# Generalized Choquard Equations Driven by Nonhomogeneous Operators 

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#### Abstract

In this work we prove the existence of solutions for a class of generalized Choquard equations involving the $\Delta_{\Phi}$-Laplacian operator. Our arguments are essentially based on variational methods. One of the main difficulties in this approach is to use the Hardy-LittlewoodSobolev inequality for nonlinearities involving N-functions. The methods developed in this paper can be extended to wide classes of nonlinear problems driven by nonhomogeneous operators.


Mathematics Subject Classification. 35A15, 35J62, 35J60.
Keywords. Choquard equation, variational methods, nonlinear elliptic equation, Hardy-Littlewood-Sobolev inequality.

## 1. Introduction

The stationary Choquard equation

$$
\begin{equation*}
-\Delta u+V(x) u=\left(\int_{\mathbb{R}^{N}} \frac{|u|^{p}}{|x-y|^{\lambda}} \mathrm{d} x\right)|u|^{p-2} u \quad \text { in } \mathbb{R}^{N}, \tag{1.1}
\end{equation*}
$$

where $N \geq 3,0<\lambda<N$, has appeared in the context of various physical models. In particular, this equation plays particularly an important role in the theory of Bose-Einstein condensation where it accounts for the finite-range many-body interactions.

For $N=3, p=2$, and $\lambda=1$, problem (1.1) was investigated by Fröhlich [1] and Pekar [2] in relationship with the quantum theory of a polaron, where free electrons in an ionic lattice interact with phonons associated to deformations of the lattice or with the polarisation that it creates on the medium (interaction of an electron with its own hole). We recall that Choquard [3] used this equation in the Hartree-Fock theory of one-component plasma. This equation was also proposed by Penrose in [4] as a model of self-gravitating matter and is known in that context as the Schrödinger-Newton equation. In fact, the Choquard equation is also known as the Schrödinger-Newton
equation in models coupling the Schrödinger equation of quantum physics together with nonrelativistic Newtonian gravity.

Recent relevant contributions included in the papers are by Ackermann [5], Alves et al. [6-8], Cingolani, Secchi and Squassina [9], Gao and Yang [10], Lions [11], Ma and Zhao [12], Moroz and van Schaftingen [13-16], van Schaftingen and Xia [17], Wang [18], and their references.

In all the above-mentioned papers, the authors used variational methods to show the existence of solution. This method works well thanks to a Hardy-Littlewood-Sobolev type inequality [19] that has the following statement.

Proposition 1.1 (Hardy-Littlewood-Sobolev inequality). Let $t, r>1$ and $0<\lambda<N$ with $1 / t+\lambda / N+1 / r=2, g \in L^{t}\left(\mathbb{R}^{N}\right)$ and $h \in L^{r}\left(\mathbb{R}^{N}\right)$. Then there exists a sharp constant $C(t, N, \lambda, r)$, independent of $f, h$, such that

$$
\begin{equation*}
\left|\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{g(x) h(y)}{|x-y|^{\lambda}} d x d y\right| \leq C(t, N, \lambda, r)\|g\|_{L^{t}\left(\mathbb{R}^{N}\right)}\|h\|_{L^{r}\left(\mathbb{R}^{N}\right)} \tag{1.2}
\end{equation*}
$$

In this paper, we are concerned with the existence of solution for the following class of quasilinear problems:

$$
\left\{\begin{align*}
-\Delta_{\Phi} u+\phi(|u|) u= & \left(\int_{\mathbb{R}^{N}} \frac{F(u(x))}{|x-y|^{\lambda}} \mathrm{d} x\right) f(u(y)) \quad \text { in } \mathbb{R}^{N},  \tag{1.3}\\
& u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) .
\end{align*}\right.
$$

This problem generalizes (1.1) in a nonhomogeneous setting, where $f$ : $\mathbb{R} \rightarrow \mathbb{R}$ is a continuous function verifying some natural hypotheses, which will be mentioned in Sect. 3. Let $F$ be the primitive of $f$, that is,

$$
F(t)=\int_{0}^{t} f(s) \mathrm{d} s
$$

We denote $\Delta_{\Phi}=\operatorname{div}(\phi(|\nabla u|) \nabla u)$, where $\Phi: \mathbb{R} \rightarrow \mathbb{R}$ is a $N$-function of the form

$$
\begin{equation*}
\Phi(t):=\int_{0}^{|t|} \phi(s) s \mathrm{~d} s \tag{1.4}
\end{equation*}
$$

with $\phi:(0,+\infty) \rightarrow(0,+\infty)$ being a $C^{1}$ function satisfying

$$
\begin{equation*}
(t \phi(t))^{\prime}>0 ; \quad \forall t>0 \tag{1}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow 0^{+}} t \phi(t)=0, \quad \lim _{t \rightarrow+\infty} t \phi(t)=+\infty . \tag{2}
\end{equation*}
$$

We also assume that there exist $l, m \in(1, N)$ such that $l \leq m<l^{*}:=\frac{N l}{N-l}$ and
$\left(\phi_{3}\right) \quad l \leq m<l^{*}:=\frac{N l}{N-l} \quad$ and $\quad l \leq \frac{\phi(t) t^{2}}{\Phi(t)} \leq m$ for all $t>0$.
Our purpose is to show that the variational method can be used to establish the existence of solutions for problem (1.3). One of the main difficulties is to show that the energy functional associated with (1.3) given
by

$$
J(u)=\int_{\mathbb{R}^{N}}(\Phi(|\nabla u|)+\Phi(|u|)) \mathrm{d} x-\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) F(u(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y
$$

is well defined and belongs to $C^{1}\left(W^{1, \Phi}\left(\mathbb{R}^{N}\right), \mathbb{R}\right)$. In fact, the main difficulty is to prove that the functional $\Psi: W^{1, \Phi}\left(\mathbb{R}^{N}\right) \rightarrow \mathbb{R}$ given by

$$
\begin{equation*}
\Psi(u)=\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) F(u(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \tag{1.5}
\end{equation*}
$$

belongs to $C^{1}\left(W^{1, \Phi}\left(\mathbb{R}^{N}\right), \mathbb{R}\right)$ with

$$
\Psi^{\prime}(u) v=\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x) f(u(y)) v(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y, \quad \forall u, v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)
$$

In what follows, we would like to point out that the operator $\Delta_{\Phi}$ arises in several physical applications, such as:
Non-Newtonian fluids: $\quad \Phi(t)=\frac{1}{p}|t|^{p}$ for $p>1$,
Plasma physics: $\quad \Phi(t)=\frac{1}{p}|t|^{p}+\frac{1}{q}|t|^{q}$ where $1<p<q<N$ with $q \in\left(p, p^{*}\right)$,
Nonlinear elasticity: $\Phi(t)=\left(1+t^{2}\right)^{\alpha}-1, \alpha \in\left(1, \frac{N}{N-2}\right)$,
Plasticity: $\Phi(t)=t^{p} \ln (1+t), 1<\frac{-1+\sqrt{1+4 N}}{2}<p<N-1, N \geq 3$,
Generalized Newtonian fluids: $\Phi(t)=\int_{0}^{t} s^{1-\alpha}\left(\sinh ^{-1} s\right)^{\beta} \mathrm{d} s, 0 \leq \alpha \leq 1$ and $\beta>0$.

The reader can find more details about the physical applications in [20], [21] and their references. The existence of solution for

$$
-\Delta_{\Phi} u+V(x) \phi(|u|) u=f(u) \quad \text { in } \quad \Omega, \quad \text { and } \quad u=0 \quad \text { on } \quad \partial \Omega,
$$

with $\Omega \subset \mathbb{R}^{N}$ being a bounded or unbounded domain has been established in some papers, see for example $[22-35]$ and the references therein.

This paper is organized as follows. In Sect. 2 we recall some facts involving Orlicz-Sobolev spaces. In Sect. 3 we show that $\Psi$ is of class $C^{1}$ under certain conditions on $f$. In Sect. 4 we study the existence of solutions to problem (1.3). Finally, in Sect. 5, we show other problems that can be studied with the approach developed in the present paper.

## 2. Orlicz-Sobolev Spaces

In this section we recall some results on Orlicz-Sobolev spaces. The results pointed below can be found in $[21,28,36,37]$.

We say that a continuous function $\Phi: \mathbb{R} \rightarrow[0,+\infty)$ is a $N$-function if:
(i) $\Phi$ is convex,
(ii) $\Phi(t)=0 \Leftrightarrow t=0$,
(iii) $\lim _{t \rightarrow 0} \frac{\Phi(t)}{t}=0$ and $\lim _{t \rightarrow+\infty} \frac{\Phi(t)}{t}=+\infty$,
(iv) $\Phi$ is even.

We say that a N -function $\Phi$ verifies the $\Delta_{2}$-condition, and we denote by $\Phi \in \Delta_{2}$, if

$$
\begin{equation*}
\Phi(2 t) \leq K \Phi(t), \quad \forall t \geq t_{0} \tag{2.1}
\end{equation*}
$$

for some constants $K, t_{0}>0$. An important fact is that if $\Phi$ verifies the $\Delta_{2}$-condition then for some $s>0$ there is a constant $C_{s}>0$ such that

$$
\Phi(s t) \leq C_{s} \Phi(t), \quad \forall t \geq 0
$$

We fix an open set $\Omega \subset \mathbb{R}^{N}$ and a $N$-function $\Phi$. We define the Orlicz space associated with $\Phi$ as follows

$$
L^{\Phi}(\Omega)=\left\{u \in L^{1}(\Omega): \quad \int_{\Omega} \Phi\left(\frac{|u|}{\lambda}\right) \mathrm{d} x<+\infty \text { for some } \lambda>0\right\}
$$

The space $L^{\Phi}(\Omega)$ is a Banach space endowed with the Luxemburg norm given by

$$
\|u\|_{\Phi}=\inf \left\{\lambda>0: \int_{\Omega} \Phi\left(\frac{|u|}{\lambda}\right) \mathrm{d} x \leq 1\right\}
$$

In the case of $|\Omega|=+\infty$ we will consider that $\Phi \in \Delta_{2}$ if $t_{0}=0$ in (2.1). The complementary function $\widetilde{\Phi}$ associated with $\Phi$ is given by the Legendre transformation, that is,

$$
\widetilde{\Phi}(s)=\max _{t \geq 0}\{s t-\Phi(t)\}, \quad \text { for all } s \geq 0
$$

We also have a Young-type inequality given by

$$
s t \leq \Phi(t)+\widetilde{\Phi}(s), \quad \text { for all } s, t \geq 0
$$

Using the above inequality, it is possible to establish the following Höldertype inequality:

$$
\left|\int_{\Omega} u v \mathrm{~d} x\right| \leq 2\|u\|_{\Phi}\|v\|_{\tilde{\Phi}}, \quad \text { for all } u \in L^{\Phi}(\Omega), v \in L^{\tilde{\Phi}}(\Omega)
$$

The following results will be often used and they can be found in [21,28].
Lemma 2.1. Consider $\Phi$ a $N$-function of the form (1.4) and satisfying $\left(\phi_{1}\right),\left(\phi_{2}\right)$ and $\left(\phi_{3}\right)$ only with the restriction $1<l \leq m$. Set

$$
\zeta_{0}(t)=\min \left\{t^{\ell}, t^{m}\right\} \text { and } \zeta_{1}(t)=\max \left\{t^{\ell}, t^{m}\right\}, t \geq 0
$$

Then $\Phi$ satisfies

$$
\begin{aligned}
& \zeta_{0}(t) \Phi(\rho) \leq \Phi(\rho t) \leq \zeta_{1}(t) \Phi(\rho), \rho, t>0 \\
& \zeta_{0}\left(\|u\|_{\Phi}\right) \leq \int_{\Omega} \Phi(u) d x \leq \zeta_{1}\left(\|u\|_{\Phi}\right), u \in L_{\Phi}(\Omega)
\end{aligned}
$$

Lemma 2.2. If $\Phi$ is a $N$-function of the form (1.4) satisfying ( $\phi_{1}$ ) and ( $\phi_{2}$ ), then

$$
\widetilde{\Phi}(\phi(|t|) t) \leq \Phi(2 t) \quad \forall t \geq 0
$$

For a $N$-function $\Phi$, the corresponding Orlicz-Sobolev space is defined as the Banach space

$$
W^{1, \Phi}(\Omega)=\left\{u \in L^{\Phi}(\Omega): \quad \frac{\partial u}{\partial x_{i}} \in L^{\Phi}(\Omega), \quad i=1, \ldots, N\right\}
$$

endowed with the norm

$$
\|u\|_{1, \Phi}=\|\nabla u\|_{\Phi}+\|u\|_{\Phi} .
$$

If $\Phi$ and $\widetilde{\Phi}$ satisfy the $\Delta_{2}$-condition, then the spaces $L^{\Phi}(\Omega)$ and $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ are reflexive and separable. Moreover, the $\Delta_{2}$-condition also implies that

$$
u_{n} \rightarrow u \text { in } L_{\Phi}(\Omega) \Longleftrightarrow \int_{\Omega} \Phi\left(\left|u_{n}-u\right|\right) \rightarrow 0
$$

and
$u_{n} \rightarrow u$ in $W^{1, \Phi}(\Omega) \Longleftrightarrow \int_{\Omega} \Phi\left(\left|u_{n}-u\right|\right) \rightarrow 0$ and $\int_{\Omega} \Phi\left(\left|\nabla u_{n}-\nabla u\right|\right) \rightarrow 0$.
An important function related to a $N$-function $\Phi$ is the Sobolev conjugate function $\Phi_{*}$ of $\Phi$ defined by
$\Phi_{*}^{-1}(t)=\int_{0}^{t} \frac{\Phi^{-1}(s)}{s^{(N+1) / N}} \mathrm{~d} s \quad$ for $\quad t>0 \quad$ when $\quad \int_{1}^{+\infty} \frac{\Phi^{-1}(s)}{s^{(N+1) / N}} \mathrm{~d} s=+\infty$.
Let $\Phi_{1}$ and $\Phi_{2}$ be $N$-functions. We say that $\Phi_{1}$ increases strictly lower than $\Phi_{2}$, and we denote by $\Phi_{1} \prec \prec \Phi_{2}$, if $\lim _{t \rightarrow+\infty} \frac{\Phi_{1}(k t)}{\Phi_{2}(t)}=0$ for all $k>0$.

We say that $\Omega \subset \mathbb{R}^{N}$ is an admissible domain, if the embedding $W^{1,1}(\Omega) \hookrightarrow L^{q}(\Omega)$ is continuous for all $q \in\left[1, \frac{N}{N-1}\right]$.

The following embedding results can be found in $[25,38]$.
Theorem 2.1. Let $\Phi$ be a $N$-function, $A$ a $N$-function and $\Omega$ be an admissible domain. Then
(i) If $\limsup _{t \rightarrow 0} \frac{A(t)}{\Phi(t)}<+\infty$ and $\limsup _{t \rightarrow+\infty} \frac{A(t)}{\Phi_{*}(t)}<+\infty$, then the embeddings $W^{1, \Phi}(\Omega) \hookrightarrow L^{A}(\Omega)$ and $W^{1, \Phi}\left(\mathbb{R}^{N}\right) \hookrightarrow L^{A}\left(\mathbb{R}^{N}\right)$ are continuous,
(ii) If $A \prec \prec \Phi_{\star}$ and $|\Omega|<+\infty$, then the embedding $W^{1, \Phi}(\Omega) \hookrightarrow L^{A}(\Omega)$ is compact.

The next result can be found in [25] and it will play an important role in this work.
Theorem 2.2. Consider $\Phi(t):=\int_{0}^{|t|} \phi(t) t d t$ a $N$-function with $\phi$ satisfying $\left(\phi_{1}\right),\left(\phi_{2}\right)$ and $\left(\phi_{3}\right)$. Let $\left(u_{n}\right)$ be a bounded sequence in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ such that there exists $R>0$ satisfying

$$
\lim _{n \rightarrow+\infty} \sup _{y \in \mathbb{R}^{N}} \int_{B_{R}(y)} \Phi\left(\left|u_{n}\right|\right)=0
$$

Then, for any $N$-function $P$ verifying the $\Delta_{2}$-condition with

$$
\begin{equation*}
\lim _{t \rightarrow 0} \frac{P(t)}{\Phi(t)}=0 \tag{1}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{t \rightarrow+\infty} \frac{P(t)}{\Phi_{*}(t)}=0 \tag{2}
\end{equation*}
$$

we have

$$
u_{n} \rightarrow 0 \quad \text { in } \quad L^{P}\left(\mathbb{R}^{N}\right)
$$

## 3. Differentiability of the Functional $\Psi$

In this section, we will study the differentiability of the functional $\Psi$ given in (1.5). To this end, we must assume some conditions on $f$.

We will consider $B: \mathbb{R} \rightarrow[0,+\infty)$ being a $N$-function given by $B(t)=\int_{0}^{|t|} b(s) s \mathrm{~d} s$, where $b:(0,+\infty) \rightarrow(0,+\infty)$ is a function satisfying the following conditions:
$\left(H_{2}\right)$

$$
\begin{array}{r}
(t b(t))^{\prime}>0, \quad \text { for all } t>0,  \tag{1}\\
\lim _{t \rightarrow+\infty} t b(t)=+\infty
\end{array}
$$

and there exist $b_{i} \in(1,+\infty), i=1,2$ such that

$$
\begin{equation*}
b_{1} \leq \frac{b(t) t^{2}}{B(t)} \leq b_{2}, \forall t>0 \tag{3}
\end{equation*}
$$

The above hypotheses permit to use Lemmas 2.1 and 2.2 by changing $\Phi$ by $B$. More precisely, $B$ satisfies the $\Delta_{2}$ condition with

$$
\zeta_{0, B}(t) B(\rho) \leq B(\rho t) \leq \zeta_{1, B}(t) B(\rho), \rho, t>0,
$$

and

$$
\zeta_{0, B}\left(\|u\|_{B}\right) \leq \int_{\Omega} B(u) \mathrm{d} x \leq \zeta_{1, B}\left(\|u\|_{B}\right), u \in L_{B}(\Omega)
$$

where

$$
\zeta_{0, B}(t)=\min \left\{t^{b_{1}}, t^{b_{2}}\right\} \text { and } \zeta_{1, B}(t)=\max \left\{t^{b_{1}}, t^{b_{2}}\right\}, t \geq 0 .
$$

Let $f: \mathbb{R} \rightarrow \mathbb{R}$ be a continuous function satisfying the growth condition

$$
\begin{equation*}
|f(t)| \leq C b(|t|)|t|, \quad \forall t \in \mathbb{R} \tag{1}
\end{equation*}
$$

where $C$ is a positive constant.
The primitive of $f$, that is, $F(t):=\int_{0}^{t} f(s) \mathrm{d} s$ is continuous and satisfies

$$
\begin{equation*}
|F(t)| \leq C B(t), \quad \forall t \in \mathbb{R} \tag{F}
\end{equation*}
$$

for some constant $C>0$.
Let $0<\lambda<N$ and consider $s \in \mathbb{R}$ with

$$
\frac{1}{s}+\frac{\lambda}{N}+\frac{1}{s}=2
$$

that is, $s=\frac{2 N}{2 N-\lambda}$. We will also suppose that the embeddings

$$
\begin{equation*}
W^{1, \Phi}\left(\mathbb{R}^{N}\right) \hookrightarrow L^{s b_{i}}\left(\mathbb{R}^{N}\right), i=1,2 \tag{E}
\end{equation*}
$$

are continuous. The above condition is not empty, because we can consider the following conditions on $B$

$$
\limsup _{t \rightarrow 0^{+}} \frac{t^{s b_{i}}}{\Phi(t)}=0
$$

and

$$
\limsup _{t \rightarrow+\infty} \frac{t^{s b_{i}}}{\Phi_{*}(t)}=0
$$

If $\Phi(t)=\frac{1}{p}|t|^{p}$ with $1<p<N$, we have that $\Delta_{\Phi} u=\Delta_{p} u$ and the function $f$ can be of the form

$$
f(t)=|t|^{q-2} t+|t|^{\beta-2} t, \quad \text { for all } t \in \mathbb{R}
$$

with $s q, s \beta \in\left(p, p^{*}\right)$.
In the proof of the differentiability of $\Psi$, we will use the following elementary property, whose proof we omit.

Lemma 3.1. Let $E$ be a normed vector space and $J: E \rightarrow \mathbb{R}$ be a functional verifying the following properties:
(i) the Fréchet derivative $\frac{\partial J(u)}{\partial v}:=\lim _{t \rightarrow 0} \frac{J(u+t v)-J(u)}{t}$ exists for all $u, v \in E$;
(ii) for each $u \in E, \frac{\partial J(u)}{\partial(.)} \in E^{\prime}$, that is, the application $v \mapsto \frac{\partial J(u)}{\partial v}$ is a continuous linear functional;
(iii) we have

$$
u_{n} \rightarrow u \quad \text { in } \quad E \Longrightarrow \frac{\partial J\left(u_{n}\right)}{\partial(.)} \rightarrow \frac{\partial J(u)}{\partial(.)} \quad \text { in } \quad E^{\prime}
$$

that is,

$$
u_{n} \rightarrow u \quad \text { in } \quad E \Longrightarrow \sup _{\|v\| \leq 1}\left|\frac{\partial J\left(u_{n}\right)}{\partial v}-\frac{\partial J(u)}{\partial v}\right| \rightarrow 0
$$

Then $J \in C^{1}(E, \mathbb{R})$ and

$$
J^{\prime}(u) v=\frac{\partial J(u)}{\partial v}, \quad \text { for all } u, v \in E
$$

We are now ready to prove the differentiability of functional $\Psi$ given by (1.5).

Lemma 3.2. Assume $(E),\left(H_{1}\right)-\left(H_{3}\right)$ and $\left(f_{1}\right)$. Then $\Psi$ given in (1.5) is well defined and belongs to $C^{1}\left(W^{1, \Phi}\left(\mathbb{R}^{N}\right), \mathbb{R}\right)$ with

$$
\Psi^{\prime}(u) v=\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) f(u(y)) v(y))}{|x-y|^{\lambda}} d x d y
$$

for all $u, v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.
Proof. Using the definition of $s$, the condition $(F)$ and Proposition 1.1, it follows that $\Psi$ is well defined. In the sequel, we will show that $\Psi$ satisfies the assumptions of Lemma 3.1. To this end, we will divide the proof into three steps.

Step 1: Existence of the Fréchet derivative:
Let $u, v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ and $t \in[-1,1]$. Note that

$$
\begin{align*}
& \frac{\Psi(u+t v)-\Psi(u)}{t} \\
& \quad=\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)+t v(x)) F(u(y)+t v(y))-F(u(x)) F(u(y))}{t|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y . \tag{3.1}
\end{align*}
$$

Denoting by $I$ the integrand in (3.1), we have

$$
\begin{aligned}
I= & \frac{F(u(x)+t v(x))(F(u(y)+t v(y))-F(u(y))}{t} \\
& +\frac{F(u(y))(F(u(x)+t v(x))-F(u(x)))}{t} .
\end{aligned}
$$

By the mean value theorem, there exist $\theta(x, t), \eta(y, t) \in[-1,1]$, such that

$$
F(u(y)+t v(y))-F(u(y))=f(u(y)+\eta(y, t) t v(y)) v(y) t
$$

and

$$
F(u(x)+t v(x))-F(u(x))=f(u(x)+\theta(x, t) t v(x)) v(x) t .
$$

The relation (3.1) allows us to estimate

$$
\begin{aligned}
& \left|\frac{\Psi(u+t v)-\Psi(u)}{t}-\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x) f(u(y)) v(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y\right| \\
& \quad \leq\left|B_{1}^{t}\right|+\left|B_{2}^{t}\right|,
\end{aligned}
$$

where

$$
B_{1}^{t}:=\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)+t v(x)) f(u(y)+\eta(y, t) t v(y)) v(y)-F(u(x) f(u(y)) v(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y
$$

and

$$
\begin{aligned}
B_{2}^{t}:= & \frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(y)) f(u(x)+\theta(x, t) t v(x)) v(x)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \\
& -\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) f(u(y)) v(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y .
\end{aligned}
$$

If $x \in \mathbb{R}^{N}$ and $u(x) \neq 0$, Lemma 2.1 gives

$$
\begin{aligned}
|f(u(x)) v(x)| & \leq C b(|u(x)|)|u(x)||v(x)| \\
& \leq C\left(\frac{B(|u(x)|)}{|u(x)|}\right)|v(x)| \\
& \leq C\left(|u(x)|^{b_{1}-1}+|u(x)|^{b_{2}-1}\right)|v(x)| .
\end{aligned}
$$

The above estimate also holds if $u(x)=0$. Such estimate, combined with $(E)$ implies that $f(u) v \in L^{s}\left(\mathbb{R}^{N}\right)$. Since $F(u) \in L^{s}\left(\mathbb{R}^{N}\right)$, we have by Proposition 1.1,

$$
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|F(u(x)) f(u(y)) v(y)|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y<+\infty .
$$

Thus, by Fubini's theorem,

$$
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) f(u(y)) v(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y=\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(y)) f(u(x)) v(x)}{|x-y|^{\lambda}} \mathrm{d} y \mathrm{~d} x .
$$

Therefore, $B_{2}^{t}$ can be rewritten as

$$
B_{2}^{t}:=\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(y)) f(u(x)+\theta(x, t) t v(x)) v(x)-F(u(y) f(u(x)) v(x))}{|x-y|} \mathrm{d} x \mathrm{~d} y .
$$

By Proposition 1.1 we obtain

$$
\left|B_{2}^{t}\right| \leq C\|F(u)\|_{L^{s+}\left(\mathbb{R}^{N}\right)}\|f(u+\theta(., t) v) v-f(u) v\|_{L^{s+}\left(\mathbb{R}^{N}\right)} .
$$

Since $\theta(x, t) \in[-1,1]$, the condition $(f)$ combined with Lemma 2.1 implies that

$$
\begin{align*}
|f(u(x)+\theta(t, x) t v(x)) v(x)-f(u(x)) v(x)|^{s} \leq & C\left((|u(x)|+|v(x)|)^{s\left(b_{1}-1\right)}|v(x)|^{s}\right. \\
& \left.+(|u(x)|+|v(x)|)^{s\left(b_{2}-1\right)}|v(x)|^{s}\right), \tag{3.2}
\end{align*}
$$

for all $x \in \mathbb{R}^{N}$. Here $C$ is a constant that does not depend on $t \in[-1,1]$. The embeddings $(E)$ ensure that the right-hand side of the inequality (3.2) is an integrable function. Thus, the Lebesgue's dominated convergence theorem yields

$$
\|f(u+\theta(., t) v) v-f(u) v\|_{L^{s}\left(\mathbb{R}^{N}\right)} \rightarrow 0 \text { as } t \rightarrow 0
$$

The last limit implies that $B_{2}^{t} \rightarrow 0$ as $t \rightarrow 0$. With respect to $B_{1}^{t}$, we have the estimate below

$$
\begin{aligned}
\left|B_{1}^{t}\right| \leq & \frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|F(u(x))||f(u(y)+\eta(y, t) t v(y)) v(y)-f(v(y)) v(y)|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \\
& +\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|f(u(y)+\eta(y, t) t v(y)) v(y)||F(u(x)+t v(x))-F(u(x))|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y .
\end{aligned}
$$

Arguing as before,

$$
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|F(u(x))||f(u(y)+\eta(y, t) t v(y)) v(y)-f(v(y)) v(y)|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \rightarrow 0
$$

as $t \rightarrow 0$. On the other hand, the Lebesgue's dominated convergence theorem also yields

$$
\begin{equation*}
\|F(u+t v)-F(u)\|_{L^{s}\left(\mathbb{R}^{N}\right)} \rightarrow 0 \text { as } t \rightarrow 0 \tag{3.3}
\end{equation*}
$$

As in (3.2), $\|f(u+\eta(., t) t v) v\|_{L^{s}\left(\mathbb{R}^{N}\right)}$ is uniformly bounded by a constant that does not depend on $t \in[-1,1]$. Thus, Proposition 1.1 combined with (3.3) gives

$$
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|f(u(y)+\eta(y, t) t v(y)) v(y)||F(u(x)+t v(x))-F(u(x))|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \rightarrow 0
$$

as $t \rightarrow 0$, and so, $B_{1}^{t} \rightarrow 0$ as $t \rightarrow 0$. From the above analysis,

$$
\lim _{t \rightarrow 0} \frac{\Psi(u+t v)-\Psi(u)}{t}=\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x) f(u(y)) v(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y
$$

showing the existence of the Fréchet derivative $\frac{\partial \Psi(u)}{\partial v}$.

Step 2: $\frac{\partial \Psi(u)}{\partial(.)} \in\left(W^{1, \Phi}\left(\mathbb{R}^{N}\right)\right)^{\prime}$ for all $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.
It is evident that $\frac{\partial \Psi(u)}{\partial v}$ is linear at $v$ for each $u$ fixed. Next, we are going to show that

$$
\left|\frac{\partial \Psi(u)}{\partial v}\right| \leq C_{u}, \quad \forall v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \quad \text { with } \quad\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1
$$

for some positive constant $C_{u}$. From $(f),(F)$ and Proposition 1.1 we have

$$
\begin{equation*}
\left|\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) f(u(y)) v(y)}{|x-y|^{\lambda(x, y)}} \mathrm{d} x \mathrm{~d} y\right| \leq C\|F(u)\|_{L^{s}\left(\mathbb{R}^{N}\right)}\|f(u) v\|_{L^{s}\left(\mathbb{R}^{N}\right)} \tag{3.4}
\end{equation*}
$$

The continuous embeddings of $(E)$ combined with Hölder inequality, $(f)$ and Proposition 2.1 give

$$
\begin{align*}
& \int_{\mathbb{R}^{N}}|f(u(y)) v(y)|^{s} \mathrm{~d} y \leq C\left(\int_{\mathbb{R}^{N}}|u(y)|^{s\left(b_{1}-1\right)}|v(y)|^{s} \mathrm{~d} y\right. \\
& \left.\quad+\int_{\mathbb{R}^{N}}|u(y)|^{s\left(b_{2}-1\right)}|v(y)|^{s} \mathrm{~d} y\right) \\
& \quad+C\left(\left(\int_{\mathbb{R}^{N}}|u(y)|^{s b_{1}} \mathrm{~d} y\right)^{\frac{b_{1}-1}{b_{1}}}\left(\int_{\mathbb{R}^{N}}|v(y)|^{s b_{1}}\right)^{\frac{1}{b_{1}}}\right. \\
& \left.\quad+\left(\int_{\mathbb{R}^{N}}|u(y)|^{s b_{2}} \mathrm{~d} y\right)^{\frac{b_{2}-1}{b_{2}}}\left(\int_{\mathbb{R}^{N}} \mid v(y)^{s b_{2}}\right)^{\frac{1}{b_{2}}}\right) \leq C_{u}, \tag{3.5}
\end{align*}
$$

where

$$
C_{u}=K\left(\max \left(\left(\int_{\mathbb{R}^{N}}|u(y)|^{s b_{1}} \mathrm{~d} y\right)^{\frac{b_{1}-1}{b_{1}}},\left(\int_{\mathbb{R}^{N}}|u(y)|^{s b_{2}} \mathrm{~d} y\right)^{\frac{b_{2}-1}{b_{2}}}\right)\right)
$$

with $K$ a constant that does not depend on $u$ and $v$. Then inequalities (3.4) and (3.5) justify Step 2.

## Step 3:

$$
u_{n} \rightarrow u \quad \text { in } \quad W^{1, \Phi}\left(\mathbb{R}^{N}\right) \Rightarrow \sup _{\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1}\left|\frac{\partial \Psi\left(u_{n}\right)}{\partial v}-\frac{\partial \Psi(u)}{\partial v}\right| \rightarrow 0 .
$$

Consider $v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with $\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1$ and note that

$$
\begin{aligned}
\left|\frac{\partial \Psi\left(u_{n}\right)}{\partial v}-\frac{\partial \Psi(u)}{\partial v}\right| \leq & \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{\left|F\left(u_{n}(x)\right)-F(u(x))\right|\left|f\left(u_{n}(y)\right) v(y)\right|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \\
& +\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|F(u(x))|\left|f\left(u_{n}(y)\right) v(y)-f(y, u(y)) v(y)\right|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \\
:= & B_{f}^{n}+B_{F}^{n} .
\end{aligned}
$$

By Proposition 1.1, we obtain

$$
\left.B_{f}^{n} \leq C \| F\left(u_{n}\right)-F(u)\right)\left\|_{L^{s}\left(\mathbb{R}^{N}\right)}\right\| f\left(u_{n}\right) v \|_{L^{s}\left(\mathbb{R}^{N}\right)}
$$

Since the sequence $\left(\left\|f\left(u_{n}\right) v\right\|_{L^{s}\left(\mathbb{R}^{N}\right)}\right)$ is bounded (see (3.5)) and

$$
\left\|F\left(u_{n}\right)-F(u)\right\|_{L^{s}\left(\mathbb{R}^{N}\right)} \rightarrow 0
$$

it follows that

$$
\begin{equation*}
\left.\sup _{\substack{v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \\\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1}} \| F\left(u_{n}\right)-F(u)\right)\left\|_{L^{s}\left(\mathbb{R}^{N}\right)}\right\| f\left(u_{n}\right) v \|_{L^{s}\left(\mathbb{R}^{N}\right)} \rightarrow 0 \quad \text { as } \quad n \rightarrow+\infty \tag{3.6}
\end{equation*}
$$

Now we will estimate $B_{F}^{n}$. Given $\varepsilon>0$, fix $R>0$ large enough such that

$$
\int_{B(0, R)^{c}}|u(x)|^{b_{1} s} \mathrm{~d} x, \int_{B(0, R)^{c}}|u(x)|^{b_{2} s} \mathrm{~d} x<\varepsilon
$$

Since $u_{n} \rightarrow u$ in $L^{b_{1} s}\left(\mathbb{R}^{N}\right)$ and $L^{b_{2} s}\left(\mathbb{R}^{N}\right)$, there is $n_{0} \in \mathbb{N}$ large enough such that

$$
\begin{equation*}
\int_{B(0, R)^{c}}\left|u_{n}(x)\right|^{b_{1} s} \mathrm{~d} x, \int_{B(0, R)^{c}}\left|u_{n}(x)\right|^{b_{2} s} \mathrm{~d} x<\varepsilon \tag{3.7}
\end{equation*}
$$

for all $n \geq n_{0}$. Note that by Proposition 1.1

$$
\begin{align*}
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|F(u(x))|\left|\left(f\left(u_{n}(y)\right)-f(u(y))\right) v(y)\right|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \leq C\|F(u)\|_{L^{s}\left(\mathbb{R}^{N}\right)}  \tag{3.8}\\
\times\left\|\left(f\left(u_{n}\right)-f(u)\right) v\right\|_{L^{s}\left(\mathbb{R}^{N}\right)}
\end{align*}
$$

where $C_{0}>0$ is a constant that does not depend on $n \in \mathbb{N}$. The condition $(f)$ together with Hölder's inequality yields

$$
\begin{align*}
& \int_{B(0, R)^{c}}\left|\left(f\left(u_{n}(y)\right)-f(u(x))\right) v(y)\right|^{s} \mathrm{~d} y \\
& \leq C_{1}\left\|u_{n}\right\|_{L^{s b_{1}\left(B(0, R)^{c}\right)}}^{s\left(b_{1}-1\right)}\|v\|_{L^{s b_{1}}\left(B(0, R)^{c}\right)}^{s}  \tag{3.9}\\
& \quad+C_{1}\left\|u_{n}\right\|_{L^{s b_{2}}\left(B(0, R)^{c}\right)}^{s\left(b_{2}-1\right)}\|v\|_{L^{s b_{2}}\left(B(0, R)^{c}\right)}^{s}
\end{align*}
$$

where $C_{1}>0$ is a constant that does not depend on $n \in \mathbb{N}$. Using the embeddings $(E)$, we have $\|v\|_{L^{s b_{i}}\left(B(0, R)^{c}\right)}^{s} \leq C_{2}, i=1,2$ where $C_{2}$ is a positive constant that does not depend on $v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with $\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1$ and $R>0$. Thus, from (3.7) and (3.9)

$$
\int_{B(0, R)^{c}}\left|\left(f\left(u_{n}(y)\right)-f(u(x))\right) v(y)\right|^{+} \mathrm{d} y \leq C_{3} \max \left\{\varepsilon^{\frac{b_{1}-1}{b_{1}}}, \varepsilon^{\frac{b_{2}-1}{b_{2}}}\right\}
$$

where $C_{3}$ is a positive constant that does not depend on $v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with $\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1$ and $R>0$. Therefore

$$
\begin{equation*}
\sup _{\substack{v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \\\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1}} \int_{B(0, R)^{c}}\left|\left(f\left(u_{n}(y)\right)-f(u(y))\right) v(y)\right|^{s} \mathrm{~d} y \leq A_{\varepsilon}, \quad \forall n \geq n_{0} \tag{3.10}
\end{equation*}
$$

where

$$
A_{\varepsilon}=C_{3} \max \left\{\varepsilon^{\frac{b_{1}-1}{b_{1}}}, \varepsilon^{\frac{b_{2}-1}{b_{2}}}\right\} \rightarrow 0 \quad \text { as } \quad \varepsilon \rightarrow 0^{+}
$$

Now we will estimate the integral

$$
\int_{B(0, R)}\left|f\left(u_{n}(y)\right)-f(u(y))\right|^{s}|v(y)|^{s} \mathrm{~d} y
$$

From $(f)$ and Lemma 2.1,

$$
\left|f\left(u_{n}(y)\right)-f(u(y))\right|^{s} \leq C\left(1+\left|u_{n}(y)\right|^{s\left(b_{2}-1\right)}+|u(y)|^{s\left(b_{2}-1\right)}\right), \forall y \in \mathbb{R}^{N}
$$

From $(E)$ we have that the embedding $W^{1, \Phi}\left(\mathbb{R}^{N}\right) \hookrightarrow L^{s b_{2}}\left(\mathbb{R}^{N}\right)$ is continuous which implies that $\left|f\left(u_{n}\right)-f(u)\right|^{\frac{s b_{2}}{b_{2}-1}},|v|^{s b_{2}} \in L^{1}(B(0, R))$.

From Hölder's inequality we get

$$
\begin{aligned}
\int_{B(0, R)}\left|f\left(u_{n}(y)\right)-f(u(y))\right|^{s}|v(y)|^{s} \mathrm{~d} y \leq & C\left\|\left|f\left(u_{n}\right)-f(u)\right|^{s}\right\|_{L^{\frac{b_{2}}{b_{2}-1}}(B(0, R))} \\
& \times\left\||v|^{s}\right\|_{L^{b_{2}}(B(0, R))} \\
\leq & C_{4}\left\|\left|f\left(u_{n}\right)-f(u)\right|^{s}\right\|_{L^{\frac{b_{2}}{b_{2}-1}}(B(0, R))}
\end{aligned}
$$

where $C_{4}$ is a positive constant that does not depend on $n \in \mathbb{N}$ and $v \in$ $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with $\|v\|_{W^{1, \Phi}\left(\mathbb{R}^{N}\right)} \leq 1$. Recalling that $\left|f\left(u_{n}(y)\right)-f(u(y))\right|^{s} \rightarrow 0$ in $L^{\frac{b_{2}}{b_{2}-1}}(B(0, R))$, it follows that

$$
\begin{equation*}
\sup _{\substack{v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \\\|v\|_{1, \Phi} \leq 1}} \int_{B(0, R)}\left|\left(f\left(u_{n}(y)\right)-f(u(y))\right) v(y)\right|^{s} \mathrm{~d} y \rightarrow 0 \text { as } n \rightarrow+\infty \tag{3.11}
\end{equation*}
$$

From (3.10) and (3.11),

$$
\begin{equation*}
\sup _{\substack{v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \\\|v\|_{1, \Phi} \leq 1}}\left\|\left(f\left(u_{n}\right)-f(u)\right) v\right\|_{L^{s}\left(\mathbb{R}^{N}\right)} \rightarrow 0 \text { as } n \rightarrow+\infty \tag{3.12}
\end{equation*}
$$

From (3.8) and (3.12),

$$
\begin{equation*}
\sup _{\substack{v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \\\|v\|_{1, \Phi} \leq 1}} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{|F(u(x))|\left|\left(f\left(u_{n}(y)\right)-f(y, u(y))\right) v(y)\right|}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \rightarrow 0 \quad \text { as } \quad n \rightarrow+\infty \tag{3.13}
\end{equation*}
$$

The step is justified, according to (3.6), and (3.13). Finally, the lemma follows from the previous three steps.

## 4. Existence of Nontrivial Solutions

To prove the existence of solution for (1.3) we will need some assumptions. In what follows, we will consider

$$
m<s b_{i}<l^{*}, i=1,2
$$

which implies that the embeddings in $(E)$ hold. Moreover, we will consider the condition $\left(f_{1}\right)$ with

$$
\begin{equation*}
b_{i}>m / 2, i=1,2 \tag{4.1}
\end{equation*}
$$

Finally, we also assume the Ambrosetti-Rabinowitz-type condition: there is $\theta>m$ such that

$$
\begin{equation*}
0<\theta F(t) \leq 2 f(t) t, \quad \forall t \neq 0 \tag{2}
\end{equation*}
$$

Our main result is the following:
Theorem 4.1. The problem (1.3) has a nontrivial solution under the conditions (4.1), $\left(\phi_{1}\right)-\left(\phi_{3}\right),\left(H_{1}\right)-\left(H_{3}\right)$ and $\left(f_{1}\right)-\left(f_{2}\right)$.

In the proof of Theorem 4.1 we will use variational methods. The energy functional $J: W^{1, \Phi}\left(\mathbb{R}^{N}\right) \rightarrow \mathbb{R}$ associated with (1.3) is given by,

$$
\begin{aligned}
J(u)= & \int_{\mathbb{R}^{N}}(\Phi(|\nabla u(x)|)+\Phi(|u(x)|)) \mathrm{d} x \\
& -\frac{1}{2} \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) F(u(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y
\end{aligned}
$$

that is,

$$
J(u)=\int_{\mathbb{R}^{N}}(\Phi(|\nabla u(x)|)+\Phi(|u(x)|)) \mathrm{d} x-\Psi(u)
$$

The analysis developed in the previous section implies that $J \in$ $C^{1}\left(W^{1, \Phi}\left(\mathbb{R}^{N}\right), \mathbb{R}\right)$ with

$$
\begin{aligned}
J^{\prime}(u) v=\int_{\mathbb{R}^{N}} \phi(|\nabla u(x)|) \nabla u(x) \nabla v(x) \mathrm{d} x & +\int_{\mathbb{R}^{N}} \phi(|u(x)|) u(x) v(x) \mathrm{d} x \\
& -\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x) f(u(y)) v(y))}{|x-y|^{\lambda(x, y)}} \mathrm{d} x \mathrm{~d} y
\end{aligned}
$$

for all $u, v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.
Our first lemma establishes the mountain pass geometry.
Lemma 4.1. The functional J verifies the following properties:
(i) There exists $\rho>0$ small enough such that $J(u) \geq \eta$ for $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with $\|u\|=\rho$, for some $\eta>0$.
(ii) There exists $e \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ such that $\|e\|>\rho$ and $J(e)<0$.

Proof. (i) By $(F)$ and Proposition 1.1 we have

$$
\left|\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) F(u(y))}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y\right| \leq C\|F(u)\|_{L^{s}\left(\mathbb{R}^{N}\right)}^{2}
$$

for all $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$. Note that

$$
\begin{aligned}
\|F(u)\|_{L^{s}\left(\mathbb{R}^{N}\right)} & \leq C\left(\int_{\mathbb{R}^{N}}\left(|u(x)|^{s b_{1}}+|u(x)|^{s b_{2}}\right) \mathrm{d} x\right)^{\frac{1}{s}} \\
& \leq C\left(\|u\|_{L^{s b_{1}}\left(\mathbb{R}^{N}\right)}^{b_{1}}+\|u\|_{L^{s b_{2}}\left(\mathbb{R}^{N}\right)}^{b_{2}}\right)
\end{aligned}
$$

where $C$ is a constant that does not depend on $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.
From $(E)$ and Theorem 2.1, it follows that the embeddings $W^{1, \Phi}\left(\mathbb{R}^{N}\right) \hookrightarrow$ $L^{s b_{i}}\left(\mathbb{R}^{N}\right), i=1,2$ are continuous. Therefore

$$
\|u\|_{L^{s b_{i}\left(\mathbb{R}^{N}\right)}} \leq L\|u\|_{1, \Phi}, u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right), i=1,2
$$

for a positive constant $L>0$ that does not depend on $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.

By using the classical inequality

$$
(x+y)^{\alpha} \leq 2^{\alpha-1}\left(x^{\alpha}+y^{\alpha}\right), x, y \geq 0 \quad \text { with } \quad \alpha>1
$$

and Lemma 2.1, we get for $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with $\|u\|_{1, \Phi}=\|\nabla u\|_{\Phi}+\|u\|_{\Phi}<1$ that

$$
\begin{aligned}
J(u) \geq & \Phi(1)\left(\min \left(\|\nabla u\|_{\Phi}^{l},\|\nabla u\|_{\Phi}^{m}\right)+\min \left(\|u\|_{\Phi}^{l},\|u\|_{\Phi}^{m}\right)\right) \\
& -C\left(\|u\|_{1, \Phi}^{2 b_{1}}+\|u\|_{1, \Phi}^{2 b_{2}}\right) \\
\geq & \bar{C}\left(\|\nabla u\|_{\Phi}^{m}+\|u\|_{\Phi}^{m}\right)-\bar{C}\left(\|u\|_{1, \Phi}^{2 b_{1}}+\|u\|_{1, \Phi}^{2 b_{2}}\right) \\
\geq & K\|u\|_{1, \Phi}^{m}-K\left(\|u\|_{1, \Phi}^{2 b_{1}}+\|u\|_{1, \Phi}^{2 b_{2}}\right)
\end{aligned}
$$

where $C, \bar{C}$, and $K$ are constants that do not depend on $u$. Since (4.1) holds, the result follows by fixing $\|u\|_{1, \Phi}=\rho$ with $\rho>0$ small enough.
(ii) The condition $\left(f_{2}\right)$ implies that

$$
F(t) \geq C_{1} t^{\frac{\theta}{2}}-C_{2} \quad \forall t \in \mathbb{R},
$$

where $C_{1}, C_{2} \geq 0$ depends only on $l$ and $\theta$. Now, considering a nonnegative function $\varphi \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right) \backslash\{0\}$, the last inequality permits to conclude that $J(t \varphi)<0$ for $t$ large enough. This finishes the proof.

Using the mountain pass theorem without the Palais-Smale condition (see [39, Theorem 5.4.1]), there is a sequence $\left(u_{n}\right) \subset W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ such that

$$
J\left(u_{n}\right) \rightarrow d \quad \text { and } \quad J^{\prime}\left(u_{n}\right) \rightarrow 0
$$

where $d>0$ is the mountain pass level defined by

$$
\begin{equation*}
d:=\inf _{\gamma \in \Gamma} \sup _{t \in[0,1]} J(\gamma(t)), \tag{4.2}
\end{equation*}
$$

with

$$
\Gamma:=\left\{\gamma \in C\left(W^{1, \Phi}\left(\mathbb{R}^{N}\right), W^{1, \Phi}\left(\mathbb{R}^{N}\right) ; \gamma(0)=0, \gamma(1)=e\right\} .\right.
$$

Regarding the above sequence we have the following auxiliary property.
Lemma 4.2. The sequence $\left(u_{n}\right)$ is bounded in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.
Proof. Note that

$$
J\left(u_{n}\right)-\frac{J^{\prime}\left(u_{n}\right) u_{n}}{\theta} \leq d+1+\left\|u_{n}\right\|_{1, \Phi}
$$

for $n$ large enough and $d$ given in (4.2). On the other hand by $\left(\phi_{3}\right)$ and $\left(f_{2}\right)$ we have

$$
\begin{aligned}
J\left(u_{n}\right)-\frac{J^{\prime}\left(u_{n}\right) u_{n}}{\theta}= & \int_{\mathbb{R}^{N}}\left(\Phi\left(\left|\nabla u_{n}(x)\right|\right)+\Phi\left(\left|u_{n}(x)\right|\right) \mathrm{d} x\right. \\
& -\frac{1}{\theta} \int_{\mathbb{R}^{N}} \phi(|\nabla u|)|\nabla u|^{2}+\phi\left(\left|u_{n}(x)\right|\right)\left|u_{n}(x)\right|^{2} \mathrm{~d} x \\
& +\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(x, u_{n}(x)\right)}{|x-y|^{\lambda}}\left(\frac{f\left(y, u_{n}(y)\right) u_{n}(y)}{\theta}-\frac{F\left(y, u_{n}(y)\right)}{2}\right) \mathrm{d} x \mathrm{~d} y \\
\geq & \left(1-\frac{m}{\theta}\right) \int_{\mathbb{R}^{N}}\left(\Phi\left(\left|\nabla u_{n}\right|\right)+\Phi\left(\left|u_{n}\right|\right)\right) \mathrm{d} x .
\end{aligned}
$$

The last two inequalities give the boundedness of $\left(u_{n}\right)$ in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.

Since $\left(u_{n}\right)$ is bounded in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$, there exists $u \in W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ such that $\nabla u_{n} \rightharpoonup \nabla u$ in $\left(L^{\Phi}\left(\mathbb{R}^{N}\right)\right)^{N}$ up to a subsequence. From (4.1) and Theorem 2.1, we have that the embedding $W^{1, \Phi}\left(B_{R}(0)\right) \hookrightarrow L^{\Phi}\left(B_{R}(0)\right)$, where $B_{R}(0)$ denotes the open ball centered at the origin with radius $R$ is compact.

Since $L^{\Phi}\left(B_{R}(0)\right) \hookrightarrow L^{1}\left(B_{R}(0)\right)$, it follows that $u_{n}(x) \rightarrow u(x)$ a.e in $B_{R}(0)$ for some subsequence. Since $R>0$ is arbitrary, we have that $u_{n}(x) \rightarrow$ $u(x)$ a.e in $\mathbb{R}^{N}$ for some subsequence.

The next two lemmas will be needed to prove that $u$ is a critical point of $J$.

Lemma 4.3. The following limits hold for a subsequence:
(i)

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(x, u_{n}(x)\right) f(y, u(y)) v(y)}{|x-y|^{\lambda}} d x d y \\
& \quad \rightarrow \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(x, u(x)) f(y, u(y)) v(y)}{|x-y|^{\lambda}} d x d y
\end{aligned}
$$

for all $v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$,
(ii)

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(x, u_{n}(x)\right)\left(f\left(y, u_{n}(y)\right) v(y)-f(y, u(y)) v(y)\right)}{|x-y|^{\lambda}} d x d y \rightarrow 0 \\
& \text { for all } v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)
\end{aligned}
$$

(iii)

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(x, u_{n}(x)\right) f\left(y, u_{n}(y)\right) v(y)}{|x-y|^{\lambda}} d x d y \\
& \quad \rightarrow \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(x, u(x)) f(y, u(y)) v(y)}{|x-y|^{\lambda}} d x d y
\end{aligned}
$$

for all $v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$.
Proof. (i) The hypothesis $(F)$, the fact that $\left(u_{n}\right)$ is bounded in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ and the continuous embeddings $W^{1, \Phi}\left(\mathbb{R}^{N}\right) \hookrightarrow L^{s b_{i}}\left(\mathbb{R}^{N}\right), i=1,2$ ensure that $\left(F\left(u_{n}().\right)\right)$ is a bounded sequence in $L^{s}\left(\mathbb{R}^{N}\right)$. Combining the previous information with the pointwise convergence $F\left(u_{n}(x)\right) \rightarrow F(u(x))$ a.e in $\mathbb{R}^{N}$, we have $F\left(u_{n}().\right) \rightharpoonup F(u)$ in $L^{s}\left(\mathbb{R}^{N}\right)$.

By Proposition 1.1 it follows that the function

$$
H(w):=\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{w(x) f(u(y)) v(y)}{|x-y|^{\lambda}}, w \in L^{s}\left(\mathbb{R}^{N}\right)
$$

defines a continuous linear functional. Since $F\left(., u_{n}\right) \rightharpoonup F(., u)$ in $L^{s}\left(\mathbb{R}^{N}\right)$, it follows that

$$
\begin{aligned}
& \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right) v(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \\
& \quad \rightarrow \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) f(u(y)) v(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y
\end{aligned}
$$

which proves $(i)$.
(ii) Denote by $I$ the integral described in (ii). Since $\left(F\left(u_{n}().\right)\right)$ is bounded in $L^{s}\left(\mathbb{R}^{N}\right)$,

$$
\begin{aligned}
|I| & \leq C\left\|F\left(u_{n}\right)\right\|_{L^{s}\left(\mathbb{R}^{N}\right)}\left\|f\left(u_{n}\right) v-f(u) v\right\|_{L^{s}\left(\mathbb{R}^{N}\right)} \\
& \leq K\left\|f\left(u_{n}\right) v-f(u) v\right\|_{L^{s}\left(\mathbb{R}^{N}\right)},
\end{aligned}
$$

for some positive constant $K>0$ that does not depend on $n \in \mathbb{N}$ and $v \in$ $C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$. Let $v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ and consider a bounded open set $\Omega$ that contains the support of $v$. Since $\Omega$ is bounded the compactness of the embeddings $W^{1, \Phi}(\Omega) \hookrightarrow L^{s b_{i}}(\Omega), i=1,2$ implies that there exist $\bar{c} \in L^{s b_{1}}(\Omega)$ and $\bar{d} \in$ $L^{s b_{2}}(\Omega)$ such that

- $u_{n}(x) \rightarrow u(x)$ a.e in $\mathbb{R}^{N}$,
- $\left|u_{n}(x)\right| \leq \bar{c}(x), \bar{d}(x)$ a.e in $\Omega$.

These information combined with Lebesgue's dominated convergence theorem give

$$
\left\|f\left(u_{n}\right) v-f(u) v\right\|_{L^{s}\left(\mathbb{R}^{N}\right)}=\left\|f\left(u_{n}\right) v-f(u) v\right\|_{L^{s}(\Omega)} \rightarrow 0
$$

This finishes the proof of (ii). (iii) is a direct consequence of (i) and (ii).
The next lemma is crucial to prove that $u$ is a critical point of $J$.
Lemma 4.4. There is a subsequence such that
(i) $\nabla u_{n}(x) \rightarrow \nabla u(x)$ a.e in $\mathbb{R}^{N}$,
(ii) $\phi\left(\left|\nabla u_{n}\right|\right) \frac{\partial u_{n}}{\partial x_{i}} \rightharpoonup \phi(|\nabla u|) \frac{\partial u}{\partial x_{i}}$ in $L^{\widetilde{\Phi}}\left(\mathbb{R}^{N}\right)$,
(iii) $\phi\left(\left|u_{n}\right|\right) u_{n} \rightharpoonup \phi(|u|) u$ in $L^{\widetilde{\Phi}}\left(\mathbb{R}^{N}\right)$.

Proof. (i) We begin this proof observing that $\left(\phi_{1}\right)$ yields

$$
\begin{equation*}
(\phi(|x|) x-\phi(|y|) y)(x-y)>0, \forall x, y \in \mathbb{R}^{N} \text { with } x \neq y \tag{4.3}
\end{equation*}
$$

where $(\cdot, \cdot)$ denotes the usual inner product in $\mathbb{R}^{N}$. Given $R>0$, let us consider $\xi=\xi_{R} \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ satisfying

$$
0 \leq \xi \leq 1, \xi \equiv 1 \text { in } B_{R}(0) \text { and } \operatorname{supp}(\xi) \subset B_{2 R}(0)
$$

Using (4.3), we get

$$
\begin{align*}
0 \leq & \int_{B_{R}(0)}\left(\phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}-\phi(|\nabla u(x)|) \nabla u(x)\right) \\
& \times\left(\nabla u_{n}(x)-\nabla u(x)\right) \mathrm{d} x \\
\leq & \int_{B_{2 R}(0)}\left(\phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x)-\phi(|\nabla u(x)|) \nabla u(x)\right) \\
& \times\left(\nabla u_{n}(x)-\nabla u(x)\right) \xi(x) \mathrm{d} x  \tag{4.4}\\
= & \int_{B_{2 R}(0)} \phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x)\left(\nabla u_{n}(x)-\nabla u(x)\right) \xi(x) \mathrm{d} x \\
- & \int_{B_{2 R}(0)} \phi(|\nabla u(x)|) \nabla u(x)\left(\nabla u_{n}(x)-\nabla u(x)\right) \xi(x) \mathrm{d} x .
\end{align*}
$$

Now, combining the boundedness of $\left\{\left(u_{n}-u\right) \xi\right\}$ in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ with the limit $\left\|J^{\prime}\left(u_{n}\right)\right\|=o_{n}(1)$, it follows that

$$
\begin{align*}
o_{n}(1)= & J^{\prime}\left(u_{n}\right)\left(\left(u_{n}-u\right) \xi\right)=\int_{B_{2 R}(0)} \phi\left(\left|\nabla u_{n}(x)\right|\right) \\
& \times \nabla u_{n}(x) \nabla\left(\left(u_{n}-u\right)(x) \xi(x)\right) \mathrm{d} x \\
& +\int_{B_{2 R}(0)} \phi\left(\left|u_{n}(x)\right|\right) u_{n}(x)\left(u_{n}-u\right)(x) \xi(x) \mathrm{d} x  \tag{4.5}\\
& -\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right)\left(u_{n}(y)-u(y)\right) \xi(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y .
\end{align*}
$$

The boundedness of $\left(u_{n}\right)$ in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ implies that

$$
\begin{align*}
& \int_{B_{2 R}(0)} \widetilde{\Phi}\left(\phi\left(\left|u_{n}(x)\right|\right) u_{n}(x)\right) \mathrm{d} x \leq \int_{B_{2 R}(0)} \Phi\left(2\left|u_{n}(x)\right|\right) \\
& \quad \times \mathrm{d} x \leq K \int_{B_{2 R}(0)} \Phi\left(\left|u_{n}(x)\right|\right) \mathrm{d} x \leq \widetilde{C} \tag{4.6}
\end{align*}
$$

where $\widetilde{C}$ is a constant that does not depend on $\left(u_{n}\right)$. Since

$$
\begin{aligned}
& \left\|\phi\left(\left|\nabla u_{n}\right|\right) u_{n}\right\|_{L^{\frac{m}{\Phi}}\left(B_{2 R}(0)\right)}^{\frac{m}{m}}-1 \leq \min \left(\left\|\phi\left(\left|\nabla u_{n}\right|\right) u_{n}\right\|_{L^{\tilde{\Phi}}\left(B_{2 R}(0)\right)}^{\frac{l}{l-1}}\right. \\
& \left.\left\|\phi\left(\left|\nabla u_{n}\right|\right) u_{n}\right\|_{L^{\frac{m}{\Phi}\left(B_{2 R}(0)\right)}}^{\frac{m}{m-1}}\right)
\end{aligned}
$$

it follows from (4.6) that the sequence $\left(\phi\left(\left|u_{n}\right|\right)\left|u_{n}\right|\right)$ is bounded in $L^{\widetilde{\Phi}}\left(B_{2 R}(0)\right)$. A similar reasoning implies that $\left(\phi\left(\left|\nabla u_{n}\right|\right)\left|\nabla u_{n}\right|\right)$ is bounded in $L^{\widetilde{\Phi}}\left(B_{2 R}(0)\right)$.

The compact embedding $W^{1, \Phi}\left(B_{2 R}(0)\right) \hookrightarrow L^{\Phi}\left(B_{2 R}(0)\right)$ implies that $\left\|u_{n}-u\right\|_{L^{\Phi}\left(B_{2 R}(0)\right)} \rightarrow 0$ for a subsequence. It follows that, up to a subsequence, we have

$$
\begin{align*}
& \left|\int_{B_{2 R}(0)} \phi\left(\left|u_{n}(x)\right|\right) u_{n}(x)\left(u_{n}-u\right)(x) \xi(x) \mathrm{d} x\right| \\
& \quad \leq \int_{B_{2 R}(0)} \phi\left(\left|u_{n}(x)\right|\right)\left|u_{n}(x) \|\left(u_{n}-u\right)(x)\right| \mathrm{d} x \\
& \quad \leq C\left\|\phi\left(\left|u_{n}\right|\right)\left|u_{n}\right|\right\|_{L^{\tilde{\Phi}}\left(B_{2 R}(0)\right)}\left\|u_{n}-u\right\|_{L^{\Phi}\left(B_{2 R}(0)\right)} \\
& \quad \rightarrow 0 \tag{4.7}
\end{align*}
$$

A similar reasoning implies that

$$
\begin{equation*}
\int_{B_{2 R}(0)} \phi\left(\left|\nabla u_{n}(x)\right|\right)\left(u_{n}-u\right)(x) \nabla u_{n}(x) \nabla \xi(x) \mathrm{d} x \rightarrow 0 \tag{4.8}
\end{equation*}
$$

for some subsequence. Standard arguments imply, for a subsequence, that

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right)\left(u_{n}-u\right)(y) \xi(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \rightarrow 0 \tag{4.9}
\end{equation*}
$$

Note that

$$
\begin{aligned}
\int_{B_{2 R}(0)} \widetilde{\Phi}\left(\xi(x) \phi(|\nabla u(x)|)\left|\frac{\partial u(x)}{\partial x_{i}}\right|\right) \mathrm{d} x & \leq C \int_{B_{2 R}(0)} \widetilde{\Phi}\left(\phi(|\nabla u(x)|)\left|\frac{\partial u(x)}{\partial x_{i}}\right|\right) \mathrm{d} x \\
& \leq \tilde{C} \int_{B_{2 R}(0)} \Phi(2|\nabla u(x)|) \mathrm{d} x \\
& \leq \tilde{K} \int_{B_{2 R}(0)} \Phi(|\nabla u(x)|) \mathrm{d} x
\end{aligned}
$$

Thus from Hölder inequality, it follows that the function

$$
L_{i}(v):=\int_{B_{2 R}(0)} \xi(x) \phi(|\nabla u(x)|) \frac{\partial u(x)}{\partial x_{i}} v(x) \mathrm{d} x, i=1, \ldots, N
$$

with $v \in L^{\Phi}\left(B_{2 R}(0)\right)$ defines a linear continuous functional. Since $\frac{\partial u_{n}}{\partial x_{i}} \rightharpoonup$ $\frac{\partial u}{\partial x_{i}}$ in $L^{\Phi}\left(B_{2 R}(0)\right)$, we get

$$
\begin{equation*}
\int_{B_{2 R}(0)} \phi(|\nabla u(x)|) \nabla u(x)\left(\nabla u_{n}(x)-\nabla u(x)\right) \xi(x) \mathrm{d} x \rightarrow 0 . \tag{4.10}
\end{equation*}
$$

From (4.4), (4.5), (4.7), (4.8) (4.9), (4.10) we obtain that

$$
\int_{B_{2 R}(0)} \phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x)\left(\nabla u_{n}(x)-\nabla u(x)\right) \xi(x) \mathrm{d} x \rightarrow 0 .
$$

Therefore

$$
\int_{B_{R}(0)} \phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x)\left(\nabla u_{n}(x)-\nabla u(x)\right) \mathrm{d} x \rightarrow 0 .
$$

Applying a result found in Dal Maso and Murat [40], it follows that

$$
\nabla u_{n}(x) \rightarrow \nabla u(x) \text { a.e in } \quad B_{R}(0),
$$

for each $R>0$. Since $R$ is arbitrary, there is a subsequence of $\left(u_{n}\right)$, still denoted by itself, such that

$$
\nabla u_{n}(x) \rightarrow \nabla u(x) \text { a.e in } \mathbb{R}^{N} .
$$

ii) We have $\phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x) \rightarrow \phi(|\nabla u(x)|) \nabla u(x)$ a.e in $\mathbb{R}^{N}$. Note that

$$
\widetilde{\Phi}\left(\phi\left(\left|\nabla u_{n}(x)\right|\right)\left|\nabla u_{n}(x)\right|\right) \leq \tilde{C} \Phi\left(2\left|\nabla u_{n}(x)\right|\right) \leq \tilde{K} \Phi\left(\left|\nabla u_{n}(x)\right|\right) .
$$

Then the boundedness of $\left(u_{n}\right)$ in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$ implies that the sequences $\left(\phi\left(\left|\nabla u_{n}\right|\right) \frac{\partial u_{n}}{\partial x_{i}}\right), i=1, \ldots, N$ are bounded in $L^{\widehat{\Phi}}\left(\mathbb{R}^{N}\right)$. By [41], it follows that

$$
\phi\left(\left|\nabla u_{n}\right|\right) \frac{\partial u_{n}}{\partial x_{i}} \rightharpoonup \phi(|\nabla u|) \frac{\partial u}{\partial x_{i}} \quad \text { in } \quad L^{\widetilde{\Phi}}\left(\mathbb{R}^{N}\right) .
$$

A similar reasoning implies (iii).
Now, we are ready to prove that $u$ is a critical point of $J$.

Lemma 4.5. The function $u$ is a critical point of $J$, that is, $J^{\prime}(u)=0$.
Proof. First of all, we claim that

$$
J^{\prime}\left(u_{n}\right) v \rightarrow J^{\prime}(u) v, \quad \forall v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)
$$

To verify such limit, note that

$$
\begin{aligned}
J^{\prime}\left(u_{n}\right) v= & \int_{\mathbb{R}^{N}} \phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x) \nabla v(x)+\phi\left(\left|u_{n}\right|\right) u_{n}(x) v(x) \mathrm{d} x \\
& -\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right) v(y)}{|x-y|^{\lambda(x, y)}} \mathrm{d} x \mathrm{~d} y .
\end{aligned}
$$

By Lemmas 4.3 and 4.4,

$$
\begin{gather*}
\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right) v(y)}{|x-y|^{\lambda(x, y)}} \mathrm{d} x \mathrm{~d} y \rightarrow \int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F(u(x)) f(u(y)) v(y)}{|x-y|^{\lambda(x, y)}} \mathrm{d} x \mathrm{~d} y,  \tag{4.11}\\
\int_{\mathbb{R}^{N}} \phi\left(\left|\nabla u_{n}(x)\right|\right) \nabla u_{n}(x) \nabla v(x) \mathrm{d} x \rightarrow \int_{\mathbb{R}^{N}} \phi(|\nabla u(x)|) \nabla u(x) \nabla v(x) \mathrm{d} x, \tag{4.12}
\end{gather*}
$$

and

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} \phi\left(\left|u_{n}(x)\right|\right) u_{n}(x) v(x) \mathrm{d} x \rightarrow \int_{\mathbb{R}^{N}} \phi(|u(x)|) u(x) v(x) \mathrm{d} x \tag{4.13}
\end{equation*}
$$

From the relations (4.11), (4.12) and (4.13) we have the claim. Since $J^{\prime}\left(u_{n}\right) v \rightarrow 0$, the claim ensures that $J^{\prime}(u) v=0$, for all $v \in C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$. Now, the lemma follows using the fact that $C_{0}^{\infty}\left(\mathbb{R}^{N}\right)$ is dense in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$.

### 4.1. Proof of Theorem 4.1

If $u \neq 0$, then $u$ is a nontrivial solution and the theorem is proved. If $u=0$, we must find another solution $v \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) \backslash\{0\}$ for the equation (1.3). For such purpose, the claim below is crucial in our argument.
Claim 4.1. There exist $r>0, \beta>0$ and a sequence $\left(y_{n}\right) \subset \mathbb{R}^{N}$ such that

$$
\liminf _{n \rightarrow+\infty} \int_{B_{r}\left(y_{n}\right)} \Phi\left(\left|u_{n}(x)\right|\right) d x \geq \beta>0
$$

Proof. In fact, if the above claim does not hold, using Theorem 2.2, we derive the limit

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} P\left(\left|u_{n}\right|\right) \mathrm{d} x \rightarrow 0 \tag{4.14}
\end{equation*}
$$

for any $N$-function $P$ satisfying $\left(P_{1}\right)-\left(P_{2}\right)$. Applying Proposition 1.1,

$$
\begin{aligned}
& \left|\int_{\mathbb{R}^{N}} \int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right) u_{n}(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y\right| \\
& \quad \leq C\left\|F\left(u_{n}(x)\right)\right\|_{L^{s}\left(\mathbb{R}^{N}\right)}\left\|f\left(u_{n}(y)\right) u_{n}(y)\right\|_{L^{s}\left(\mathbb{R}^{N}\right)}
\end{aligned}
$$

By $\left(f_{1}\right),(F),(E)$ and Theorem 2.2,

$$
\int_{\mathbb{R}^{N}}\left|F\left(u_{n}(x)\right)\right|^{s} \mathrm{~d} x \rightarrow 0
$$

and

$$
\int_{\mathbb{R}^{N}}\left|f\left(u_{n}(y)\right) u_{n}(y)\right|^{s} \mathrm{~d} y \rightarrow 0
$$

Therefore

$$
\int_{\mathbb{R}^{N}} \frac{F\left(u_{n}(x)\right) f\left(u_{n}(y)\right) u_{n}(y)}{|x-y|^{\lambda}} \mathrm{d} x \mathrm{~d} y \rightarrow 0
$$

The above limit together with the fact that $J^{\prime}\left(u_{n}\right) u_{n}=o_{n}(1)$ give

$$
\int_{\mathbb{R}^{N}} \phi\left(\left|\nabla u_{n}(x)\right|\right)\left|\nabla u_{n}(x)\right|^{2}+\phi\left(\left|u_{n}(x)\right|\right)\left|\nabla u_{n}(x)\right|^{2} \mathrm{~d} x \rightarrow 0 .
$$

From $\left(\phi_{3}\right)$,

$$
\int_{\mathbb{R}^{N}} \Phi\left(\left|\nabla u_{n}(x)\right|\right)+\Phi\left(\left|u_{n}(x)\right|\right) \mathrm{d} x \rightarrow 0
$$

This limit leads to $J\left(u_{n}\right) \rightarrow 0$, which contradicts the limit $J\left(u_{n}\right) \rightarrow d>0$.
Using standard arguments, we can assume in Claim 4.1 that $\left(y_{n}\right) \subset \mathbb{Z}^{N}$. By setting $v_{n}(x)=u_{n}\left(x+y_{n}\right)$, it follows that

$$
J\left(v_{n}\right)=J\left(u_{n}\right),\left\|J^{\prime}\left(v_{n}\right)\right\|=\left\|J^{\prime}\left(u_{n}\right)\right\| \quad \text { and } \quad\left\|u_{n}\right\|_{1, \Phi}=\left\|v_{n}\right\|_{1, \Phi} \quad \forall n \in \mathbb{N}
$$

From the above information, we have that $J\left(v_{n}\right) \rightarrow d$ and $J^{\prime}\left(v_{n}\right) \rightarrow 0$. Since $\left(v_{n}\right)$ is bounded in $W^{1, \Phi}\left(\mathbb{R}^{N}\right)$, up to a subsequence, $v_{n} \rightarrow v$ in $L^{\Phi}\left(B_{r}(0)\right)$, for some $v \in W^{1, \Phi}$. To verify that $v \neq 0$, note that by Claim 4.1, we have for some subsequence

$$
\begin{aligned}
0 & <\beta \leq \lim _{n \rightarrow+\infty} \int_{B_{r}\left(y_{n}\right)} \Phi\left(\left|u_{n}(x)\right|\right) \mathrm{d} x \\
& =\lim _{n \rightarrow+\infty} \int_{B_{r}(0)} \Phi\left(\left|v_{n}(x)\right|\right) \mathrm{d} x=\int_{B_{r}(0)} \Phi(|v(x)|) \mathrm{d} x .
\end{aligned}
$$

Applying the same arguments as in the proofs of Lemmas 4.3, 4.4 and 4.5 for the sequence $\left(v_{n}\right)$ we obtain the desired result.

## 5. Final Comments

The same arguments used in this paper can be applied to study the existence of solutions for related problems of the following type:

$$
\left\{\begin{align*}
-\Delta_{\Phi} u+V(x) \phi(|u|) u & =\left(\int_{\mathbb{R}^{N}} \frac{F(u(x))}{|x-y|^{\lambda}}\right) f(u(y)), \text { in } \mathbb{R}^{N},  \tag{5.1}\\
u & \in W^{1, \Phi}\left(\mathbb{R}^{N}\right) .
\end{align*}\right.
$$

The potential $V: \mathbb{R}^{N} \rightarrow \mathbb{R}$ is a continuous functions with $\inf _{x \in \mathbb{R}^{N}} V(x)>0$ that belongs to one of the following classes:
Class 1: $V$ is periodic: $V$ is a $\mathbb{Z}^{N}$-periodic function, that is,

$$
V(x+y)=V(x), \quad \forall x \in \mathbb{R}^{N} \quad \text { and } \quad y \in \mathbb{Z}^{N}
$$

Class 2: $V$ is asymptotic periodic function: There is a $\mathbb{Z}^{N}$-periodic function $V_{p}: \mathbb{R}^{N} \rightarrow \mathbb{R}$ such that

$$
V(x)<V_{p}(x), \quad \forall x \in \mathbb{R}^{N}
$$

and

$$
\left|V(x)-V_{p}(x)\right| \rightarrow 0 \quad \text { as } \quad|x| \rightarrow+\infty .
$$

Class 3: $V$ is coercive $V$ is a coercive function, that is,

$$
V(x) \rightarrow+\infty \quad \text { as } \quad|x| \rightarrow+\infty .
$$

Class 4: $V$ is a Bartsch-Wang-like potential: The potential $V$ verifies the following property

$$
\operatorname{med}\left(\left\{x \in \mathbb{R}^{N}: V(x) \leq M\right\}\right)<+\infty, \quad \text { for all } M>0
$$

## Acknowledgements

This work started when Leandro S. Tavares was visiting the Federal University of Campina Grande. He thanks the hospitality of Professor Claudianor Alves and of the other members of the department. V. D. Rădulescu was supported by the Slovenian Research Agency grants P1-0292, J1-8131, J17025, N1-0064, and N1-0083. He also acknowledges the support through the Project MTM2017-85449-P of the DGISPI (Spain). C. O. Alves was partially supported by CNPq/Brazil 301807/2013-2.

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## References

[1] Fröhlich, H.: Theory of electrical breakdown in ionic crystal. Proc. R. Soc. Ser. A 160(901), 230-241 (1937)
[2] Pekar, S.: Untersuchung über die Elektronentheorie der Kristalle. Akademie, Berlin (1954)
[3] Lieb, E. H.: Existence and uniqueness of the minimizing solution of Choquard's nonlinear equation. Stud. Appl. Math. 57, 93-105 (1976/77)
[4] Moroz, I.M., Penrose, R., Tod, P.: Spherically-symmetric solutions of the Schrödinger-Newton equations. Class. Quant. Grav. 15, 2733-2742 (1998)
[5] Ackermann, N.: On a periodic Schrödinger equation with nonlocal superlinear part. Math. Z. 248, 423-443 (2004)
[6] Alves, C.O., Figueiredo, G.M., Yang, M.: Existence of solutions for a nonlinear Choquard equation with potential vanishing at infinity. Adv. Nonlinear Anal. 5(4), 331-345 (2016)
[7] Alves, C.O., Yang, M.: Multiplicity and concentration behavior of solutions for a quasilinear Choquard equation via penalization method. Proc. R. Soc. Edinburgh Sect. A 146, 23-58 (2016)
[8] Alves, C.O., Yang, M.: Existence of semiclassical ground state solutions for a generalized Choquard equation. J. Differ. Equ. 257, 4133-4164 (2014)
[9] Cingolani, S., Secchi, S., Squassina, M.: Semi-classical limit for Schrödinger equations with magnetic field and Hartree-type nonlinearities. Proc. R. Soc. Edinburgh Sect. A 140, 973-1009 (2010)
[10] Gao, F., Yang, M.: On the Brezis-Nirenberg type critical problem for nonlinear Choquard equation. Sci. China Math. 61, 1219-1242 (2018)
[11] Lions, P.L.: The Choquard equation and related questions. Nonlinear Anal. 4, 1063-1072 (1980)
[12] Ma, L., Zhao, L.: Classification of positive solitary solutions of the nonlinear Choquard equation. Arch. Ration. Mech. Anal. 195, 455-467 (2010)
[13] Moroz, V., van Schaftingen, J.: Existence of ground states for a class of nonlinear Choquard equations. Trans. Am. Math. Soc. 367, 6557-6579 (2015)
[14] Moroz, V., van Schaftingen, J.: Semi-classical states for the Choquard equation. Calc. Var. Partial Differ. Equ. 52, 199-235 (2015)
[15] Moroz, V., van Schaftingen, J.: Ground states of nonlinear Choquard equations: Hardy-Littlewood-Sobolev critical exponent. Commun. Contemp. Math. 17, 12 (2015)
[16] Moroz, V., van Schaftingen, J.: A guide to the Choquard equation. J. Fixed Point Theory Appl. 19(1), 773-813 (2017)
[17] van Schaftingen, J., Xia, J.: Standing waves with a critical frequency for nonlinear Choquard equations. Nonlinear Anal. 161, 87-107 (2017)
[18] Wang, T.: Existence of positive ground-state solution for Choquard-type equations. Mediterr. J. Math. 14(1), 15 (2017)
[19] Lieb, E., Loss, M.: Analysis, Graduate Studies in Mathematics. Amer. Math. Soc, Providence (2001)
[20] DiBenedetto, E.: $C^{1, \gamma}$ local regularity of weak solutions of degenerate elliptic equations. Nonlinear Anal. 7, 827-850 (1985)
[21] Fukagai, N., Narukawa, K.: On the existence of multiple positive solutions of quasilinear elliptic eigenvalue problems. Ann. Mat. Pura Appl. 186, 539-564 (2007)
[22] Azzollini, A., d'Avenia, P., Pomponio, A.: Quasilinear elliptic equations in $\mathbb{R}^{N}$ via variational methods and Orlicz-Sobolev embeddings. Calc. Var. Partial Differ. Equ. 49, 197-213 (2014)
[23] Alves, C.O., da Silva, A.R.: Multiplicity and concentration of positive solutions for a class of quasilinear problems through Orlicz-Sobolev space. J. Math. Phys. 57, 111502 (2016). https://doi.org/10.1063/1.4966534
[24] Alves, C.O., da Silva, A.R.: Multiplicity and concentration behavior of solutions for a quasilinear problem involving N -functions via penalization method. Electron. J. Differ. Equ. 2016(158), 1-24 (2016)
[25] Alves, C.O., Figueiredo, G.M., Santos, J.A.: Strauss and Lions type results for a class of Orlicz-Sobolev spaces and applications. Topol. Methods Nonlinear Anal. 44, 435-456 (2014)
[26] Bonanno, G., Molica Bisci, G., Rădulescu, V.D.: Quasilinear elliptic nonhomogeneous dirichlet problems through Orlicz-Sobolev spaces. Nonlinear Anal. 75, 4441-4456 (2012)
[27] Bonanno, G., Molica Bisci, G., Rădulescu, V.D.: Existence and multiplicity of solutions for a quasilinear nonhomogeneous problems: An Orlicz-Sobolev space setting. J. Math. Anal. Appl. 330, 416-432 (2007)
[28] Fukagai, N., Ito, M., Narukawa, K.: Quasilinear elliptic equations with slowly growing principal part and critical Orlicz-Sobolev nonlinear term. Proc. R. Soc. Edinburgh Sect. A 139, 73-106 (2009)
[29] Le, V.K., Motreanu, D., Motreanu, V.V.: On a non-smooth eigenvalue problem in Orlicz-Sobolev spaces. Appl. Anal. 89, 229-242 (2010)
[30] Mihailescu, M., Rădulescu, V., Repovš, D.: On a non-homogeneous eigenvalue problem involving a potential: an Orlicz-Sobolev space setting. J. Math. Pures Appl. 93, 132-148 (2010)
[31] Rădulescu, V.D.: Nonlinear elliptic equations with variable exponent: old and new. Nonlinear Anal. 121, 336-369 (2015)
[32] Rădulescu, V.D., Repovš, D.: Partial Differential Equations with Variable Exponents Variational Methods and Qualitative Analysis, Monographs and Research Notes in Mathematics. CRC, Boca Raton (2015)
[33] Repovš, D.: Stationary waves of Schrödinger-type equations with variable exponent. Anal. Appl. 13, 645-661 (2015)
[34] Santos, J.A.: Multiplicity of solutions for quasilinear equations involving critical Orlicz-Sobolev nonlinear terms. Electron. J. Differ. Equ. 249(2013), 1-13 (2013)
[35] Santos, J.A., Soares, S.H.M.: Radial solutions of quasilinear equations in OrliczSobolev type spaces. J. Math. Anal. Appl. 428, 1035-1053 (2015)
[36] Adams, A., Fournier, J.F.: Sobolev Spaces, 2nd edn. Academic, Oxford (2003)
[37] Rao, M.N., Ren, Z.D.: Theory of Orlicz Spaces. Marcel Dekker, New York (1985)
[38] Donaldson, T.K., Trudinger, N.S.: Orlicz-Sobolev spaces and embedding theorems. J. Funct. Anal. 8, 52-75 (1971)
[39] Chabrowski, J.: Variational Methods for Potential Operator Equations with Applications to Nonlinear Elliptic Equations. Walter de Gruyter, Berlin-New York (1997)
[40] Dal Maso, G., Murat, F.: Almost everywhere convergence of gradients of solutions to nonlinear elliptic systems. Nonlinear Anal. 31, 405-412 (1998)
[41] Gossez, J.P.: Orlicz-Sobolev spaces and nonlinear elliptic boundary value problems. In: Nonlinear Analysis, Function Spaces and Applications (Proc. Spring School, Horni Bradlo, 1978). Teubner, Leipzig, pp. 59-94 (1979)

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Received: June 28, 2018.
Revised: September 28, 2018.
Accepted: December 21, 2018.
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