# Infinitely many solutions for a class of sublinear Schrödinger equations with indefinite potentials 

Anouar Bahrouni and Hichem Ounaies<br>Mathematics, Faculty of Sciences, University of Monastir, 5019 Monastir, Tunisia (bahrounianouar@yahoo.fr; hichem.ounaies@fsm.rnu.tn)

Vicenţiu D. Rădulescu
Department of Mathematics, Faculty of Science, King Abdulaziz University, PO Box 80203, Jeddah 21589, Saudi Arabia (vicentiu.radulescu@math.cnrs.fr)
and
'Simion Stoilow' Institute of Mathematics of the Romanian Academy, PO Box 1-764, 014700 Bucharest, Romania
(MS received 10 July 2013; accepted 11 February 2014)


#### Abstract

In this paper we are concerned with qualitative properties of entire solutions to a Schrödinger equation with sublinear nonlinearity and sign-changing potentials. Our analysis considers three distinct cases and we establish sufficient conditions for the existence of infinitely many solutions.


## 1. Historical perspective of the Schrödinger equation

The Schrödinger equation plays the role of Newton's laws and conservation of energy in classical mechanics, that is, it predicts the future behaviour of a dynamic system. The linear Schrödinger equation is a central tool of quantum mechanics, which provides a thorough description of a particle in a non-relativistic setting. Schrödinger's linear equation is

$$
\Delta \psi+\frac{8 \pi^{2} m}{\hbar^{2}}(E-V(x)) \psi=0
$$

where $\psi$ is the Schrödinger wave function, $m$ is the mass, $\hbar$ denotes Planck's constant, $E$ is the energy and $V$ stands for the potential energy.

The structure of the nonlinear Schrödinger equation is much more complicated. This equation is a prototypical dispersive nonlinear partial differential equation that has been central for almost four decades to a variety of areas in mathematical physics. The relevant fields of application vary from Bose-Einstein condensates and nonlinear optics (see [15]), propagation of the electric field in optical fibres (see $[26,32]$ ) to the self-focusing and collapse of Langmuir waves in plasma physics (see [43]) and the behaviour of deep water waves and freak waves (so-called rogue waves) in the ocean (see [34]). The nonlinear Schrödinger equation also describes various phenomena arising in the theory of Heisenberg ferromagnets and magnons,
the self-channelling of a high-power ultra-short laser in matter, condensed matter theory, dissipative quantum mechanics, electromagnetic fields (see [5]) and plasma physics (for example, the Kurihara superfluid film equation). We refer to Ablowitz et al. [1] and Sulem [36] for a modern overview, including applications.

Schrödinger also established the classical derivation of his equation based upon the analogy between mechanics and optics and closer to de Broglie's ideas. He developed a perturbation method, inspired by the work of Lord Rayleigh in acoustics, proved the equivalence between his wave mechanics and Heisenberg's matrix mechanics and introduced the time-dependent Schrödinger equation

$$
\begin{equation*}
\mathrm{i} \hbar \psi_{t}=-\frac{\hbar^{2}}{2 m} \Delta \psi+V(x) \psi-\gamma|\psi|^{p-1} \psi \quad \text { in } \mathbb{R}^{N}(N \geqslant 2) \tag{1.1}
\end{equation*}
$$

where $p<2 N /(N-2)$ if $N \geqslant 3$ and $p<+\infty$ if $N=2$. In physical problems a cubic nonlinearity corresponding to $p=3$ is common; in this case (1.1) is called the Gross-Pitaevskii equation. In the study of (1.1), Floer and Weinstein [24] and Oh [33] supposed that the potential $V$ is bounded and possesses a non-degenerate critical point at $x=0$. More precisely, it is assumed that $V$ belongs to the class $\left(V_{a}\right)$ (for some real number $a$ ) introduced by Kato [29]. Taking $\gamma>0$ and $\hbar>0$ sufficiently small and using a Lyapunov-Schmidt type reduction, Oh [33] proved the existence of standing wave solutions of (1.1), that is, a solution of the form

$$
\begin{equation*}
\psi(x, t)=\mathrm{e}^{-\mathrm{i} E t / \hbar} u(x) \tag{1.2}
\end{equation*}
$$

Using the ansatz (1.2), we reduce the nonlinear Schrödinger equation (1.1) to the semilinear elliptic equation

$$
-\frac{\hbar^{2}}{2 m} \Delta u+(V(x)-E) u=|u|^{p-1} u
$$

The change of variable $y=\hbar^{-1} x$ (and replacing $y$ by $x$ ) yields

$$
\begin{equation*}
-\Delta u+2 m\left(V_{\hbar}(x)-E\right) u=|u|^{p-1} u \quad \text { in } \mathbb{R}^{N} \tag{1.3}
\end{equation*}
$$

where $V_{\hbar}(x)=V(\hbar x)$.
If, for some $\xi \in \mathbb{R}^{N} \backslash\{0\}, V(x+s \xi)=V(x)$ for all $s \in \mathbb{R}$, then (1.1) is invariant under the Galilean transformation

$$
\psi(x, t) \mapsto \psi(x-\xi t, t) \exp \left(\mathrm{i} \xi \cdot x / \hbar-\frac{1}{2} \mathrm{i}|\xi|^{2} t / \hbar\right) \psi(x-\xi t, t)
$$

Thus, in this case, standing waves reproduce solitary waves travelling in the direction of $\xi$. In other words, Schrödinger discovered that the standing waves are scalar waves rather than vector electromagnetic waves. This is an important difference: vector electromagnetic waves are mathematical waves that describe a direction (vector) of force, whereas the wave motions of space are scalar waves, which are simply described by their wave amplitude. The importance of this discovery was pointed out by Einstein [23], who wrote:

The Schrödinger method, which has in a certain sense the character of a field theory, does indeed deduce the existence of only discrete states, in surprising agreement with empirical facts. It does so on the basis of differential equations applying a kind of resonance argument.

In a celebrated paper, Rabinowitz [35] proved that (1.3) has a ground-state solution (mountain pass solution) for $\hbar>0$ small, under the assumption that $\inf _{x \in \mathbb{R}^{N}} V(x)>E$. After making a standing wave ansatz, Rabinowitz reduced the problem to that of studying the semilinear elliptic equation

$$
\begin{equation*}
-\Delta u+V(x) u=f(x, u) \quad \text { in } \mathbb{R}^{N} \tag{1.4}
\end{equation*}
$$

under suitable conditions on $V$ and assuming that $f$ is smooth, superlinear and has a subcritical growth.

## 2. Introduction and main results

In the present paper we are concerned with the existence of infinitely many solutions of the semilinear Schrödinger equation

$$
\begin{equation*}
-\Delta u+V(x) u=a(x) g(u) \quad x \in \mathbb{R}^{N}(N \geqslant 3) \tag{2.1}
\end{equation*}
$$

where $V$ and $a$ are functions changing sign and the nonlinearity $g$ has a sublinear growth. Such problems in $\mathbb{R}^{N}$ arise naturally in various branches of physics and present challenging mathematical difficulties.

If (2.1) is considered in a bounded domain $\Omega$, with the Dirichlet boundary condition, then there is a large literature on existence and a multiplicity of solutions (see [4, 14, 27, 28, 37, 38, 41]). In particular, Kajikiya [27] has considered such sublinear cases with sign-changing nonlinearity and has proved the existence of infinitely many solutions.

If $\Omega$ is an unbounded domain, and especially if $\Omega=\mathbb{R}^{N}$, then the existence and multiplicity of non-trivial solutions for (2.1) have been extensively investigated in the literature over the past several decades, both for sublinear and superlinear nonlinearities.

In the superlinear case, we can cite $[2,3,6,8,17,19,21,22,25,35,42]$. In particular, Costa and Tehrani [17] have considered the problem

$$
\begin{equation*}
-\Delta u-\lambda h(x) u=a(x) g(u), \quad u>0, \text { in } \mathbb{R}^{N} \tag{2.2}
\end{equation*}
$$

where $\lambda>0, h$ is a positive function, $a$ changes the sign in $\mathbb{R}^{N}, N \geqslant 3$, and $g$ is a superlinear function. With further assumptions on $h, a$ and $g$, they proved the existence of $\lambda_{1}(h)>0$ such that $(2.2)$ admits one positive solution for $0<\lambda<\lambda_{1}(h)$ and two positive solutions for $\lambda_{1}(h)<\lambda<\lambda_{1}(h)+\varepsilon$ for some $\varepsilon>0$.

In recent years, many authors have studied the question of existence and multiplicity of solutions for (2.1) with sublinear nonlinearity (see [7,10-12,16,18, 30, 39]). In most of the problems studied in these papers, $V$ and $a$ are considered to be positive. In particular, Brezis and Kamin [12] gave a sufficient and necessary condition for the existence of bounded positive solutions of (2.1) with $V=0$ and $a>0$.

Balabane et al. [7] proved that for each integer $k$, (2.1) has a radially compactly supported solution that has $k$ zeros in its support provided that $V=a=-1$ and $g(u)=|u|^{-2 \theta} u$, where $\left.\theta \in\right] 0, \frac{1}{2}[$.

Zhang and Wang [44] proved the existence of infinitely many solutions for (2.1) with $g(u)=|u|^{p-1} u, 0<p<1$, and the potentials $V>0, a>0$ satisfy the following assumptions:
$\left(\mathrm{S}_{1}\right) V \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$ and there exists $r>0$ such that

$$
m\{x \in B(y, r) ; V(x) \leqslant M\} \rightarrow 0 \text { as }|y| \rightarrow+\infty \quad \forall M>0
$$

where $m$ is the Lebesgue measure in $\mathbb{R}^{N}$;
$\left(\mathrm{S}_{2}\right) a: \mathbb{R}^{N} \rightarrow \mathbb{R}$ is a continuous function and $a \in L^{2 /(1-p)}\left(\mathbb{R}^{N}\right), 0<p<1$.
If $V$ and $a$ both change sign on $\mathbb{R}^{N}$, various difficulties arise. To the authors' knowledge, few results are known in this case. On this subject, Costa and Tehrani [18] have proved the existence of at least one non-trivial solution for the equation

$$
-\Delta u+V(x) u=\lambda u+g(x, u)
$$

under the following conditions:
$\left(\mathrm{VC}_{1}\right) V \in C^{\beta}\left(\mathbb{R}^{N}\right)(0<\beta<1)$ and $\lim _{|x| \rightarrow+\infty} V(x)=0 ;$
$\left(\mathrm{VC}_{2}\right) \int_{\mathbb{R}^{N}}\left(|\nabla \varphi|^{2}+V(x) \varphi^{2}\right) \mathrm{d} x<0$ for some $\varphi \in C_{0}^{1}\left(\mathbb{R}^{N}\right) ;$
$\left(\mathrm{GC}_{1}\right)|g(x, s)| \leqslant b_{1}(x)|s|^{\alpha}+b_{2}(x)$ for some $0<\alpha<1$ and a class of integrable functions $b_{1}$ and $b_{2}$;
$\left(\mathrm{GC}_{2}\right) \lambda<0$ is an eigenvalue of the Schrödinger operator $L_{V}=-\Delta+V(x)$ in $\mathbb{R}^{N}$.
$\left(\mathrm{GC}_{2}\right) \lim _{\substack{\left\|u_{0}\right\| \rightarrow+\infty, u_{0} \in \operatorname{Ker}(-\Delta+V-\lambda)}} \frac{1}{\left\|u_{0}\right\|^{2 \alpha}} \int_{\mathbb{R}^{N}} G\left(x, u_{0}(x)\right) \mathrm{d} x= \pm \infty$.
Tehrani [39] studied the perturbed equation

$$
\begin{equation*}
-\Delta u+V(x) u=a(x) g(u)+f \tag{2.3}
\end{equation*}
$$

where $a$ and $V$ change sign on $\mathbb{R}^{N}, f \in L^{2}\left(\mathbb{R}^{N}\right)$ and $g$ is a sublinear function. With further assumptions on $a, V, f$ and $g$, he proved the existence of at least one non-trivial solution.

Costa and Chabrowski [16] considered the $p$-Laplacian equation

$$
\begin{equation*}
-\Delta_{p} u-\lambda V(x)|u|^{p-2} u=a(x)|u|^{q-2} u, \quad x \in \mathbb{R}^{N} \tag{2.4}
\end{equation*}
$$

where $\lambda \in \mathbb{R}$ is a parameter, $1<q<p<p^{*}=N p /(N-p), V \in L^{N / p}\left(\mathbb{R}^{N}\right) \cap$ $L^{\infty}\left(\mathbb{R}^{N}\right), a \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and $\lim _{|x| \rightarrow+\infty} a(x)=a_{\infty}<0$. With further assumptions on $a$ and $V$, they proved the existence of $\lambda_{1}>0$ and $\lambda_{-1}<0$ such that (2.4) admits at least one positive solution for $\lambda_{-1}<\lambda<\lambda_{1}$ and two positive solutions for $\lambda>\lambda_{1}$ and $\lambda<\lambda_{-1}$.

Benrhouma [9] proved the existence of at least three solutions for (2.3) with $g(u)=|u|^{p} \operatorname{sgn}(u), 0<p<1, V$ changing sign and $a<0$.

In all works cited above, where $a$ and $V$ change sign the authors proved the existence of at most three solutions. In this paper, we prove the existence of infinitely many solutions of (2.1) with $a$ and $V$ changing sign, under various assumptions on these potential functions.

Denote by $s$ the best Sobolev constant,

$$
s=\inf \left\{\|\nabla u\|_{2}^{2}, u \in W^{1,2}\left(\mathbb{R}^{N}\right), \int_{\mathbb{R}^{N}}|u(x)|^{2 N /(N-2)} \mathrm{d} x=1\right\}, \quad N \geqslant 3
$$

We suppose the following hypotheses on $g$ :
$\left(\mathrm{G}_{1}\right) g \in C(\mathbb{R}, \mathbb{R}), g$ is odd and there exist $\left.c>0, q \in\right] 0,1[$ such that

$$
|g(x)| \leqslant c|x|^{q} \quad \text { for all } x \in \mathbb{R}
$$

$\left(\mathrm{G}_{2}\right) \lim _{x \rightarrow 0} \frac{G(x)}{|x|^{2}}=+\infty, \quad$ where $G(x)=\int_{0}^{x} g(t) \mathrm{d} t \quad \forall x \in \mathbb{R} ;$
$\left(\mathrm{G}_{3}\right) G$ is positive on $\mathbb{R} \backslash\{0\}$.
We give three theorems on the existence of infinitely many solutions to the nonlinear problem (2.1).
Theorem 2.1. Assume that $g(x)=|x|^{q-1} x, 0<q<1$, and that $V$ satisfies:
$\left(\mathrm{V}_{1}\right) V \in L^{\infty}\left(\mathbb{R}^{N}\right), \lim _{|x| \rightarrow+\infty} V(x)=v_{\infty}>0$ and

$$
\left\|V^{-}\right\|_{N / 2}<s
$$

where $u^{\mp}(x)=\max \{\mp u(x), 0\}$ for all $x \in \mathbb{R}^{N}$ and for all $u \in E$.
Assume also that a satisfies:
$\left(\mathrm{A}_{1}\right) a \in L^{\infty}\left(\mathbb{R}^{N}\right), \lim _{|x| \rightarrow+\infty} a(x)=a_{\infty}<0$ and there exist $y=\left(y_{1}, \ldots, y_{N}\right) \in$ $\mathbb{R}^{N}, R_{0}>0$ such that

$$
a(x)>0 \quad \text { for all } x \in B\left(y, R_{0}\right)
$$

Then (2.1) possesses a sequence of non-trivial solutions converging to 0 .
In the next two theorems we change the assumption of boundedness of $a$ by the integrability condition. The last assumption was supported to make the energy functional associated with (2.1) well defined and to guarantee that the functional $F(u)=\int_{\mathbb{R}^{N}} a(x) G(u(x)) \mathrm{d} x$ has a compact gradient. This compactness property in turn was used to prove the required Palais-Smale condition, which is essential in the application of the critical point theory. We then have the following two multiplicity properties.

Theorem 2.2. Suppose that g satisfies $\left(G_{1}\right)-\left(G_{3}\right)$ and the potentials $V$ and a satisfy the following hypotheses:
$\left(\mathrm{V}_{2}\right) V \in L^{N / 2}\left(\mathbb{R}^{N}\right)$ and

$$
\left\|V^{-}\right\|_{N / 2}<s
$$

$\left(\mathrm{A}_{2}\right) a \in L^{2^{*} /\left(2^{*}-(q+1)\right)}\left(\mathbb{R}^{N}\right)$ and there exist $y \in \mathbb{R}^{N}$ and $R_{0}>0$ such that

$$
a(x)>0 \quad \forall x \in B\left(y, R_{0}\right)
$$

Then (2.1) possesses a bounded sequence of non-trivial solutions.

Theorem 2.3. Assume that $g$ satisfies $\left(G_{1}\right)-\left(G_{3}\right), V$ satisfies $\left(V_{1}\right)$ and a satisfies:
$\left(\mathrm{A}_{3}\right) a \in L^{2 /(1-q)}\left(\mathbb{R}^{N}\right)$ and there exist $y=\left(y_{1}, \ldots, y_{N}\right) \in \mathbb{R}^{N}$ and $R_{0}>0$ such that

$$
a(x)>0 \quad \forall x \in B\left(y, R_{0}\right)
$$

Then (2.1) possesses a bounded sequence of non-trivial solutions.
This paper is organized as follows. In $\S 2$ we give some notation, we present the variational framework and we recall some definitions and standard results. Then $\S \S 3-5$ are dedicated to the proof of theorems 2.1, 2.2 and 2.3.

## 3. Notation and preliminary results

In this section we present some notation and preliminaries that will be useful in the following. We make the following definitions:

- $\|u\|_{m}=\left(\int_{\mathbb{R}^{N}}|u(x)|^{m} \mathrm{~d} x\right)^{1 / m}, \quad 1 \leqslant m<+\infty ;$
- $2^{*}=\frac{2 N}{N-2}$ if $N \geqslant 3$ and $2^{*}=+\infty$ if $n \in\{1,2\}$;
- $B_{R}$ denotes the ball centred at the origin of radius $R>0$ in $\mathbb{R}^{N}$ and $B_{R}^{c}=$ $\mathbb{R}^{N} \backslash B_{R}$;
- $F^{\prime}(u)$ is the Fréchet derivative of $F$ at $u$.

Let $F_{1}, F_{2}$ be Banach spaces and let $T: F_{1} \rightarrow F_{2} . T$ is said to be a sequentially compact operator if, given any bounded sequence $\left(x_{n}\right)$ in $F_{1},\left(T\left(x_{n}\right)\right)$ has a convergent subsequence in $F_{2}$.

Let $E=H^{1}\left(\mathbb{R}^{N}\right) \cap L^{q+1}\left(\mathbb{R}^{N}\right)(0<q<1)$ be the reflexive Banach space endowed with the norm

$$
\|u\|=\|\nabla u\|_{2}+\|u\|_{q+1}
$$

Let $X=D^{1,2}\left(\mathbb{R}^{N}\right)=\left\{u \in L^{2^{*}}\left(\mathbb{R}^{N}\right) ; \nabla u \in\left(L^{2}\left(\mathbb{R}^{N}\right)\right)^{N}\right\}$, endowed with the norm

$$
\|u\|_{X}=\left(\int_{\mathbb{R}^{N}}|\nabla u(x)|^{2} \mathrm{~d} x\right)^{1 / 2}
$$

be a Hilbert space. Moreover, the embedding $X \subset L^{2^{*}}\left(\mathbb{R}^{N}\right)$ is continuous, which implies that

$$
S:=\inf \left\{\int_{\mathbb{R}^{N}}|\nabla u(x)|^{2} \mathrm{~d} x ; u \in X, \int_{\mathbb{R}^{N}}|u(x)|^{2^{*}} \mathrm{~d} x=1\right\}>0
$$

We refer the reader to [40, pp. 8 and 9] for more details.
Let

$$
Y=\left\{u \in H^{1}\left(\mathbb{R}^{N}\right) ; \int_{\mathbb{R}^{N}} V^{+}(x) u^{2}(x) \mathrm{d} x<+\infty\right\}
$$

under the hypotheses $V \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and $\operatorname{esslim}_{x \rightarrow+\infty} V(x)>0$. We endow $Y$ with the inner product

$$
\langle u, v\rangle=\int_{\mathbb{R}^{N}} \nabla u \nabla v+\int_{\mathbb{R}^{N}} V^{+}(x) u v \mathrm{~d} x
$$

and the associated norm $\|\cdot\|_{Y}$, which is equivalent to the usual norm

$$
\|u\|_{H^{1}}=\|\nabla u\|_{2}+\|u\|_{2} .
$$

Consider the functionals

$$
\begin{aligned}
& I(u)=\frac{1}{2} \int_{\mathbb{R}^{N}}\left(|\nabla u|^{2}(x)+V(x) u^{2}(x)\right) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x) G(u(x)) \mathrm{d} x \\
& \varphi(u)=\frac{1}{2} \int_{\mathbb{R}^{N}} V(x) u^{2}(x) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x) G(u(x)) \mathrm{d} x \\
& \psi(u)=-\frac{1}{2} \int_{\mathbb{R}^{N}} V^{-}(x) u^{2}(x) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x) G(u(x)) \mathrm{d} x .
\end{aligned}
$$

Under suitable assumptions on $a, G$ and $V$ (to be fixed later), $I, \varphi$ and $\psi$ are well defined and of class $C^{1}$ on $X, Y$ or $E$. A critical point of $I$ is a weak solution of (2.1).

Next, let us recall that a Palais-Smale (PS) sequence for the functional $I$ is a sequence $\left(u_{n}\right)$ such that

$$
I\left(u_{n}\right) \text { is bounded and }\left\|I^{\prime}\left(u_{n}\right)\right\| \rightarrow 0
$$

The functional $I$ is said to satisfy the PS condition if any PS sequence possesses a convergent subsequence.

A first main difficulty that appears in the study of (2.1) is the loss of compactness. In order to overcome this difficulty, we use the Lions compactness principle [31]. A second main difficulty is to satisfy the geometric conditions required by the Ambrosetti-Rabinowitz theorem [4]. We use a geometrical construction of subsets to overcome this difficulty. Let us give a definition and recall the mountain pass theorem of Ambrosetti and Rabinowitz.

Definition 3.1. Let $E$ be a Banach space. A subset $A$ of $E$ is said to be symmetric if $u \in E$ implies $-u \in E$. For a closed symmetric set $A$ that does not contain the origin, we define the genus $\gamma(A)$ of $A$ by the smallest integer $k$ such that there exists an odd continuous mapping from $A$ to $\mathbb{R}^{k} \backslash\{0\}$. If there does not exist such a $k$, we define $\gamma(A)=+\infty$. We set $\gamma(\emptyset)=0$. Let $\Gamma_{k}$ denote the family of closed symmetric subsets $A$ of $E$ such that $0 \notin A$ and $\gamma(A) \geqslant k$.

ThEOREM 3.2 (Ambrosetti and Rabinowitz [4]). Let $E$ be an infinite-dimensional Banach space and let $I \in C^{1}(E, \mathbb{R})$ satisfy the following conditions.
(1) I is even, bounded from below, $I(0)=0$ and I satisfies the PS condition.
(2) For each $k \in \mathbb{N}$ there exists $A_{k} \in \Gamma_{k}$ such that

$$
\sup _{u \in A_{k}} I(u)<0
$$

Under assumptions (1) and (2) we define $c_{k}$ by

$$
c_{k}=\inf _{A \in \Gamma_{k}} \sup _{u \in A} I(u)
$$

Then each $c_{k}$ is a critical value of $I, c_{k} \leqslant c_{k+1}<0$ for $k \in \mathbb{N}$ and $\left(c_{k}\right)$ converges to zero. Moreover, if $c_{k}=c_{k+1}=\cdots=c_{k+p}=c$, then $\gamma\left(K_{c}\right) \geqslant p+1$. The critical set $K_{c}$ is defined by

$$
K_{c}=\left\{u \in E ; I^{\prime}(u)=0, I(u)=c\right\}
$$

## 4. Proof of theorem 2.1

In this section we consider the case in which $a$ is bounded and we define $I$ on the function space $E=H^{1}\left(\mathbb{R}^{N}\right) \cap L^{q+1}\left(\mathbb{R}^{N}\right)$.

Lemma 4.1. Assume that $\left(A_{1}\right)$ and $\left(V_{1}\right)$ hold. Then any PS sequence of $I$ is bounded in $E$.

Proof. By standard arguments, $I$ is well defined and of class $C^{1}$ on $E$.
Let $\left(u_{n}\right)$ be a PS sequence of $I$. There then exists $\alpha>0$ such that $I\left(u_{n}\right) \leqslant \alpha$. Applying Hölder's inequality and conditions $\left(\mathrm{A}_{1}\right)$ and $\left(\mathrm{V}_{1}\right)$, we have

$$
\begin{align*}
\alpha \geqslant I\left(u_{n}\right)= & \frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|\nabla u_{n}\right|^{2}(x)+V(x) u_{n}(x)^{2}\right) \mathrm{d} x-\frac{1}{q+1} \int_{\mathbb{R}^{N}} a(x)\left|u_{n}(x)\right|^{q+1} \mathrm{~d} x \\
\geqslant & \frac{1}{2} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}(x)\right|^{2} \mathrm{~d} x-\frac{1}{2} \int_{\mathbb{R}^{N}} V^{-}(x) u_{n}(x)^{2} \mathrm{~d} x \\
& -\frac{1}{q+1} \int_{\mathbb{R}^{N}} a^{+}(x)\left|u_{n}\right|^{q+1}(x) \mathrm{d} x+\frac{1}{q+1} \int_{\mathbb{R}^{N}} a^{-}(x)\left|u_{n}\right|^{q+1}(x) \mathrm{d} x \\
\geqslant & \frac{1}{2} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2}(x) \mathrm{d} x-\frac{\left\|V^{-}\right\|_{N / 2}}{2 s}\left\|\nabla u_{n}\right\|_{2}^{2} \\
& -\frac{1}{q+1} \int_{\mathbb{R}^{N}} a^{+}(x)\left|u_{n}\right|^{q+1}(x) \mathrm{d} x . \tag{4.1}
\end{align*}
$$

By $\left(\mathrm{A}_{1}\right)$, there exists $R>0$ such that

$$
\begin{equation*}
-\|a\|_{\infty} \leqslant a(x) \leqslant \frac{a_{\infty}}{2}<0 \quad \forall|x| \geqslant R \quad \text { and } \quad a^{+} \in L^{m}\left(\mathbb{R}^{N}\right) \quad \forall 1 \leqslant m \leqslant+\infty \tag{4.2}
\end{equation*}
$$

Combining (4.1) and (4.2), we infer that

$$
\begin{aligned}
\alpha \geqslant I\left(u_{n}\right) & \geqslant \frac{1}{2}\left\|\nabla u_{n}\right\|_{2}^{2}-\frac{\left\|V^{-}\right\|_{N / 2}}{2 s}\left\|\nabla u_{n}\right\|_{2}^{2}-s^{(-q-1) / 2}\left\|a^{+}\right\|_{2^{*} /\left(2^{*}-(q+1)\right)}\left\|\nabla u_{n}\right\|_{2}^{q+1} \\
& \geqslant\left(\frac{1}{2}-\frac{\left\|V^{-}\right\|_{N / 2}}{2 s}\right)\left\|\nabla u_{n}\right\|_{2}^{2}-s^{(-q-1) / 2}\left\|a^{+}\right\|_{2^{*} /\left(2^{*}-(q+1)\right)}\left\|\nabla u_{n}\right\|_{2}^{q+1}
\end{aligned}
$$

and hence there exists $\beta>0$ such that

$$
\begin{equation*}
\left\|\nabla u_{n}\right\|_{2} \leqslant \beta \quad \forall n \in \mathbb{N} \tag{4.3}
\end{equation*}
$$

On the other hand, there exists $c>0$ such that

$$
\begin{aligned}
c+\frac{\left\|u_{n}\right\|}{2} \geqslant & -\frac{1}{2}\left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle+I\left(u_{n}\right) \\
= & \left(\frac{1}{2}-\frac{1}{q+1}\right) \int_{\mathbb{R}^{N}} a(x)\left|u_{n}\right|^{q+1}(x) \mathrm{d} x \\
= & \left(\frac{1}{q+1}-\frac{1}{2}\right) \int_{\mathbb{R}^{N}} a^{-}(x)\left|u_{n}\right|^{q+1} \mathrm{~d} x \\
& -\left(\frac{1}{q+1}-\frac{1}{2}\right) \int_{\mathbb{R}^{N}} a^{+}(x)\left|u_{n}\right|^{q+1} \mathrm{~d} x \\
= & \left(\frac{1}{q+1}-\frac{1}{2}\right) \int_{\mathbb{R}^{N}}\left(a^{-}(x)+\chi_{B_{R}}(x)\right)\left|u_{n}\right|^{q+1} \mathrm{~d} x \\
& -\left(\frac{1}{q+1}-\frac{1}{2}\right) \int_{\mathbb{R}^{N}}\left(a^{+}(x)+\chi_{B_{R}}(x)\right)\left|u_{n}\right|^{q+1} \mathrm{~d} x \\
\geqslant & \left(\frac{1}{q+1}-\frac{1}{2}\right) \min \left\{\frac{-a_{\infty}}{2}, 1\right\} \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{q+1}(x) \mathrm{d} x \\
& -s^{(-q-1) / 2}\left(\frac{1}{q+1}-\frac{1}{2}\right)\left\|a^{+}+\chi_{B_{R}}\right\|_{2^{*} /\left(2^{*}-(q+1)\right)}\left\|\nabla u_{n}\right\|_{2}^{q+1} .
\end{aligned}
$$

Thus, there is a constant $c>0$ such that

$$
\int_{\mathbb{R}^{N}}\left|u_{n}\right|^{q+1} \mathrm{~d} x \leqslant c\left(\left\|\nabla u_{n}\right\|_{2}+\left\|u_{n}\right\|_{q+1}+\left\|a^{+}+\chi_{R}\right\|_{2^{*} /\left(2^{*}-(q+1)\right)}\left\|\nabla u_{n}\right\|_{2}^{q+1}\right) .
$$

Relation (4.3) yields

$$
\begin{equation*}
\left\|u_{n}\right\|_{q+1}^{q+1} \leqslant c+c\left\|u_{n}\right\|_{q+1} \quad \text { for all } n \in \mathbb{N} . \tag{4.4}
\end{equation*}
$$

Combining (4.3) and (4.4), we get

$$
\left\|u_{n}\right\| \leqslant c \quad \forall n \in \mathbb{N}
$$

The proof is complete.
We need the following lemma to prove that the PS condition is satisfied for $I$ on $E$.

Lemma 4.2. There exists a constant $c>0$ such that for all real numbers $x, y$,

$$
\begin{equation*}
\| x+\left.y\right|^{q+1}-|x|^{q+1}-\left.|y|^{q+1}|\leqslant c| x\right|^{q}|y| . \tag{4.5}
\end{equation*}
$$

Proof. If $x=0$, the inequality (4.5) is trivial.
Suppose that $x \neq 0$. We consider the continuous function $f$ defined on $\mathbb{R} \backslash\{0\}$ by

$$
f(t)=\frac{|1+t|^{q+1}-|t|^{q+1}-1}{|t|}
$$

Then $\lim _{|t| \rightarrow+\infty} f(t)=0$ and $\lim _{t \rightarrow 0 \pm} f(t)= \pm(q+1)$. Thus, there exists a constant $c>0$ such that $|f(t)| \leqslant c$ for all $t \in \mathbb{R} \backslash\{0\}$. In particular, $|f(y / x)| \leqslant c$, so

$$
\left|\left|1+\frac{y}{x}\right|^{q+1}-\left|\frac{y}{x}\right|^{q+1}-1\right| \leqslant c\left|\frac{y}{x}\right|
$$

Multiplying by $|x|^{q+1}$, we obtain the desired result.
Lemma 4.3. Assume $\left(A_{1}\right)$ and $\left(V_{1}\right)$ hold. Then I satisfies the $P S$ condition on $E$.
Proof. Let $\left(u_{n}\right)$ be a PS sequence. By lemma 4.1, $\left(u_{n}\right)$ is bounded in $E$. There then exists a subsequence $u_{n} \rightharpoonup u$ in $E, u_{n} \rightarrow u$ in $L_{\mathrm{loc}}^{p}\left(\mathbb{R}^{N}\right)$ for all $1 \leqslant p \leqslant 2^{*}$ and $u_{n} \rightarrow u$ almost everywhere (a.e.) in $\mathbb{R}^{N}$.

Fix $\varphi \in D\left(\mathbb{R}^{N}\right)$. By the weak convergence of $\left(u_{n}\right)$ to $u$, we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} \nabla u_{n} \nabla \varphi(x)+V(x) u_{n} \varphi(x) \mathrm{d} x \rightarrow \int_{\mathbb{R}^{N}} \nabla u \nabla \varphi+V(x) u \varphi(x) \mathrm{d} x \tag{4.6}
\end{equation*}
$$

By compactness Sobolev embedding, $u_{n} \rightarrow u$ in $L^{q+1}(\operatorname{supp}(\varphi))$, and hence there exists a function $h \in L^{q+1}\left(\mathbb{R}^{N}\right)$ such that

$$
a(x)\left|u_{n}\right|^{q-1} u_{n} \varphi \rightarrow a(x)|u|^{q-1} u \varphi \quad \text { a.e. in } \mathbb{R}^{N}
$$

and

$$
|a|\left|u_{n}\right|^{q}|\varphi| \leqslant\|a\|_{\infty}|h||\varphi| \quad \text { in } \mathbb{R}^{N}
$$

Using the Lebesgue dominated convergence theorem, we deduce that

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} a(x)\left|u_{n}\right|^{q-1} u_{n} \varphi(x) \mathrm{d} x \rightarrow \int_{\mathbb{R}^{N}} a(x)|u|^{q-1} u \varphi(x) \mathrm{d} x . \tag{4.7}
\end{equation*}
$$

Combining (4.6) and (4.7), we obtain

$$
0=\lim _{n \rightarrow+\infty}\left\langle I^{\prime}\left(u_{n}\right), \varphi\right\rangle=\left\langle I^{\prime}(u), \varphi\right\rangle \quad \forall \varphi \in D\left(\mathbb{R}^{N}\right)
$$

Then,

$$
\begin{equation*}
\left\langle I^{\prime}(u), u\right\rangle=0 \tag{4.8}
\end{equation*}
$$

Since $u_{n} \rightharpoonup u$ in $E$, we have $\|u\| \leqslant \liminf _{n \rightarrow+\infty}\left\|u_{n}\right\|=\lim _{n \rightarrow+\infty}\left\|u_{n}\right\|$. We distinguish two cases.

CASE 1 (compactness). $\|u\|=\lim _{n \rightarrow+\infty}\left\|u_{n}\right\|$, so

$$
\limsup _{n \rightarrow+\infty}\left\|u_{n}\right\|_{q+1} \leqslant\|u\|_{q+1}+\|\nabla u\|_{2}-\liminf _{n \rightarrow+\infty}\left\|\nabla u_{n}\right\|_{2}
$$

Since

$$
\|\nabla u\|_{2} \leqslant \liminf _{n \rightarrow+\infty}\left\|\nabla u_{n}\right\|_{2}, \quad\|u\|_{q+1} \leqslant \liminf _{n \rightarrow+\infty}\left\|u_{n}\right\|_{q+1}
$$

we obtain

$$
\|u\|_{q+1} \leqslant \liminf _{n \rightarrow+\infty}\left\|u_{n}\right\|_{q+1} \leqslant \limsup _{n \rightarrow+\infty}\left\|u_{n}\right\|_{q+1} \leqslant\|u\|_{q+1}
$$

and thus

$$
\begin{aligned}
u_{n} & \rightarrow u \text { a.e. in } \mathbb{R}^{N}, \\
\left\|u_{n}\right\|_{q+1} & \rightarrow\|u\|_{q+1} .
\end{aligned}
$$

By the Brezis-Lied lemma [13], we infer that

$$
\begin{equation*}
u_{n} \rightarrow u \quad \text { in } L^{q+1}\left(\mathbb{R}^{N}\right) . \tag{4.9}
\end{equation*}
$$

Therefore, $\left\|\nabla u_{n}\right\|_{2} \rightarrow\|\nabla u\|_{2}$. On the other hand,

$$
\int_{\mathbb{R}^{N}}\left|\nabla u_{n}-\nabla u\right|^{2} \mathrm{~d} x=\int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2} \mathrm{~d} x+\int_{\mathbb{R}^{N}}|\nabla u|^{2} \mathrm{~d} x-2 \int_{\mathbb{R}^{N}} \nabla u_{n} \nabla u \mathrm{~d} x,
$$

and hence

$$
\int_{\mathbb{R}^{N}} \nabla u_{n} \nabla u \mathrm{~d} x \rightarrow \int_{\mathbb{R}^{N}}|\nabla u|^{2} \mathrm{~d} x .
$$

Therefore,

$$
\begin{equation*}
\left\|\nabla u_{n}-\nabla u\right\|_{2} \rightarrow 0 . \tag{4.10}
\end{equation*}
$$

Combining (4.9) and (4.10), we deduce that $u_{n} \rightarrow u$ in $E$ and the PS condition for $I$ is satisfied.

CASE 2 (dichotomy). $\|u\|<\lim _{n \rightarrow+\infty}\left\|u_{n}\right\|$. We prove that this case cannot occur. Set $v_{n}=u_{n}-u$.

Step 1 (There exists $\left(y_{n}\right) \subset \mathbb{R}^{N}$ such that $v_{n}\left(\cdot+y_{n}\right) \rightharpoonup v \neq 0$ in $E$ ). If not, for all $\left(y_{n}\right) \subset \mathbb{R}^{N}, v_{n}\left(\cdot+y_{n}\right) \rightharpoonup 0$ in $E$. Then

$$
\forall R>0 \quad \sup _{y \in \mathbb{R}^{N}} \int_{B(y, R)}\left|v_{n}\right|^{q+1}(x) \mathrm{d} x \rightarrow 0 .
$$

By [31, lemma I.1, p. 231],

$$
\begin{equation*}
v_{n} \rightarrow 0 \text { in } L^{p}\left(\mathbb{R}^{N}\right) \quad \forall q+1<p<2^{*} . \tag{4.11}
\end{equation*}
$$

On the other hand,

$$
\begin{align*}
& \left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
& \quad=\int_{\mathbb{R}^{N}}\left(\left|\nabla u_{n}\right|^{2}+V(x) u_{n}^{2}\right) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x)\left|u_{n}\right|^{q+1} \mathrm{~d} x \\
& =\int_{\mathbb{R}^{N}}\left(\left|\nabla v_{n}\right|^{2}+V(x) v_{n}^{2}\right) \mathrm{d} x+\int_{\mathbb{R}^{N}}|\nabla u|^{2} \mathrm{~d} x+\int_{\mathbb{R}^{N}}\left(V(x) u^{2}+2 \nabla v_{n} \nabla u\right) \mathrm{d} x \\
& \quad \quad \quad \int_{\mathbb{R}^{N}} 2 V(x) v_{n} u \mathrm{~d} x-\int_{\mathbb{R}^{N}} a(x)\left(\left|u_{n}\right|^{q+1}-|u|^{q+1}\right) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x)|u|^{q+1} \mathrm{~d} x . \tag{4.12}
\end{align*}
$$

By (4.5) in lemma 4.2 , we obtain

$$
\begin{aligned}
|a(x)|\left|\left|u_{n}\right|^{q+1}-|u|^{q+1}-\left|v_{n}\right|^{q+1}\right| & =|a(x)|| | v_{n}+\left.u\right|^{q+1}-|u|^{q+1}-\left|v_{n}\right|^{q+1} \mid \\
& \leqslant c|a(x)||u|^{q} v_{n} .
\end{aligned}
$$

Since $v_{n} \rightharpoonup 0$ in $E$, we deduce that

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} a(x)\left(\left|u_{n}\right|^{q+1}-|u|^{q+1}\right) \mathrm{d} x=\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} a(x)\left|v_{n}\right|^{q+1}(x) \mathrm{d} x \tag{4.13}
\end{equation*}
$$

Using Hölder's inequality in combination with (4.2) and (4.11), we obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left(a^{+}(x)+\chi_{B_{R}}(x)\right)\left|v_{n}\right|^{q+1} \mathrm{~d} x \leqslant\left\|a^{+}+\chi_{B_{R}}\right\|_{2 /(1-q)}\left\|v_{n}\right\|_{L^{2}(B(0, R))}^{q+1} \rightarrow 0 \tag{4.14}
\end{equation*}
$$

Passing to the limit in (4.12) and using (4.2), (4.11), (4.13) and (4.14), we obtain

$$
\begin{aligned}
0=\lim _{n \rightarrow+\infty}\left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle= & \left\langle I^{\prime}(u), u\right\rangle+\lim _{n \rightarrow+\infty}\left(\int_{\mathbb{R}^{N}}\left(\left|\nabla v_{n}\right|^{2}-a(x)\left|v_{n}\right|^{q+1}\right) \mathrm{d} x\right) \\
= & \left\langle I^{\prime}(u), u\right\rangle+\lim _{n \rightarrow+\infty}\left(\int_{\mathbb{R}^{N}}\left|\nabla v_{n}\right|^{2}+\left(a^{-}(x)+\chi_{B_{R}}\right)\left|v_{n}\right|^{q+1}\right) \\
& \quad-\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left(a^{+}(x)+\chi_{B_{R}}\right)\left|v_{n}\right|^{q+1} \mathrm{~d} x \\
= & \lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left|\nabla v_{n}\right|^{2}+\left(a^{-}(x)+\chi_{B_{R}}\right)\left|v_{n}\right|^{q+1} \mathrm{~d} x \\
\geqslant & \lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left(\left|\nabla v_{n}\right|^{2}+\min \left(\frac{1}{2}-a_{\infty}, 1\right)\left|v_{n}\right|^{q+1}\right) \mathrm{d} x \\
\geqslant & \lim _{n \rightarrow+\infty} \min \left(1, \min \left(\frac{1}{2}-a_{\infty}, 1\right)\right) \int_{\mathbb{R}^{N}}\left(\left|\nabla v_{n}\right|^{2}+\left|v_{n}\right|^{q+1}\right) \mathrm{d} x
\end{aligned}
$$

Then $v_{n} \rightarrow 0$ in $E$, which yields a contradiction.
Step 2. $\left(y_{n}\right)$ is not bounded. Indeed, suppose that $\left(y_{n}\right)$ is bounded and there exists a subsequence of $\left(y_{n}\right)$, also denoted by $\left(y_{n}\right)$, such that $y_{n} \rightarrow y_{0}$. Then, for all $\varphi \in D\left(\mathbb{R}^{N}\right)$,
$0=\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} \varphi\left(x-y_{n}\right) v_{n} \mathrm{~d} x=\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} \varphi(x) v_{n}\left(x+y_{n}\right) \mathrm{d} x=\int_{\mathbb{R}^{N}} \varphi(x) v(x) \mathrm{d} x$.
Hence, $v=0$ a.e. in $\mathbb{R}^{N}$, a contradiction.
STEP 3. We show that $v$ is a solution of the following problem:

$$
\left.\begin{array}{rl}
-\Delta u+v_{\infty} u & =a_{\infty}|u|^{q-1} u \quad \text { in } \mathbb{R}^{N} \\
u & \in E .
\end{array}\right\}
$$

We first prove that $\left(P_{\infty}\right)$ admits only the trivial solution. Thus, since $v$ solves $\left(P_{\infty}\right)$, we will obtain a contradiction.

Since $\left(y_{n}\right)$ is not bounded, $u_{n}\left(\cdot+y_{n}\right) \rightharpoonup v$ is in $E$. In fact, $u\left(\cdot+y_{n}\right) \rightharpoonup \psi \in E$, and hence

$$
0=\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{\mathbb{N}}} u\left(x+y_{n}\right) \varphi(x) \mathrm{d} x=\int_{\mathbb{R}^{\mathbb{N}}} \psi(x) \varphi(x) \mathrm{d} x \quad \forall \varphi \in D\left(\mathbb{R}^{N}\right)
$$

It follows that $\psi=0$ a.e. Therefore,

$$
\begin{equation*}
u_{n}\left(\cdot+y_{n}\right) \rightharpoonup v \quad \text { in } E \tag{4.15}
\end{equation*}
$$

Let $\varphi \in D\left(\mathbb{R}^{N}\right)$. We have

$$
\begin{aligned}
\left\langle I^{\prime}\left(u_{n}\right), \varphi\left(\cdot-y_{n}\right)\right\rangle= & \int_{\mathbb{R}^{N}}\left(\nabla u_{n} \nabla \varphi\left(x-y_{n}\right)+V(x) u_{n} \varphi\left(x-y_{n}\right)\right) \mathrm{d} x \\
& \quad-\int_{\mathbb{R}^{N}} a(x)\left|u_{n}\right|^{q-1} u_{n} \varphi\left(x-y_{n}\right) \mathrm{d} x \\
= & \int_{\mathbb{R}^{N}} \nabla u_{n}\left(x+y_{n}\right) \nabla \varphi(x)+V\left(x+y_{n}\right) u_{n}\left(x+y_{n}\right) \varphi(x) \mathrm{d} x \\
& -\int_{\mathbb{R} N} a\left(x+y_{n}\right)\left|u_{n}\right|^{q-1}\left(x+y_{n}\right) u_{n}\left(x+y_{n}\right) \varphi(x) \mathrm{d} x .
\end{aligned}
$$

Relation (4.15) yields

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} \nabla u_{n}\left(x+y_{n}\right) \nabla \varphi(x) \mathrm{d} x \rightarrow \int_{\mathbb{R}^{N}} \nabla v(x) \nabla \varphi(x) \mathrm{d} x . \tag{4.16}
\end{equation*}
$$

Since $\left(u_{n}\left(\cdot+y_{n}\right)\right)$ is bounded in $E, u_{n}\left(\cdot+y_{n}\right) \rightarrow v$ in $L_{\mathrm{loc}}^{p}\left(\mathbb{R}^{N}\right)$ for all $1 \leqslant p \leqslant 2^{*}$ (up to a subsequence), $u_{n}\left(x+y_{n}\right) \rightarrow v$ a.e. in $\mathbb{R}^{N}$ and there exists $K \in L^{p}\left(\mathbb{R}^{N}\right)$ such that $\varphi\left|u_{n}\left(\cdot+y_{n}\right)\right| \leqslant|K|$ in $\mathbb{R}^{N}, 1 \leqslant p \leqslant 2^{*}$. Then, by $\left(\mathrm{V}_{1}\right)$, we obtain

$$
\begin{aligned}
V\left(x+y_{n}\right) u_{n}\left(x+y_{n}\right) \varphi & \rightarrow v_{\infty} v \varphi \quad \text { a.e. in } \mathbb{R}^{N} \\
\left|V\left(x+y_{n}\right) u_{n}\left(x+y_{n}\right) \varphi\right| & \leqslant\|V\|_{\infty}|K||\varphi| \in L^{1}\left(\mathbb{R}^{N}\right) .
\end{aligned}
$$

Applying Lebesgue's dominated convergence theorem, we deduce that

$$
\begin{equation*}
\int_{\mathbb{R}^{N}} V\left(x+y_{n}\right) u_{n}\left(x+y_{n}\right) \varphi(x) \mathrm{d} x \rightarrow v_{\infty} \int_{\mathbb{R}^{N}} v(x) \varphi(x) \mathrm{d} x \tag{4.17}
\end{equation*}
$$

From hypothesis $\left(\mathrm{A}_{1}\right)$, we find

$$
\begin{gathered}
a\left(x+y_{n}\right)\left|u_{n}\left(x+y_{n}\right)\right|^{q-1} u_{n}\left(x+y_{n}\right) \varphi \rightarrow a_{\infty}|v|^{q-1} v \varphi \quad \text { a.e. in } \mathbb{R}^{N}, \\
\left|a\left(x+y_{n}\right)\right|\left|u_{n}\left(x+y_{n}\right)\right|^{q}|\varphi| \leqslant\|a\|_{\infty}|K|^{q}|\varphi| \in L^{1}\left(\mathbb{R}^{N}\right)
\end{gathered}
$$

Next, by Lebesgue's dominated convergence theorem, we obtain

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}} a\left(x+y_{n}\right)\left|u_{n}\right|^{q-1}\left(x+y_{n}\right) u_{n}\left(x+y_{n}\right) \mathrm{d} x=a_{\infty} \int_{\mathbb{R}^{N}}|v|^{q-1} v \varphi(x) \mathrm{d} x \tag{4.18}
\end{equation*}
$$

Combining (4.16)-(4.18), we deduce that for all $\varphi \in D\left(\mathbb{R}^{N}\right)$,

$$
\begin{aligned}
0 & =\lim _{n \rightarrow+\infty}\left\langle I^{\prime}\left(u_{n}\right), \varphi\left(\cdot-y_{n}\right)\right\rangle \\
& =\int_{\mathbb{R}^{N}}\left(\nabla v(x) \nabla \varphi(x)+v_{\infty} v \varphi\right) \mathrm{d} x-a_{\infty} \int_{\mathbb{R}^{N}}|v|^{q-1} v \varphi(x) \mathrm{d} x .
\end{aligned}
$$

Thus, $v$ is a weak solution of $\left(P_{\infty}\right)$, and hence $v=0$, which yields a contradiction. From steps 1, 2, and 3, we conclude that the dichotomy does not occur. The proof is complete.

Lemma 4.4. Assume that $\left(A_{1}\right)$ and ( $V_{1}$ ) are fulfilled. Then, for each $k \in \mathbb{N}$, there exists $A_{k} \in \Gamma_{k}$ such that $\sup _{u \in A_{k}} I(u)<0$.

Proof. We use some ideas developed in [27].
Let $R_{0}$ and $y_{0}$ be fixed by assumption $\left(\mathrm{A}_{1}\right)$ and consider the cube

$$
D\left(R_{0}\right)=\left\{\left(x_{1}, \ldots, x_{N}\right) \in \mathbb{R}^{N}:\left|x_{i}-y_{i}\right|<R_{0}, 1 \leqslant i \leqslant N\right\}
$$

Fix $k \in \mathbb{N}$ arbitrarily. Let $n \in \mathbb{N}$ be the smallest integer such that $n^{N} \geqslant k$. We divide $D\left(R_{0}\right)$ equally into $n^{N}$ small cubes (denote them by $D_{i}$ with $1 \leqslant i \leqslant n^{N}$ ) with planes parallel to each face of $D\left(R_{0}\right)$. The edge of $D_{i}$ has the length of $a=R_{0} / n$. We construct new cubes $E_{i}$ in $D_{i}$ such that $E_{i}$ has the same centre as that of $D_{i}$. The faces of $E_{i}$ and $D_{i}$ are parallel and the edge of $E_{i}$ has the length $\frac{1}{2} a$. Thus, we can construct a function $\psi_{i}, 1 \leqslant i \leqslant k$, such that

$$
\begin{aligned}
& \sup \left(\psi_{i}\right) \subset D_{i}, \quad \operatorname{supp}\left(\psi_{i}\right) \cap \operatorname{supp}\left(\psi_{j}\right)=\emptyset \quad(i \neq j) \\
& \psi_{i}(x)=1 \quad \text { for } x \in E_{i}, \quad 0 \leqslant \psi_{i}(x) \leqslant 1 \quad \forall x \in \mathbb{R}^{N}
\end{aligned}
$$

We define

$$
\begin{align*}
S^{k-1} & =\left\{\left(t_{1}, \ldots, t_{k}\right) \in \mathbb{R}^{k}: \max _{1 \leqslant i \leqslant k}\left|t_{i}\right|=1\right\}  \tag{4.19}\\
W_{k} & =\left\{\sum_{i=1}^{k} t_{i} \psi_{i}(x):\left(t_{1}, \ldots, t_{k}\right) \in S^{k-1}\right\} \subset E \tag{4.20}
\end{align*}
$$

Since the mapping $\left(t_{1}, \ldots, t_{k}\right) \rightarrow \sum_{i=1}^{k} t_{i} \psi_{i}$ from $S^{k-1}$ to $W_{k}$ is odd and homeomorphic, we have $\gamma\left(W_{k}\right)=\gamma\left(S^{k-1}\right)=k$. But $W_{k}$ is compact in $E$, and thus there is a constant $\alpha_{k}>0$ such that

$$
\|u\|^{2} \leqslant \alpha_{k} \quad \text { for all } u \in W_{k}
$$

We recall the inequality

$$
\begin{equation*}
\|u\|_{2} \leqslant c\|\nabla u\|_{2}^{r}\|u\|_{q+1}^{1-r} \leqslant c\|u\| \tag{4.21}
\end{equation*}
$$

with $r=2^{*}(q-1) / 2\left(2^{*}-q-1\right)$. Then there is a constant $c_{k}>0$ such that

$$
\|u\|_{2}^{2} \leqslant c_{k} \quad \text { for all } u \in W_{k}
$$

Let $z>0$ and $u=\sum_{i=1}^{k} t_{i} \psi_{i}(x) \in W_{k}$. We have

$$
\begin{equation*}
I(z u) \leqslant \frac{z^{2}}{2} \alpha_{k}+z^{2} \frac{\|V\|_{\infty}}{2} c_{k}-\frac{1}{q+1} \sum_{i=1}^{k} \int_{D_{i}} a(x)\left|z t_{i} \psi_{i}\right|^{q+1} \mathrm{~d} x \tag{4.22}
\end{equation*}
$$

By (4.19), there exists $j \in[1, k]$ such that $\left|t_{j}\right|=1$ and $\left|t_{i}\right| \leqslant 1$ for $i \neq j$. Then

$$
\begin{align*}
\sum_{i=1}^{k} \int_{D_{i}} a(x)\left|z t_{i} \psi_{i}\right|^{q+1} \mathrm{~d} x= & \int_{E_{j}} a(x)\left|z t_{j} \psi_{j}\right|^{q+1} \mathrm{~d} x+\int_{D_{j} \backslash E_{j}} a(x)\left|z t_{j} \psi_{j}(x)\right|^{q+1} \mathrm{~d} x \\
& +\sum_{i \neq j} \int_{D_{i}} a(x)\left|z t_{i} \psi_{i}\right|^{q+1} \mathrm{~d} x \tag{4.23}
\end{align*}
$$

Since $\psi_{j}(x)=1$ for $x \in E_{j}$ and $\left|t_{j}\right|=1$, we have

$$
\begin{equation*}
\int_{E_{j}} a(x)\left|z t_{j} \psi_{j}\right|^{q+1} \mathrm{~d} x=|z|^{q+1} \int_{E_{j}} a(x) \mathrm{d} x \tag{4.24}
\end{equation*}
$$

On the other hand, by $\left(\mathrm{A}_{1}\right)$,

$$
\begin{equation*}
\int_{D_{j} \backslash E_{j}} a(x)\left|z t_{j} \psi_{j}\right|^{q+1} \mathrm{~d} x+\sum_{i \neq j} \int_{D_{i}} a(x)\left|z t_{i} \psi_{i}\right|^{q+1} \mathrm{~d} x \geqslant 0 \tag{4.25}
\end{equation*}
$$

Relations (4.22)-(4.25) yield

$$
\begin{equation*}
\frac{I(z u)}{z^{2}} \leqslant \frac{\alpha_{k}}{2}+\frac{\|V\|_{\infty}}{2} c_{k}-\frac{|z|^{q+1}}{z^{2}} \inf _{1 \leqslant i \leqslant k}\left(\int_{E_{i}} a(x) \mathrm{d} x\right) . \tag{4.26}
\end{equation*}
$$

By (4.26), we conclude that

$$
\lim _{z \rightarrow 0} \sup _{u \in W_{k}} \frac{I(z u)}{z^{2}}=-\infty
$$

We fix $z$ small enough such that

$$
\sup \left\{I(u), u \in A_{k}\right\}<0, \quad \text { where } A_{k}=z W_{k} \in \Gamma_{k}
$$

This concludes the proof.
Lemma 4.5. Assume that $\left(A_{1}\right)$ and $\left(V_{1}\right)$ hold. Then $I$ is bounded from below.
Proof. By $\left(\mathrm{A}_{1}\right)$, we obtain

$$
\begin{equation*}
a^{+} \in L^{p}\left(\mathbb{R}^{N}\right) \quad \text { for all } 1 \leqslant p \leqslant+\infty \tag{4.27}
\end{equation*}
$$

Then

$$
\begin{aligned}
I(u) & =\frac{1}{2} \int_{\mathbb{R}^{N}}\left(|\nabla u|^{2}+V(x) u^{2}\right) \mathrm{d} x-\frac{1}{q+1} \int_{\mathbb{R}^{N}} a(x)|u|^{q+1} \mathrm{~d} x \\
& \geqslant \frac{1}{2} \int_{\mathbb{R}^{N}}\left(|\nabla u|^{2}-V^{-}(x) u^{2}\right) \mathrm{d} x-\frac{1}{q+1} \int_{\mathbb{R}^{N}} a^{+}(x)|u|^{q+1} \mathrm{~d} x \\
& \geqslant\left(\frac{1}{2}-\frac{\left\|V^{-}\right\|_{N / 2}}{2 s}\right)\|\nabla u\|_{2}^{2}-\frac{\left\|a^{+}\right\|_{2^{*} /\left(2^{*}-q-1\right)}}{s^{(q+1) / 2}}\|\nabla u\|_{2}^{q+1} .
\end{aligned}
$$

In view of $\left(\mathrm{V}_{1}\right)$, we conclude the proof.
Proof of theorem 2.1 concluded. We have $I(0)=0$ and $I$ is even. Combining lemmas $4.3,4.4$ and 4.5 , we deduce that theorem $3.2(1)$ and (2) are satisfied. Thus, there exists a sequence $\left(u_{n}\right) \subset E$ such that $I\left(u_{n}\right)<0, I^{\prime}\left(u_{n}\right)=0$ and $I\left(u_{n}\right) \rightarrow 0$ for all $n \geqslant 0$, and hence $u_{n}$ is a weak solution of (2.1).

By $\left(\mathrm{V}_{1}\right)$, we deduce that

$$
\begin{aligned}
\frac{1}{q+1}\left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle-I\left(u_{n}\right) & =\left(\frac{1}{q+1}-\frac{1}{2}\right) \int_{\mathbb{R}^{N}}\left(\left|\nabla u_{n}\right|^{2}+V(x) u_{n}^{2}\right) \mathrm{d} x \\
& \geqslant\left(\frac{1}{q+1}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2 s}\left\|V^{-}\right\|_{N / 2}\right)\left\|\nabla u_{n}\right\|_{2}^{2}
\end{aligned}
$$

It follows that

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2} \mathrm{~d} x=0 \tag{4.28}
\end{equation*}
$$

On the other hand, by Hölder's inequality, (4.2) and (4.28), we have

$$
\begin{aligned}
0 & =\lim _{n \rightarrow+\infty}\left(I\left(u_{n}\right)-\frac{1}{2}\left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle\right) \\
& =\lim _{n \rightarrow+\infty}\left(\frac{1}{2}-\frac{1}{q+1}\right) \int_{\mathbb{R}^{N}} a(x)\left|u_{n}\right|^{q+1} \\
& =\left(\frac{1}{q+1}-\frac{1}{2}\right) \lim _{n \rightarrow+\infty}\left(\int_{\mathbb{R}^{N}}\left(a^{-}(x)+\chi_{B_{R}}\right)\left|u_{n}\right|^{q+1} \mathrm{~d} x\right. \\
& \left.\quad-\int_{\mathbb{R}^{N}}\left(a^{+}(x)+\chi_{B_{R}}\right)\left|u_{n}\right|^{q+1}\right) \\
& \geqslant\left(\frac{1}{q+1}-\frac{1}{2}\right) \lim _{n \rightarrow+\infty}\left(\int_{\mathbb{R}^{N}}\left(a^{-}(x)+\chi_{B_{R}}\right)\left|u_{n}\right|^{q+1} \mathrm{~d} x\right. \\
& \left.\quad-\frac{\left\|a^{+}+\chi_{B_{R}}\right\|^{2^{*} /\left(2^{*}-q-1\right)}}{s^{(q+1) / 2}}\left\|\nabla u_{n}\right\|_{2}^{q+1}\right) \\
& =\left(\frac{1}{q+1}-\frac{1}{2}\right) \lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left(a^{-}(x)+\chi_{B_{R}}\right)\left|u_{n}\right|^{q+1} \mathrm{~d} x \\
& \geqslant\left(\frac{1}{q+1}-\frac{1}{2}\right) \min \left(\frac{-a_{\infty}}{2}, 1\right) \lim _{n \rightarrow+\infty}\left\|u_{n}\right\|_{q+1}^{q+1} .
\end{aligned}
$$

This shows that

$$
\begin{equation*}
\lim _{n \rightarrow+\infty} \int_{\mathbb{R}^{N}}\left|u_{n}\right|^{q+1} \mathrm{~d} x=0 \tag{4.29}
\end{equation*}
$$

and hence $\lim _{n \rightarrow+\infty} u_{n}=0$ in $E$. This concludes the proof.

## 5. Proof of theorem 2.2

In this section we define $I$ and $\varphi$ on $X$. We use standard arguments based on the fact that $I^{\prime}$ is a sequentially compact operator in order to prove that $I$ satisfies the PS condition. We then deduce that (2.1) admits infinitely many non-trivial solutions in $X$.

To prove theorem 2.2 , we need the following auxiliary results.
Lemma 5.1. Assume that $\left(A_{2}\right),\left(V_{2}\right)$ and $\left(G_{1}\right)$ are satisfied. Then $\varphi^{\prime}$ is a sequentially compact operator on $X$.

Proof. By standard arguments, the functionals $I$ and $\varphi$ are well defined and of class $C^{1}$ on $X$.

Let $\left(u_{n}\right) \subset X$ be a bounded sequence. Then, for all $h \in X$,

$$
\left\langle\varphi^{\prime}\left(u_{n}\right)-\varphi^{\prime}(u), h\right\rangle=\int_{\mathbb{R}^{N}}\left[V(x)\left(u_{n}-u\right)-a(x)\left(g\left(u_{n}\right)-g(u)\right)\right] h(x) \mathrm{d} x
$$

Let $R>0$ and $h \in X$ be such that $\|h\|=1$. We have

$$
\left\langle\varphi^{\prime}\left(u_{n}\right)-\varphi^{\prime}(u), h\right\rangle=J_{1}(n, h, R)+J_{2}(n, h, R)
$$

where

$$
\begin{aligned}
& J_{1}(n, h, R)=\int_{B_{R}}\left[V(x)\left(u_{n}-u\right)-a(x)\left(g\left(u_{n}\right)-g(u)\right)\right] h(x) \mathrm{d} x \\
& J_{2}(n, h, R)=\int_{B_{R}^{c}}\left[V(x)\left(u_{n}-u\right)-a(x)\left(g\left(u_{n}\right)-g(u)\right)\right] h(x) \mathrm{d} x
\end{aligned}
$$

By Hölder's inequality, $\left(\mathrm{V}_{2}\right),\left(\mathrm{A}_{2}\right)$ and $\left(\mathrm{G}_{1}\right)$, we obtain

$$
\begin{aligned}
& \left|J_{2}(n, h, R)\right| \\
& \leqslant \\
& \leqslant \int_{B_{R}^{c}}\left|V(x)\left(u_{n}-u\right) h(x)-a(x)\left(g\left(u_{n}\right)-g(u)\right) h(x)\right| \mathrm{d} x \\
& \leqslant \\
& \quad+c\left(\int_{B_{R}^{c}}|V(x)|^{N / 2} \mathrm{~d} x\right)^{2 / N}\left(\int_{B_{R}^{c}}|a(x)|^{2^{*} /\left(2^{*}-(q+1)\right)} \mathrm{d} x\right)^{\left(2^{*}-(q+1)\right) / 2^{*}}\left(\int_{B_{R}^{c}}\left(\left|u_{n}(x)+u(x)\right| 2^{2^{*}} \mathrm{~d} x\right)^{1 / 2^{*}}\left(\int_{B_{R}^{c}}|h(x)|^{2^{*}} \mathrm{~d} x\right)^{q / 2^{*}}\right. \\
& \\
& \quad \times\left(\int_{B_{R}^{c}}|h(x)|^{2^{*}} \mathrm{~d} x\right)^{1 / 2^{*}} \\
& \leqslant \\
& \leqslant\left(\int_{B_{R}^{c}}|V(x)|^{N / 2} \mathrm{~d} x\right)^{2 / N}+\left(\int_{B_{R}^{c}}|a(x)|^{2^{*} /\left(2^{*}-(q+1)\right)} \mathrm{d} x\right)^{\left(2^{*}-(q+1)\right) / 2^{*}}
\end{aligned}
$$

The last expression can be made arbitrarily small by taking $R>0$ large enough.
For $J_{1}$, since $V \in L^{N / 2}\left(\mathbb{R}^{N}\right)$ and $a \in L^{2^{*} /\left(2^{*}-(q+1)\right)}\left(\mathbb{R}^{N}\right)$, we deduce that for all $\varepsilon>0$ there exists $\eta>0$ such that

$$
\left(\int_{K}|a(x)|^{2^{*} /\left(2^{*}-(q+1)\right)} \mathrm{d} x\right)^{\left(2^{*}-(q+1)\right) / 2^{*}}+\left(\int_{K}|V(x)|^{N / 2} \mathrm{~d} x\right)^{2 / N}<\varepsilon
$$

for all $K \subset B_{R}$ with $m(K)<\eta$ (see [20]). Moreover,

$$
\begin{aligned}
\int_{K} \mid V(x) & \left(u_{n}-u\right)-a(x)\left(g\left(u_{n}\right)-g(u)\right)| | h(x) \mid \mathrm{d} x \\
& \leqslant c\left(\int_{K}|V(x)|^{N / 2} \mathrm{~d} x\right)^{2 / N}+c\left(\int_{K}|a(x)|^{2^{*} /\left(2^{*}-(q+1)\right)} \mathrm{d} x\right)^{\left(2^{*}-(q+1)\right) / 2^{*}} \\
& \leqslant c \varepsilon
\end{aligned}
$$

where $c$ is independent of $n$ and $h$. By using the Vitali convergence theorem, we deduce that $J_{1}(n, h, R) \rightarrow 0$ as $n \rightarrow+\infty$ uniformly for $\|h\|=1$. We conclude that $\varphi^{\prime}\left(u_{n}\right) \rightarrow \varphi^{\prime}(u)$ strongly in $X^{\prime}$. The proof is complete.

Lemma 5.2. Assume that $\left(V_{2}\right),\left(A_{2}\right)$ and $\left(G_{1}\right)$ are satisfied. Then any $P S$ sequence of $I$ is bounded in $X$.

Proof. Let $\left(u_{n}\right) \subset X$ be a PS sequence. There then exists $\alpha>0$ such that $I\left(u_{n}\right) \leqslant$ $\alpha$. By Hölder's inequality and conditions $\left(\mathrm{A}_{2}\right),\left(\mathrm{V}_{2}\right)$ and $\left(\mathrm{G}_{1}\right)$, we have

$$
\begin{aligned}
\alpha \geqslant I\left(u_{n}\right) & \geqslant \frac{1}{2} \int_{\mathbb{R}^{N}}\left(\left|\nabla u_{n}\right|^{2}(x)-V^{-}(x) u_{n}(x)^{2}\right) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x) G\left(u_{n}(x)\right) \mathrm{d} x \\
& \geqslant\left(\frac{1}{2}-\frac{1}{2 s}\left\|V^{-}\right\|_{N / 2}\right)\left\|u_{n}\right\|_{X}^{2}-s^{(-q-1) / 2}\|a\|_{2^{*} /\left(2^{*}-(q+1)\right)}\left\|u_{n}\right\|_{X}^{q+1}
\end{aligned}
$$

Since $0<q<1$, the last inequality shows that $\left(u_{n}\right)$ is bounded in $X$. The proof is complete.

As a consequence, we obtain the following result.
Lemma 5.3. Assume that $\left(V_{2}\right),\left(A_{2}\right)$ and $\left(G_{1}\right)$ are satisfied. Then I satisfies the $P S$ condition in $X$.

Proof. Set

$$
\begin{aligned}
& F: D^{1,2}\left(\mathbb{R}^{N}\right) \rightarrow\left(D^{1,2}\left(\mathbb{R}^{N}\right)\right)^{\prime} \\
& u \mapsto F(u) \\
&\langle F(u), v\rangle=\int_{\mathbb{R}^{N}} \nabla u \nabla v \mathrm{~d} x \quad \forall v \in D^{1,2}\left(\mathbb{R}^{N}\right) .
\end{aligned}
$$

Then $F$ is an isomorphism. Let $\left(u_{n}\right)$ be a PS sequence of $I$; hence,

$$
\begin{equation*}
u_{n}=F^{-1}\left(\varphi^{\prime}\left(u_{n}\right)\right)+o(1) \tag{5.1}
\end{equation*}
$$

By lemma 5.2, ( $u_{n}$ ) is bounded in $X$. Since $\varphi^{\prime}$ is a compact operator and using (5.1), we deduce that $\left(u_{n}\right)$ is strongly convergent in $X$ (up to a subsequence).
Lemma 5.4. Assume that $\left(G_{1}\right),\left(V_{2}\right)$ and $\left(A_{2}\right)$ are satisfied. Then $I$ is bounded from below.

Proof. By $\left(\mathrm{G}_{1}\right),\left(\mathrm{V}_{2}\right)$ and $\left(\mathrm{A}_{2}\right)$, we have

$$
\begin{aligned}
I(u) & =\frac{1}{2} \int_{\mathbb{R}^{N}}\left(|\nabla u(x)|^{2}+V(x) u^{2}(x)\right) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x) G(u(x)) \mathrm{d} x \\
& \geqslant \frac{1}{2} \int_{\mathbb{R}^{N}}|\nabla u(x)|^{2} \mathrm{~d} x-\frac{1}{2} \int_{\mathbb{R}^{N}} V^{-}(x) u^{2}(x) \mathrm{d} x-\int_{\mathbb{R}^{N}} a(x) G(u(x)) \mathrm{d} x \\
& \geqslant \frac{1}{2} \int_{\mathbb{R}^{N}}|\nabla u(x)|^{2} \mathrm{~d} x-\frac{1}{2 s}\left\|V^{-}\right\|_{N / 2}\|u\|_{X}^{2}-s^{(-q-1) / 2}\|a\|_{2^{*} /\left(2^{*}-(q+1)\right)}\|u\|_{X}^{q+1} \\
& \geqslant\left(\frac{1}{2}-\frac{1}{2 s}\left\|V^{-}\right\|_{N / 2}\right)\|u\|_{X}^{2}-s^{(-q-1) / 2}\|a\|_{2^{*} /\left(2^{*}-(q+1)\right)}\|u\|_{X}^{q+1} .
\end{aligned}
$$

Since $1<q+1<2$, we deduce that $I$ is bounded from below. The proof is complete.

Next, we prove the geometric condition required by theorem 3.2.
Lemma 5.5. Assume that $\left(A_{2}\right),\left(V_{2}\right),\left(G_{1}\right),\left(G_{2}\right)$ and $\left(G_{3}\right)$ are satisfied. Then, for each $k \in \mathbb{N}$ there exists an $A_{k} \in \Gamma_{k}$ such that $\sup _{u \in A_{k}} I(u)<0$.

Proof. By using conditions $\left(\mathrm{G}_{2}\right)$ and $\left(\mathrm{G}_{3}\right)$, the proof is similar to that of lemma 4.4.

Proof of theorem 2.2 concluded. The energy functional $I$ is even and $I(0)=0$. By lemmas 5.3 and 5.4 , theorem 3.2(1) is satisfied. In view of lemma 5.5, theorem 3.2(2) is also satisfied. Thus, there exists a sequence $\left(u_{k}\right)$ such that $c_{k}=I\left(u_{k}\right)$ is a critical value of $I, c_{k}<0, c_{k} \rightarrow 0$ for all $k \geqslant 0$. This means that $\left(u_{k}\right)$ are weak solutions of (2.1) and $\left(u_{k}\right)$ is a PS sequence of $I$. Then, by lemma 5.2, $\left(u_{k}\right)$ is bounded.

REmARK 5.6. If $g(x)=|x|^{q-1} x, 0<q<1$, then $u_{n} \rightarrow 0$ in $X$. In fact, by $\left(\mathrm{V}_{2}\right)$, we have

$$
\begin{aligned}
0=\frac{1}{q+1}\left\langle I^{\prime}\left(u_{n}\right), u_{n}\right\rangle-I\left(u_{n}\right) & =\left(\frac{1}{q+1}-\frac{1}{2}\right) \int_{\mathbb{R}^{N}}\left(\left|\nabla u_{n}\right|^{2}+V(x)\left|u_{n}\right|^{2}\right) \mathrm{d} x \\
& \geqslant\left(\frac{1}{q+1}-\frac{1}{2}\right)\left(1-\frac{\left\|V^{-}\right\|_{N / 2}}{s}\right) \int_{\mathbb{R}^{N}}\left|\nabla u_{n}\right|^{2}(x) \mathrm{d} x
\end{aligned}
$$

Since $I^{\prime}\left(u_{n}\right)=0$ and $\lim _{n \rightarrow+\infty} I\left(u_{n}\right)=0$, we deduce that $u_{n} \rightarrow 0$ in $X$.

## 6. Proof of theorem 2.3

In this section we change the condition $\left(\mathrm{V}_{2}\right)$ to $\left(\mathrm{V}_{1}\right)$ and we suppose that $a$ satisfies $\left(\mathrm{A}_{3}\right)$. Under the last conditions, if the functional $I$ is not well defined on either $X$ or $E$, then we define it on the space $Y$. We first establish that $(Y,\langle\cdot\rangle)$ is a Hilbert space and it is embedded into $L^{p}\left(\mathbb{R}^{N}\right)$ for $2 \leqslant p \leqslant 2^{*}$. The proof of the following result relies on standard arguments and we will omit it.

Lemma 6.1. Assume that ( $V_{1}$ ) holds. Then

$$
u \rightarrow\left(\int_{\mathbb{R}^{N}}\left(|\nabla u(x)|^{2}+V^{+}(x) u^{2}(x)\right) \mathrm{d} x\right)^{1 / 2}
$$

defines a norm on $Y$, which is equivalent to the usual norm in $H^{1}\left(\mathbb{R}^{N}\right)$,

$$
\|u\|_{H^{1}}=\|\nabla u\|_{2}+\|u\|_{2} .
$$

By using lemma 6.1, the proof of theorem 2.3, with slight modifications, is similar to that of theorem 2.2.

Remark 6.2. If $g(x)=|x|^{q-1} x, 0<q<1$, then $u_{n} \rightarrow 0$ in $Y$.
REMARK 6.3. In theorems 2.1, 2.2 and 2.3 we can suppose that $u_{0}$ is a non-negative solution of (2.1), since

$$
I\left(u_{0}\right)=I\left(\left|u_{0}\right|\right)=c_{0}
$$

In such a case, $u_{0}$ is called a ground state for $I$.

## Acknowledgements

The authors are grateful to the anonymous referee and to the handling editor for the careful reading of the paper and for their constructive remarks. V.R. acknowledges support from grant no. CNCS PCE-47/2011.

## References

1 M. J. Ablowitz, B. Prinari and A. D. Trubatch. Discrete and continuous nonlinear Schrödinger systems (Cambridge University Press, 2004).
$2 \quad$ S. Adachi and K. Tanaka. Four positive solutions for the equation $-\Delta u+u=a(x) u^{p}+f(x)$ in $\mathbb{R}^{N}$. Calc. Var. PDEs 11 (2000), 63-95.
3 A. Ambrosetti and M. Badiale. The dual variational principle and elliptic problems with discontinuous nonlinearities. J. Math. Analysis Applic. 140 (1989), 363-373.
4 A. Ambrosetti and P. H. Rabinowitz. Dual variational methods in critical point theory and applications. J. Funct. Analysis 14 (1973), 349-381.
5 J. Avron, I. Herbst and B. Simon. Schrödinger operators with electromagnetic fields. III. Atoms in homogeneous magnetic field. Commun. Math. Phys. 79 (1981), 529-572.
6 A. Bahri and P.-L. Lions. On the existence of a positive solution of semilinear elliptic equations in unbounded domains. Annales Inst. H. Poincaré Analyse Non Linéaire 14 (1997), 365-413.

7 M. Balabane, J. Dolbeault and H. Ounaies. Nodal solutions for a sublinear elliptic equation. Nonlin. Analysis 52 (2003), 219-237.
8 T. Bartsch, Z. Liu and T. Weth. Sign changing solutions for superlinear Schrödinger equations. Commun. PDEs 29 (2004), 25-42.
9 M. Benrhouma. Study of multiplicity and uniqueness of solutions for a class of nonhomogeneous sublinear elliptic equations. Nonlin. Analysis 74 (2011), 2682-2694.
10 M. Benrhouma and H. Ounaies. Existence and uniqueness of positive solution for nonhomogeneous sublinear elliptic equation. J. Math. Analysis Applic. 358 (2009), 307-319.
11 M. Benrhouma and H. Ounaies. Existence of solutions for a perturbation sublinear elliptic equation in $\mathbb{R}^{N}$. Nonlin. Diff. Eqns Applic. 5 (2010), 647-662.
12 H. Brezis and S. Kamin. Sublinear elliptic equations in $\mathbb{R}^{n}$. Manuscr. Math. 74 (1992), 87-106.
13 H. Brezis and E. H. Lieb. A relation between pointwise convergence of functions and convergence of functionals. Proc. Am. Math. Soc. 88 (1983), 486-490.
14 H. Brezis and L. Oswald. Remarks on sublinear elliptic equations. Nonlin. Analysis 10 (1986), 55-64.

15 J. Byeon and Z. Q. Wang. Standing waves with a critical frequency for nonlinear Schrödinger equations. Arch. Ration. Mech. Analysis 165 (2002), 295-316.
16 J. Chabrowski and D. G. Costa. On a class of Schrödinger-type equations with indefinite weight functions. Commun. PDEs 33 (2008), 1368-1394.
17 D. G. Costa and H. Tehrani. Existence of positive solutions for a class of indefinite elliptic problems in $\mathbb{R}^{N}$. Calc. Var. PDEs 13 (2001), 159-189.
18 D. G. Costa and H. Tehrani. Unbounded perturbations of resonant Schrödinger equations. In Variational methods: open problems, recent progress, and numerical algorithms. Contemporary Mathematics, vol. 357, pp. 101-110 (Providence, RI: American Mathematical Society, 2004).
19 D. G. Costa, H. Tehrani and M. Ramos. Non-zero solutions for a Schrödinger equation with indefinite linear and nonlinear terms. Proc. R. Soc. Edinb. A 134 (2004), 249-258.
20 C. Dellacherie and P. A. Meyer. Probabilités et potentiel (Paris: Hermann, 1983).
21 W. Dong and L. Mei. Multiple solutions for an indefinite superlinear elliptic problem on $\mathbb{R}^{N}$. Nonlin. Analysis 73 (2010), 2056-2070.
22 Y. Du. Multiplicity of positive solutions for an indefinite superlinear elliptic problem on $\mathbb{R}^{N}$. Annales Inst. H. Poincaré Analyse Non Linéaire 21 (2004), 657-672.
23 A. Einstein. Ideas and opinions (New York: Crown Trade Paperbacks, 1954).
24 A. Floer and A. Weinstein. Nonspreading wave packets for the cubic Schrödinger equation with a bounded potential. J. Funct. Analysis 69 (1986), 397-408.
25 M. F. Furtado, L. A. Maia and E. S. Medeiros. Positive and nodal solutions for a nonlinear Schrödinger equation with indefinite potential. Adv. Nonlin. Studies 8 (2008), 353-373.
26 A. Hasegawa and Y. Kodama. Solitons in optical communications (Academic Press, 1995).
27 R. Kajikiya. A critical point theorem related to the symmetric mountain pass lemma and its applications to elliptic equations. J. Funct. Analysis 225 (2005), 352-370.
28 R. Kajikiya. Multiple solutions of sublinear Lane-Emden elliptic equations. Calc. Var. PDEs 26 (2006), 29-48.

32 B. A. Malomed. Variational methods in nonlinear fiber optics and related fields. Prog. Opt 43 (2002), 69-191.
33 Y. G. Oh. Existence of semi-classical bound states of nonlinear Schrödinger equations with potentials of the class $\left(V_{a}\right)$. Commun. PDEs 13 (1988), 1499-1519.

35 P. H. Rabinowitz. On a class of nonlinear Schrödinger equations. Z. Angew. Math. Phys. 43 (1992), 270-291.
36 C. Sulem and P.-L. Sulem. The nonlinear Schrödinger equation. Self-focusing and wave collapse. Applied Mathematical Sciences, vol. 139 (Springer, 1999).
37 K. Taira and K. Umezu. Positive solutions of sublinear elliptic boundary value problems. Nonlin. Analysis 29 (1997), 711-761.
38 H. Tehrani. Infinitely many solutions for indefinite semilinear elliptic equations without symmetry. Commun. PDEs 21 (1996), 541-557.
$39 \quad$ H. Tehrani. Existence results for an indefinite unbounded perturbation of a resonant Schrödinger equation. J. Diff. Eqns 236 (2007), 1-28.
40 M. Willem. Minimax theorems. Progress in Nonlinear Differential Equations and Their Applications, vol. 24 (Birkhäuser, 1996).
41 T. F. Wu. On semilinear elliptic equations involving concave-convex nonlinearities and sign-changing weight function. J. Math. Analysis Applic. 318 (2006), 253-276.
42 T. F. Wu. Multiple positive solutions for a class of concave-convex elliptic problems in $\mathbb{R}^{N}$ involving sign-changing weight. J. Funct. Analysis 258 (2010), 99-131.
43 V. E. Zakharov. Collapse and self-focusing of Langmuir waves. In Basic plasma physics $I I$ (ed. A. A. Galeev and R. N. Sudan). Handbook of Plasma Physics, vol. 2, pp. 81-121 (Elsevier, 1984).
44 Q. Zhang and Q. Wang. Multiple solutions for a class of sublinear Schrödinger equations. J. Math. Analysis Applic. 389 (2012), 511-518.

