frac{1}{n}\left(C_{8}\left(1+m^{*}\right)+1\right) \\
\Rightarrow & \hat{\Gamma}_{c} \text { is dense in } \Gamma .
\end{aligned}
$$

This proves the Claim.

Using the Claim and Proposition 2.1, we see that given $\eta \in(0, \vartheta)$, we can find $\hat{\gamma} \in \hat{\Gamma}_{c}$ such that

$$
\begin{equation*}
\|D \hat{\gamma}(t)\|_{q}^{q} \leq \hat{\lambda}_{2}(q)+\eta \tag{31}
\end{equation*}
$$

Hypothesis $H_{1}(i i i)$ implies that we can find $\delta>0$ such that

$$
\begin{equation*}
F(z, x) \geq \frac{\eta}{q}|x|^{q} \text { for a.a. } z \in \Omega, \text { all }|x| \leq \delta \tag{32}
\end{equation*}
$$

The set $\hat{\gamma}([-1,1]) \subseteq M_{c}$ is compact. Recall that $u_{\lambda}^{*} \in \operatorname{int} C_{+}, v_{\lambda}^{*} \in-\operatorname{int} C_{+}$. So, using Proposition 4.1.22 of [21, p. 274] we can find $\mu \in(0,1)$ small such that

$$
\begin{cases}\mu \hat{\gamma}(t) \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega}), & \text { for all } t \in[-1,1]  \tag{33}\\ |\mu \hat{\gamma}(t)(z)| \leq \delta, & \text { for all } z \in \bar{\Omega}, \text { all } t \in[-1,1]\end{cases}
$$

Let $u \in \mu \hat{\gamma}([-1,1])$. We have $u=\mu \hat{u}$ with $\hat{u} \in \gamma([-1,1])$. Then

$$
\begin{aligned}
\hat{\varphi}_{\lambda}(u) \leq & \frac{\mu^{p}}{p} k_{p}(\hat{u})+\frac{\mu^{q}}{q}\left(\|D \hat{u}\|_{q}^{q}-\lambda \eta\|\hat{u}\|_{q}^{q}\right)(\text { see }(32),(33)) \\
\leq & \frac{\mu^{p}}{p} k_{p}(\hat{u})+\frac{\mu^{q}}{q}\left(\hat{\lambda}_{2}(q)+\eta-\lambda \eta\right) \\
& \quad\left(\text { see }(31) \text { and recall that }\|\hat{u}\|_{q}=1\right) .
\end{aligned}
$$

But $\lambda>\frac{\hat{\lambda}_{2}(q)}{\vartheta}+1 \Rightarrow \vartheta(\lambda-1)>\hat{\lambda}_{2}(q) \Rightarrow \eta(\lambda-1)>\hat{\lambda}_{2}(q)$ for $\eta$ near $\vartheta$. Therefore we have

$$
\hat{\varphi}_{\lambda}(u) \leq C_{10} \mu^{p}-C_{11} \mu^{q} \text { for some } C_{10}, C_{11}>0
$$

Since $q<p$, choosing $\mu \in(0,1)$ even smaller if necessary we have

$$
\begin{equation*}
\hat{\varphi}_{\lambda}(u)<0 \text { for all } u \in \mu \hat{\gamma}([-1,1]) . \tag{34}
\end{equation*}
$$

We set $\gamma_{0}=\mu \hat{\gamma}$. Then $\gamma_{0}$ is a continuous path connecting $-\mu \hat{u}_{1}(q)$ and $\mu \hat{u}_{1}(q)$ and

$$
\begin{equation*}
\left.\hat{\varphi}_{\lambda}\right|_{\gamma_{0}}<0(\operatorname{see}(34)) \tag{35}
\end{equation*}
$$

Next, we produce a continuous path connecting $\mu \hat{u}_{1}(q)$ and $u_{\lambda}^{*}$ and along this path $\hat{\varphi}_{\lambda}$ is negative.

So, let $a=\hat{\varphi}_{\lambda}^{+}\left(u_{\lambda}^{*}\right)=\varphi_{\lambda}\left(u_{\lambda}^{*}\right), b=0=\hat{\varphi}_{\lambda}^{+}(0)=\varphi_{\lambda}(0)$. Recall that $a<0=b$. Using Proposition 4.1, we have

$$
K_{\hat{\varphi}_{\lambda}^{+}}^{a}=\left\{u_{\lambda}^{*}\right\}, K_{\hat{\varphi}_{\lambda}^{+}}^{\circ}=\{0\} \text { and } \hat{\varphi}_{\lambda}^{+}\left(K_{\hat{\varphi}_{\lambda}^{+}}\right) \cap(a, 0)=\emptyset .
$$

Using the Second Deformation Theorem (see [21], Theorem 5.3.12, p. 386), we produce a deformation $\hat{h}:[0,1] \times\left(\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash\{0\} \rightarrow\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ}\right)$ such that

$$
\begin{align*}
& \hat{h}(0, u)=u \text { for all } u \in\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash\{0\}  \tag{36}\\
& \hat{h}(1, u)=u_{\lambda}^{*} \text { for all } u \in\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash\{0\}  \tag{37}\\
& \hat{h}\left(t, u_{\lambda}^{*}\right)=u_{\lambda}^{*} \text { for all } t \in[0,1]  \tag{38}\\
& \hat{\varphi}_{\lambda}^{+}(\hat{h}(t, u)) \leq \hat{\varphi}_{\lambda}^{+}(\hat{h}(s, u))  \tag{39}\\
& \text { for all } 0 \leq s \leq t \leq 1, \text { all } u \in\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash\{0\}
\end{align*}
$$

(recall from Section 2, that $\left(\hat{\varphi}_{\lambda}^{+}\right)^{0}=\left\{u \in W^{1, p}(\Omega): \hat{\varphi}_{\lambda}^{+}(u) \leq 0\right\}$ ).
From (36), (37), (38) we see that $K_{\hat{\varphi}_{\lambda}^{+}}^{a}=\left\{u_{\lambda}^{*}\right\}$ is a strong deformation retract of $\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash\{0\}=\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash K_{\hat{\varphi}_{\lambda}^{+}}^{\circ}$ and from (39) it follows that the deformation is $\hat{\varphi}_{\lambda}^{+}$-decreasing.

We set $\gamma_{+}(t)=\hat{h}\left(t, \mu \hat{u}_{1}(q)\right)^{+}$for all $0 \leq t \leq 1$. This is a continuous path in $W^{1, p}(\Omega)$ and since $\mu \hat{u}_{1}(q) \in\left(\hat{\varphi}_{\lambda}^{+}\right)^{\circ} \backslash\{0\}$ (see (35)), we have

$$
\begin{align*}
& \gamma_{+}(0)=\mu \hat{u}_{1}(q)(\text { see }(36)), \gamma_{+}(1)=u_{\lambda}^{*}(\text { see }(37)) \\
& \hat{\varphi}_{\lambda}\left(\gamma_{+}(t)\right)=\hat{\varphi}_{\lambda}^{+}\left(\gamma_{+}(t)\right) \leq \hat{\varphi}_{\lambda}^{+}\left(\gamma_{+}(0)\right)=\hat{\varphi}_{\lambda}^{+}\left(\mu \hat{u}_{1}(q)\right) \\
= & \hat{\varphi}_{\lambda}\left(\mu \hat{u}_{1}(q)\right)<0 \text { for all } t \in[0,1](\text { see }(35)), \\
\Rightarrow & \left.\hat{\varphi}_{\lambda}\right|_{\gamma_{+}}<0 . \tag{40}
\end{align*}
$$

In a similar fashion we produce another continuous path $\gamma_{-}(\cdot)$ in $W^{1, p}(\Omega)$ connecting $-\mu \hat{u}_{1}(q)$ and $v_{\lambda}^{*}$ such that

$$
\begin{equation*}
\left.\hat{\varphi}_{\lambda}\right|_{\gamma_{-}}<0 \tag{41}
\end{equation*}
$$

We concatenate $\gamma_{-}, \gamma_{0}, \gamma_{+}$and produce a path $\gamma_{*} \in \Gamma$ such that

$$
\begin{aligned}
& \left.\hat{\varphi}_{\lambda}\right|_{\gamma_{*}}<0(\operatorname{see}(35),(40),(41)), \\
\Rightarrow & \hat{\varphi}_{\lambda}\left(y_{\lambda}\right)<0=\hat{\varphi}_{\lambda}(0)(\operatorname{see}(28)), \\
\Rightarrow & y_{\lambda} \neq 0 .
\end{aligned}
$$

Therefore $y_{\lambda} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega})$ is a nodal solution of problem $\left(P_{\lambda}\right)$.
So, summarizing, we can state the following multiplicity theorem for problem $\left(P_{\lambda}\right)$. Note that we provide sign information for all solutions and the solutions are ordered.

Theorem 4.3. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}$ hold, then
(a) for all $\lambda>0$ problem $\left(P_{\lambda}\right)$ has constant sign solutions

$$
u_{\lambda} \in \operatorname{int} C_{+} \text {and } v_{\lambda} \in-\operatorname{int} C_{+}
$$

(b) for all $\lambda>\frac{\hat{\lambda}_{2}(q)}{\vartheta}+1$ problem $\left(P_{\lambda}\right)$ has at least three nontrivial solutions

$$
u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+}, y_{\lambda} \in\left[v_{\lambda}, u_{\lambda}\right] \cap C^{1}(\bar{\Omega}) \text { nodal. }
$$

If we introduce a symmetry hypothesis on $f(z, \cdot)$, we can have a whole sequence of nodal solutions converging to zero in $C^{1}(\bar{\Omega})$ and the result is valid for every parameter value $\lambda>0$. We introduce the following stronger version of hypothesis $\mathbf{H}_{\mathbf{1}}$.
$\mathbf{H}_{\mathbf{1}}^{\prime}$ : for a.a. $z \in \Omega, f(z, \cdot)$ is odd, hypotheses $\mathbf{H}_{\mathbf{1}}(i)$, (ii) hold and
(iii) $\lim _{x \rightarrow 0} \frac{f(z, x)}{|x|^{q-2} x}=+\infty$ uniformly for a.a. $z \in \Omega$.

Proposition 4.4. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}^{\prime}$ hold and $\lambda>0$, then problem $\left(P_{\lambda}\right)$ has a sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq C^{1}(\bar{\Omega})$ of nodal solutions such that $u_{n} \rightarrow 0$ in $C^{1}(\bar{\Omega})$.

Proof. From Proposition 3.4, we know that there exist extremal constant sign solutions

$$
u_{\lambda}^{*} \in \operatorname{int} C_{+} \text {and } v_{\lambda}^{*} \in-\operatorname{int} C_{+} .
$$

The energy functional $\varphi_{\lambda}$ is even (see hypotheses $\mathbf{H}_{\mathbf{1}}^{\prime}$ ) and coercive, thus it is bounded below. Hypothesis $\mathbf{H}_{\mathbf{1}}^{\prime}(i i i)$ implies that given any $\eta>0$, we can find $\delta=\delta(\eta)>0$ such that

$$
\begin{equation*}
F(z, x) \geq \frac{\eta}{q}|x|^{q} \text { for a.a. } z \in \Omega, \text { all }|x| \leq \delta \tag{42}
\end{equation*}
$$

Let $V \subseteq W^{1, p}(\Omega)$ be a finite dimensional subspace. Then on $V$ all norms are equivalent and so we can find $\rho_{V} \in(0,1)$ such that

$$
\begin{equation*}
u \in V \text { and }\|u\| \leq \rho_{V} \Rightarrow|u(z)| \leq \delta \text { for a.a. } z \in \Omega \tag{43}
\end{equation*}
$$

If $u \in V$ with $\|u\|=\rho_{V}$, then using (42) and (43) we have

$$
\begin{aligned}
\varphi_{\lambda}(u) & \leq \frac{1}{p} k_{p}(u)+\frac{1}{q}\left(\|D u\|_{q}^{q}-\eta\|u\|_{q}^{q}\right) \\
& \leq C_{10} \rho_{V}^{p}+\frac{1}{q}\left(C_{11}-\eta C_{V}\right) \rho_{V}^{q}
\end{aligned}
$$

for some $C_{10}, C_{11}, C_{V}>0$ (since all norms on $V$ are equivalent).

Recall that $\eta>0$ is arbitrary. So, we choose $\eta>\frac{C_{11}}{C_{V}}$ and have

$$
\varphi_{\lambda}(u) \leq C_{10} \rho_{V}^{p}-C_{12} \rho_{V}^{q} \text { for some } C_{12}>0 .
$$

Since $q<p$, choosing $\rho_{V} \in(0,1)$ small we have

$$
\sup \left\{\varphi_{\lambda}(u): u \in V,\|u\|=\rho_{V}\right\}<0 .
$$

Then we can apply Theorem 1 of Kajikiya [10] and produce a sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq$ $K_{\varphi_{\lambda}}$ such that

$$
\begin{equation*}
\varphi_{\lambda}\left(u_{n}\right) \leq 0 \text { and }\left\|u_{n}\right\| \rightarrow 0 \tag{44}
\end{equation*}
$$

The nonlinear regularity theory (see Lieberman [11]) implies that we can find $\alpha \in(0,1)$ and $C_{13}>0$ such that

$$
u_{n} \in C^{1, \alpha}(\bar{\Omega}),\left\|u_{n}\right\|_{C^{1, \alpha}(\bar{\Omega})} \leq C_{13} \text { for all } n \in \mathbb{N} .
$$

Exploiting the compact embedding of $C^{1, \alpha}(\bar{\Omega})$ into $C^{1}(\bar{\Omega})$ and using (44), we have

$$
\begin{aligned}
& u_{n} \rightarrow 0 \text { in } C^{1}(\bar{\Omega}), \\
\Rightarrow & u_{n} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega}) \text { for all } n \geq n_{0}, \\
\Rightarrow & \left\{u_{n}\right\}_{n \geq n_{0}} \text { is a sequence of nodal solutions of problem }\left(P_{\lambda}\right) .
\end{aligned}
$$

This completes the proof.
Using the same tools we can also treat the Dirichlet problem. So, now the problem under consideration is the following:

$$
\left\{\begin{array}{l}
-\Delta_{p} u(z)-\Delta_{q} u(z)=\lambda f(z, u(z)) \text { in } \Omega \\
\left.u\right|_{\partial \Omega}=0,1<q<p, \lambda>0 .
\end{array}\right.
$$

We know that the $q$-Laplace differential operator with Dirichlet boundary condition, has a smallest eigenvalue $\hat{\lambda}_{1}(q)>0$. Then Theorem 4.3 takes the following form.

Theorem 4.5. If hypotheses $\mathbf{H}_{\mathbf{1}}$ hold, then
(a) for all $\lambda>\hat{\lambda}_{1}(q)$ problem $\left(P_{\lambda}^{\prime}\right)$ has constant sign solutions

$$
u_{\lambda} \in \operatorname{int} C_{+} \text {and } v_{\lambda} \in-\operatorname{int} C_{+}
$$

(b) for all $\lambda>\frac{\hat{\lambda}_{2}(q)}{\vartheta}+1$ problem $\left(P_{\lambda}^{\prime}\right)$ has at least three nontrivial solutions

$$
u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+} \text {and } y_{\lambda} \in\left[v_{\lambda}, u_{\lambda}\right] \cap C^{1}(\bar{\Omega}) \text { nodal. }
$$

Similarly Proposition 4.4 is also valid but with $\lambda>\hat{\lambda}_{1}(q)$.
Proposition 4.6. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{1}}^{\prime}$ hold and $\lambda>\hat{\lambda}_{1}(q)$, then problem $\left(P_{\lambda}^{\prime}\right)$ has a sequence $\left\{u_{n}\right\}_{n \in \mathbb{N}} \subseteq C^{1}(\bar{\Omega})$ of nodal solutions such that $u_{n} \rightarrow 0$ in $C^{1}(\bar{\Omega})$.

## 5. ( $p, 2$ )-equations

When $q=2$ (that is, we deal with a $(p, 2)$-equation) and we strengthen the regularity of $f(z, \cdot)$, then we can produce a second nodal solution, for a total of four nontrivial smooth solutions all with sign information.

So, the Robin problem under consideration, is the following

$$
\left\{\begin{array}{l}
-\Delta_{p} u(z)-\Delta u(z)+\xi(z)|u(z)|^{p-2} u(z)=\lambda f(z, u(z)) \text { in } \Omega \\
\frac{\partial u}{\partial n_{p 2}}+\beta(z)|u|^{p-2} u=0 \text { on } \partial \Omega, 1<2<p, \lambda>0
\end{array}\right.
$$

Now the hypotheses of the reaction $f(z, x)$ are the following:
$\mathbf{H}_{\mathbf{2}}: f: \Omega \times \mathbb{R} \mapsto \mathbb{R}$ is a measurable function such that for a.a. $z \in \Omega, f(z, 0)=0$, $f(z, \cdot) \in C^{1}(\mathbb{R})$ and
(i) $\left|f_{x}^{\prime}(z, x)\right| \leq a(z)\left(1+|x|^{r-2}\right)$ for a.a. $z \in \Omega$, all $x \in \mathbb{R}$ with $a \in L^{\infty}(\Omega)$ and $p<r<$ $p^{*}$;
(ii) $\limsup _{x \rightarrow \pm \infty} \frac{f(z, x)}{|x|^{p-2} x} \leq 0$ uniformly for a.a. $z \in \Omega$;
(iii) there exists $m \in \mathbb{N}, m \geq 2$ such that

$$
\begin{aligned}
& f_{x}^{\prime}(z, 0) \in\left[\hat{\lambda}_{m}(2), \hat{\lambda}_{m+1}(2)\right] \text { for a.a. } z \in \Omega \\
& f_{x}^{\prime}(\cdot, 0) \not \equiv \hat{\lambda}_{m}(2), f_{x}^{\prime}(\cdot, 0) \not \equiv \hat{\lambda}_{m+1}(2) \\
& f_{x}^{\prime}(z, 0)=\lim _{x \rightarrow 0} \frac{f(z, x)}{x} \text { uniformly for a.a. } z \in \Omega
\end{aligned}
$$

We introduce the functional $\hat{\tau}_{\lambda}: H^{1}(\Omega) \mapsto \mathbb{R}$ defined by

$$
\hat{\tau}_{\lambda}(u)=\frac{1}{2}\|D u\|_{2}^{2}-\lambda \int_{\Omega} F(z, u) d z \text { for all } u \in H^{1}(\Omega)
$$

Note that $\hat{\tau}_{\lambda} \in C^{2}\left(H^{1}(\Omega)\right)$. We consider the functional

$$
\tau_{\lambda}=\left.\hat{\tau}_{\lambda}\right|_{W^{1, p}(\Omega)}(\text { recall that } 2<p)
$$

Proposition 5.1. If hypotheses $\mathbf{H}_{\mathbf{2}}$ hold, then $C_{k}\left(\tau_{\lambda}, 0\right)=\delta_{k, d_{m}} \mathbb{Z}$ for all $k \in \mathbb{N}_{0}$, with $d_{m}=\operatorname{dim} \bar{H}_{m}$.

Proof. As we already mentioned, $\hat{\tau}_{k} \in C^{2}\left(H^{1}(\Omega)\right)$ and if by $\langle\cdot, \cdot\rangle_{H^{1}}$ we denote the duality brackets for the pair $\left(H^{1}(\Omega), H^{1}(\Omega)^{*}\right)$, we have

$$
\begin{align*}
\left\langle\hat{\tau}_{\lambda}^{\prime \prime}(u) v, h\right\rangle_{H^{1}}= & \int_{\Omega}(D v, D h)_{\mathbb{R}^{N}} d z-\lambda \int_{\Omega} f_{x}^{\prime}(z, u) v h d z  \tag{45}\\
& \text { for all } u, v, h \in H^{1}(\Omega)
\end{align*}
$$

Suppose that $v \in N\left(\hat{\tau}_{\lambda}^{\prime \prime}(0)\right)=\operatorname{ker}\left(\hat{\tau}_{\lambda}^{\prime \prime}(0)\right)$. We have the unique orthogonal decomposition $v=\bar{v}+\hat{v}$ with $\bar{v} \in \bar{H}_{m}$ and $\hat{v} \in \hat{H}_{m+1}=\bar{H}_{m}^{\perp}$. In (45) let $u=0, v \in N\left(\hat{\tau}_{\lambda}^{\prime \prime}(0)\right)$ and choose $h=\hat{v}$. Exploiting the orthogonality of $\bar{H}_{m}$ and $\hat{H}_{m+1}$ and hypothesis $\mathbf{H}_{\mathbf{2}}(i i i)$, we obtain

$$
\begin{align*}
& \|D \hat{v}\|_{2}^{2}=\int_{\Omega} f_{x}^{\prime}(z, 0) \hat{v}^{2} d z \leq \hat{\lambda}_{m+1}(2)\|\hat{v}\|_{2}^{2}  \tag{46}\\
\Rightarrow & \hat{v} \in E\left(\hat{\lambda}_{m+1}(2)\right)(\text { see }(2))
\end{align*}
$$

If $\hat{v} \neq 0$, then by the UCP (see de Figueiredo \& Gossez [5]) we have that $\hat{v}(z) \neq 0$ for a.a. $z \in \Omega$ and so from (46) and hypothesis $\mathbf{H}_{2}(i i i)$, we have

$$
\|D \hat{v}\|_{2}^{2}<\hat{\lambda}_{m+1}(2)\|\hat{v}\|_{2}^{2}
$$

a contradiction (see (2)). Hence $\hat{v}=0$. Similarly, we show that $\bar{v}=0$ and so finally $v=0$. Therefore $u=0$ is nondegenerate critical point of $\hat{\tau}_{\lambda}$ with Morse index $\hat{d}_{m}$ and so from Proposition 6.2.6 of [21, p. 479], we have

$$
\begin{equation*}
C_{k}\left(\hat{\tau}_{\lambda}, 0\right)=\delta_{k, d_{m}} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0} \tag{47}
\end{equation*}
$$

We know that $W^{1, p}(\Omega) \hookrightarrow H^{1}(\Omega)$ densely and so by Theorem 6.6 .26 of [21, p. 545], we have

$$
\begin{aligned}
C_{k}\left(\tau_{\lambda}, 0\right) & =C_{k}\left(\hat{\tau}_{\lambda}, 0\right) \text { for all } k \in \mathbb{N}_{0}, \\
\Rightarrow C_{k}\left(\tau_{\lambda}, 0\right) & =\delta_{k, d_{m}} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0}(\text { see }(47))
\end{aligned}
$$

The proof is now complete.
Using this proposition, we can have a second nodal solution.
Proposition 5.2. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{2}}$ hold and $\lambda>\frac{\hat{\lambda}_{2}(2)}{\hat{\lambda}_{m}(2)}+1$, then problem ( $Q_{\lambda}$ ) has at least two nodal solutions

$$
y_{\lambda}, \hat{y}_{\lambda} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] .
$$

Proof. From Theorem 4.3 we already have a nodal solution

$$
y_{\lambda} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega})
$$

Let $a: \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$ be the map defined by

$$
a(y)=|y|^{p-2} y+y \text { for all } y \in \mathbb{R}^{N} .
$$

Since $p>2$, we see that $a \in C^{1}\left(\mathbb{R}^{N}, \mathbb{R}^{N}\right)$ and

$$
\nabla a(y)=|y|^{p-2}\left[i d+(p-2) \frac{y \otimes y}{|y|^{2}}\right]+i d \text { for all } y \in \mathbb{R}^{N} \backslash\{0\}
$$

We have

$$
(\nabla a(y) \xi, \xi)_{\mathbb{R}^{N}} \geq|\xi|^{2} \text { for all } y \in \mathbb{R}^{N} \backslash\{0\}, \xi \in \mathbb{R}^{N}
$$

Since $u_{\lambda}^{*} \in \operatorname{int} C_{+}$and $v_{\lambda}^{*} \in-\operatorname{int} C_{+}$, using the tangency principle of Pucci \& Serrin [24, p. 35], we have

$$
v_{\lambda}^{*}(z)<y_{\lambda}(z)<u_{\lambda}^{*}(z) \text { for all } z \in \Omega
$$

Consider the following open cone in $C^{1}(\bar{\Omega})$

$$
D_{+}=\left\{u \in C^{1}(\bar{\Omega}): u(z)>0 \text { for all } z \in \Omega,\left.\frac{\partial u}{\partial n}\right|_{\partial \Omega \cap u^{-1}(0)}<0\right\}
$$

From Proposition 3.2 of Gasiński \& Papageorgiou [8], we have $u_{\lambda}^{*}-y_{\lambda} \in D_{+}$and $y_{\lambda}-v_{\lambda}^{*} \in D_{+}$. Therefore

$$
\begin{equation*}
y_{\lambda} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] . \tag{48}
\end{equation*}
$$

Using (48) and the standard homotopy invariance argument, we obtain

$$
\begin{equation*}
C_{k}\left(\varphi_{\lambda}, y_{\lambda}\right)=C_{k}\left(\hat{\varphi}_{\lambda}, y_{\lambda}\right) \text { for all } k \in \mathbb{N}_{0} \tag{49}
\end{equation*}
$$

with $\varphi_{\lambda}(\cdot)$ and $\hat{\varphi}_{\lambda}(\cdot)$ as before, only now $q=2$. Recall that $y_{\lambda}$ is a critical point of mountain pass-type for $\hat{\varphi}_{\lambda}(\cdot)$, hence

$$
\begin{equation*}
C_{1}\left(\hat{\varphi}_{\lambda}, y_{\lambda}\right) \neq 0(\text { see }[21, \text { p. } 527]) \tag{50}
\end{equation*}
$$

We assume that $K_{\hat{\varphi}_{\lambda}}$ is finite or otherwise we already have an infinity of nodal solutions and so we are done. Since now on account of hypotheses $\mathbf{H}_{\mathbf{2}}, \varphi_{\lambda} \in C^{2}\left(W^{1, p}(\Omega)\right)$, as in Papageorgiou \& Rădulescu [16] (p. 414, Claim 3), using (49) and (50), we have

$$
\begin{align*}
& C_{k}\left(\varphi_{\lambda}, y_{\lambda}\right)=\delta_{k, 1} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0}, \\
& \Rightarrow C_{k}\left(\hat{\varphi}_{\lambda}, y_{\lambda}\right)=\delta_{k, 1} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0} \text { (see (49)). } \tag{51}
\end{align*}
$$

The $C^{1}$-continuity property of critical groups (see [21, p. 503]) implies that

$$
\begin{align*}
& C_{k}\left(\varphi_{\lambda}, 0\right)=C_{k}\left(\tau_{\lambda}, 0\right) \text { for all } k \in \mathbb{N}_{0} \\
\Rightarrow & C_{k}\left(\varphi_{\lambda}, 0\right)=\delta_{k, d_{m}} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0} \\
& (\text { see Proposition } 5.1), \\
\Rightarrow & C_{k}\left(\hat{\varphi}_{\lambda}, 0\right)=\delta_{k, d_{m}} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0}(\text { see (49)). } \tag{52}
\end{align*}
$$

From the proof of Proposition 4.2, we know that $u_{\lambda}^{*}$ and $v_{\lambda}^{*}$ are local minimizers of $\hat{\varphi}_{\lambda}(\cdot)$. Hence

$$
\begin{equation*}
C_{k}\left(\hat{\varphi}_{\lambda}, u_{\lambda}^{*}\right)=C_{k}\left(\hat{\varphi}_{\lambda}, v_{\lambda}^{*}\right)=\delta_{k, 0} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0} . \tag{53}
\end{equation*}
$$

Recall that $\hat{\varphi}_{\lambda}(\cdot)$ is coercive (see (20)). Therefore

$$
\begin{equation*}
C_{k}\left(\hat{\varphi}_{\lambda}, \infty\right)=\delta_{k, 0} \mathbb{Z} \text { for all } k \in \mathbb{N}_{0}(\text { see }[21, \text { p. 491]) } \tag{54}
\end{equation*}
$$

Suppose $K_{\hat{\varphi}_{\lambda}}=\left\{y_{\lambda}, 0, u_{\lambda}^{*}, v_{\lambda}^{*}\right\}$. From (51), (52), (53), (54) and the Morse relation (see Section 2), with $t=-1$, we have

$$
\begin{aligned}
& (-1)^{1}+(-1)^{d_{m}}+2(-1)^{0}=(-1)^{0} \\
\Rightarrow & (-1)^{d_{m}}=0, \text { a contradiction }
\end{aligned}
$$

So, there exists $\hat{y}_{\lambda} \in K_{\hat{\varphi}_{\lambda}}, \hat{y}_{\lambda} \notin\left\{y_{\lambda}, 0, u_{\lambda}^{*}, v_{\lambda}^{*}\right\}$. We have
$\hat{y}_{\lambda} \in\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] \cap C^{1}(\bar{\Omega})$ (see Proposition 4.1),
$\Rightarrow \hat{y}_{\lambda} \in C^{1}(\bar{\Omega})$ is a nodal solution of problem $\left(Q_{\lambda}\right)$.

Moreover, as we did for $y_{\lambda}$, we show that

$$
\hat{y}_{\lambda} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left[v_{\lambda}^{*}, u_{\lambda}^{*}\right] .
$$

This completes the proof.

So, for the problem $\left(Q_{\lambda}\right)$ we can state the following multiplicity theorem.
Theorem 5.3. If hypotheses $\mathbf{H}_{\mathbf{0}}, \mathbf{H}_{\mathbf{2}}$ hold and $\lambda>\frac{\hat{\lambda}_{2}(2)}{\hat{\lambda}_{m}(2)}+1$, then problem $\left(Q_{\lambda}\right)$ has at least four nontrivial solutions

$$
\begin{aligned}
& u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+}, \\
& y_{\lambda}, \hat{y}_{\lambda} \in \operatorname{int}_{C^{1}(\bar{\Omega})}\left[v_{\lambda}, u_{\lambda}\right] \text { nodal. }
\end{aligned}
$$

The same multiplicity theorem is also true for the Dirichlet problem

$$
\left\{\begin{array}{l}
-\Delta_{p} u(z)-\Delta u(z)=\lambda f(z, u(z)) \text { in } \Omega \\
\left.u\right|_{\partial \Omega}=0,2<p, \lambda>0
\end{array}\right.
$$

Theorem 5.4. If hypotheses $\mathbf{H}_{2}$ hold and $\lambda>\frac{\hat{\lambda}_{2}(2)}{\hat{\lambda}_{m}(2)}+1$, then problem $\left(Q_{\lambda}^{\prime}\right)$ has at least four nontrivial solutions

$$
\begin{aligned}
& u_{\lambda} \in \operatorname{int} C_{+}, v_{\lambda} \in-\operatorname{int} C_{+} \\
& y_{\lambda}, \hat{y}_{\lambda} \in \operatorname{int}_{C_{0}^{1}}(\bar{\Omega})\left[v_{\lambda}, u_{\lambda}\right] \text { nodal. }
\end{aligned}
$$

Remark 5.5. Another multiplicity theorem for ( $p, 2$ )-equations under different hypotheses can be found in [22].

## Declaration of competing interest

The authors declare that there is no conflict of interests regarding the publication of this paper.

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